

Lecture 4

Numerical methods for combustion research

Course *Block-structured Adaptive Finite Volume Methods in C++*

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Outline

Complex geometry

- Boundary aligned meshes
- Cartesian techniques
- Implicit geometry representation
- Accuracy / verification
- Implementation

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- Equations
- Upwind schemes for combustion
- Tests with one-step chemistry
- Shock-induced combustion with real chemistry

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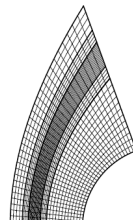
Combustion

- Equations
- Upwind schemes for combustion
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- Shock-induced combustion with real chemistry

SAMR on boundary aligned meshes

Analytic or stored geometric mapping of the coordinates
(graphic from [Yamaleev and Carpenter, 2002])

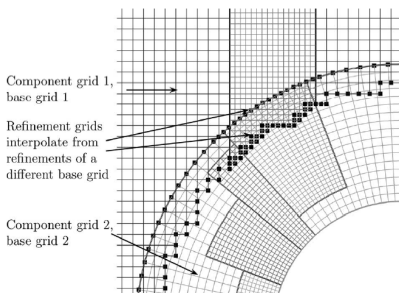
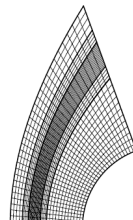
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Overlapping adaptive meshes
[Henshaw and Schwendeman, 2003],
[Meakin, 1995]

- ▶ Idea is to use a non-Cartesian structured grids only near boundary
- ▶ Very suitable for moving objects with boundary layers
- ▶ Interpolation between meshes is usually non-conservative

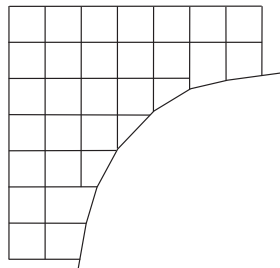
Cut-cell techniques

Accurate embedded boundary method

$$V_j^{n+1} \mathbf{Q}_j^{n+1} = V_j^n \mathbf{Q}_j^n - \Delta t \left(A_{j+1/2}^{n+1/2} \mathbf{F}(\mathbf{Q}, j) - A_{j-1/2}^{n+1/2} \mathbf{F}(\mathbf{Q}, j-1) \right)$$

Methods that represent the boundary sharply:

- ▶ Cut-cell approach constructs appropriate finite volumes
- ▶ Conservative by construction. Correct boundary flux



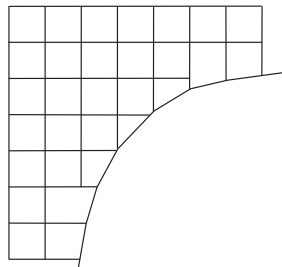
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Methods that represent the boundary sharply:

- ▶ Cut-cell approach constructs appropriate finite volumes
- ▶ Conservative by construction. Correct boundary flux
- ▶ Key question: How to avoid small-cell time step restriction in explicit methods?
 - ▶ Cell merging: [Quirk, 1994a]
- ▶ Usually explicit geometry representation used [Aftosmis, 1997], but can also be implicit, cf. [Nourgaliev et al., 2003], [Murman et al., 2003]



Embedded boundary techniques

Volume of fluid methods that resemble a cut-cell technique on purely Cartesian mesh

- ▶ Redistribution of boundary flux achieves conservation and bypasses time step restriction: [Pember et al., 1999], [Berger and Helzel, 2002]

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Methods that diffuse the boundary in one cell (good overview in [Mittal and Iaccarino, 2005]):

- ▶ Related to the immersed boundary method by Peskin, cf. [Roma et al., 1999]
- ▶ Boundary prescription often by internal ghost cell values, cf. [Tseng and Ferziger, 2003]
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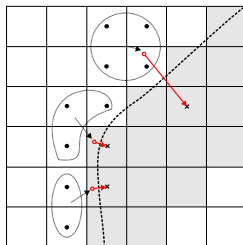
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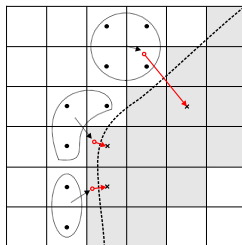
K. J. Richards et al., On the use of the immersed boundary method for engine modeling

Level-set method for boundary embedding



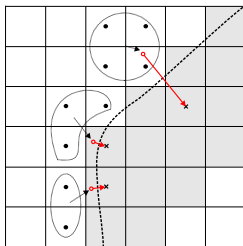
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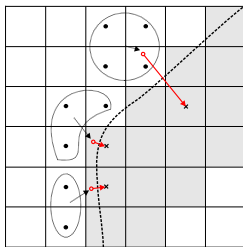
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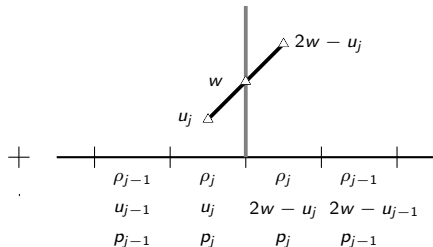
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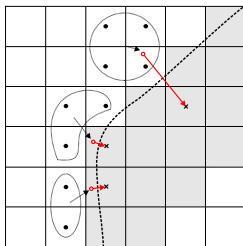
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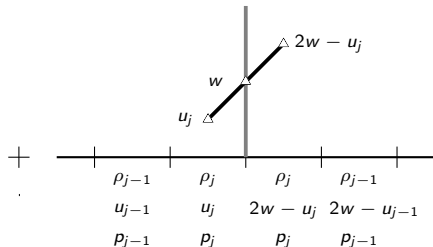
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Velocity in ghost cells

$$\begin{aligned}\mathbf{u}' &= (2\mathbf{w} \cdot \mathbf{n} - \mathbf{u} \cdot \mathbf{n})\mathbf{n} + (\mathbf{u} \cdot \mathbf{t})\mathbf{t} \\ &= 2((\mathbf{w} - \mathbf{u}) \cdot \mathbf{n})\mathbf{n} + \mathbf{u}\end{aligned}$$



Closest point transform algorithm

The signed distance φ to a surface $\mathcal{I}$ satisfies the eikonal equation [Sethian, 1999]

$$|\nabla\varphi| = 1 \quad \text{with} \quad \varphi|_{\mathcal{I}} = 0$$

Solution smooth but non-differentiable across characteristics.

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Distance computation trivial for non-overlapping elementary shapes but difficult to do efficiently for triangulated surface meshes:

- ▶ Geometric solution approach with closest-point-transform algorithm [Mauch, 2003]

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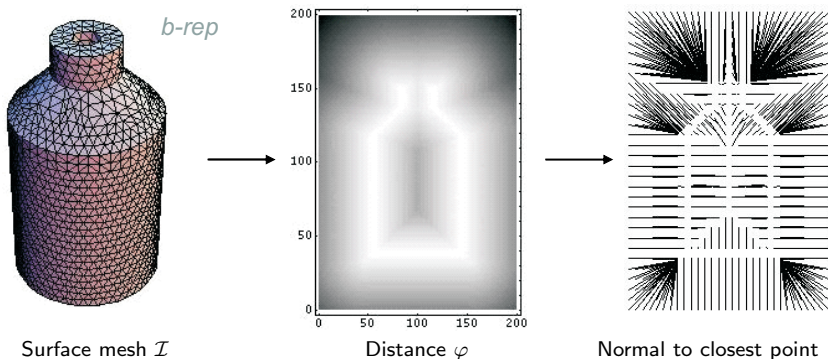
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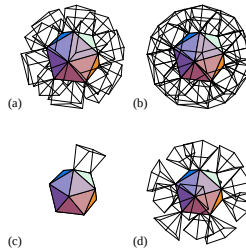
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The characteristic / scan conversion algorithm

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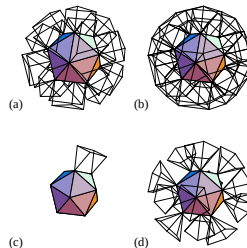
Characteristic polyhedra for faces, edges, and vertices



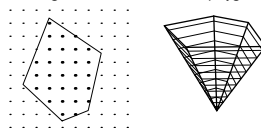
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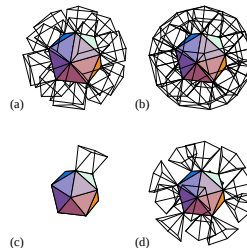
Slicing and scan conversion of apolygon



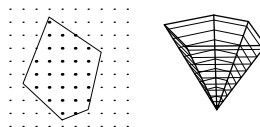
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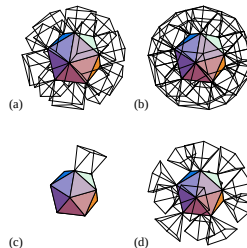
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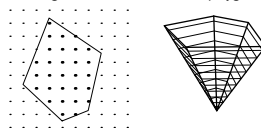
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 - ▶ $O(m)$ to build the b-rep and the polyhedra.
 - ▶ $O(n)$ to scan convert the polyhedra and compute the distance, etc.

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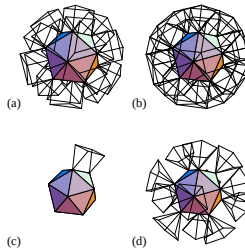
Slicing and scan conversion of a polygon



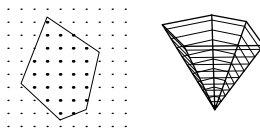
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4. Problem reduction by evaluation only within specified max. distance

Characteristic polyhedra for faces, edges, and vertices



Slicing and scan conversion of apolygon



[Mauch, 2003], see also
[Deiterding et al., 2006]

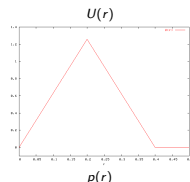
Accuracy test: stationary vortex

Construct non-trivial *radially symmetric* and *stationary* solution by balancing hydrodynamic pressure and centripetal force per volume element, i.e.

$$\frac{d}{dr}p(r) = \rho(r)\frac{U(r)^2}{r}$$

For $\rho_0 \equiv 1$ and the velocity field

$$U(r) = \alpha \cdot \begin{cases} 2r/R & \text{if } 0 < r < R/2, \\ 2(1 - r/R) & \text{if } R/2 \leq r \leq R, \\ 0 & \text{if } r > R, \end{cases}$$



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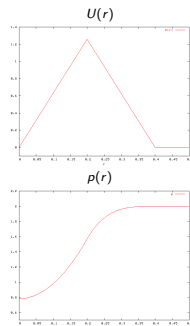
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one gets with boundary condition $p(R) = p_0 \equiv 2$ the pressure distribution

$$p(r) = p_0 + 2\rho_0\alpha^2 \cdot \begin{cases} r^2/R^2 + 1 - 2\log 2 & \text{if } 0 < r < R/2, \\ r^2/R^2 + 3 - 4r/R + 2\log(r/R) & \text{if } R/2 \leq r \leq R, \\ 0 & \text{if } r > R. \end{cases}$$



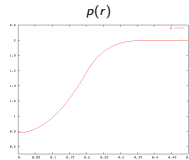
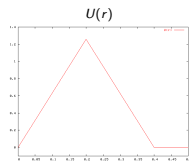
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Entire solution for Euler equations reads

$$\rho(x_1, x_2, t) = \rho_0, \quad u_1(x_1, x_2, t) = -U(r) \sin \phi, \quad u_2(x_1, x_2, t) = U(r) \cos \phi, \quad p(x_1, x_2, t) = p(r)$$

$$\text{for all } t \geq 0 \text{ with } r = \sqrt{(x_1 - x_{1,c})^2 + (x_2 - x_{2,c})^2} \text{ and } \phi = \arctan \frac{x_2 - x_{2,c}}{x_1 - x_{1,c}}$$

Stationary vortex: results

Compute one full rotation, Roe solver, embedded slip wall boundary conditions
 $x_{1,c} = 0.5$, $x_{2,c} = 0.5$, $R = 0.4$, $t_{end} = 1$, $\Delta h = \Delta x_1 = \Delta x_2 = 1/N$, $\alpha = R\pi$

No embedded boundary

N	Wave Propagation		Godunov Splitting	
	Error	Order	Error	Order
20	0.0111235		0.0182218	
40	0.0037996	1.55	0.0090662	1.01
80	0.0013388	1.50	0.0046392	0.97
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Marginal shear flow along embedded boundary, $\alpha = R\pi$, $R_G = R$, $U_W = 0$

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20	0.0120056		0.0079236	0.0144203		0.0020241
40	0.0035074	1.78	0.0011898	0.0073070	0.98	0.0001300
80	0.0014193	1.31	0.0001588	0.0038401	0.93	-0.0001036
160	0.0005032	1.50	5.046e-05	0.0018988	1.02	-2.783e-06

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Major shear flow along embedded boundary, $\alpha = R\pi$, $R_G = R/2$, $U_W = 0$

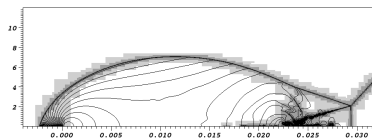
N	Wave Propagation			Godunov Splitting		
	Error	Order	Mass loss	Error	Order	Mass loss
20	0.0423925		0.0423925	0.0271446		0.0271446
40	0.0358735	0.24	0.0358735	0.0242260	0.16	0.0242260
80	0.0212340	0.76	0.0212340	0.0128638	0.91	0.0128638
160	0.0121089	0.81	0.0121089	0.0070906	0.86	0.0070906

Verification: shock reflection

- ▶ Reflection of a Mach 2.38 shock in nitrogen at 43° wedge
- ▶ 2nd order MUSCL scheme with Roe solver, 2nd order multidimensional wave propagation method

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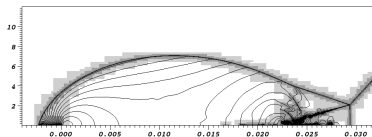
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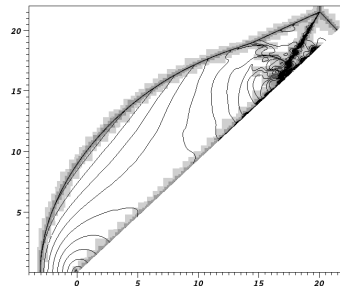
Cartesian base grid 360×160 cells on domain of $36 \text{ mm} \times 16 \text{ mm}$ with up to 3 refinement levels with $r_l = 2, 4, 4$ and $\Delta x_{1,2} = 3.125 \mu\text{m}$, 38 h CPU

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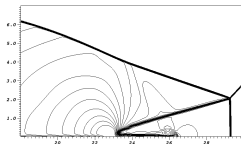


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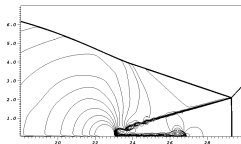


GFM base grid 390×330 cells on domain of $26 \text{ mm} \times 22 \text{ mm}$ with up to 3 refinement levels with $r_l = 2, 4, 4$ and $\Delta x_{e,1,2} = 2.849 \mu\text{m}$, 200 h CPU

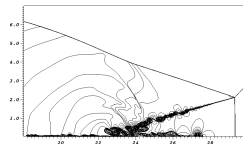
Shock reflection: SAMR solution for Euler equations



$\Delta x = 25$ mm

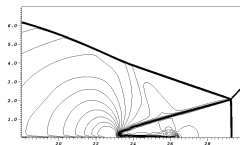


$\Delta x = 12.5$ mm

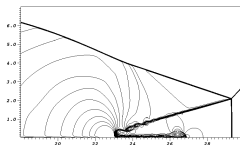


$\Delta x = 3.125$ mm

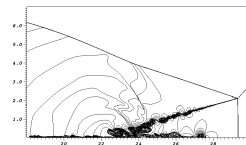
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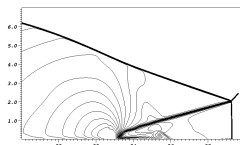
$\Delta x = 25$ mm



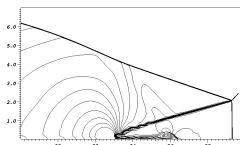
$\Delta x = 12.5$ mm



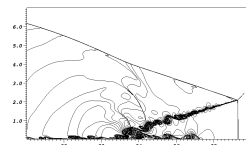
$\Delta x = 3.125$ mm



$\Delta x_e = 22.8$ mm

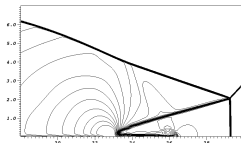


$\Delta x_e = 11.4$ mm

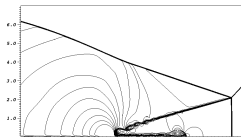


$\Delta x_e = 2.849$ mm

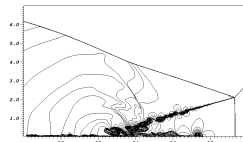
Shock reflection: SAMR solution for Euler equations



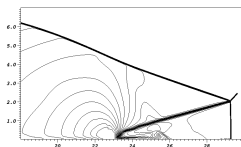
$\Delta x = 25$ mm



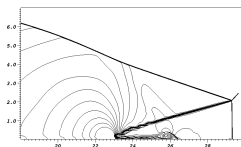
$\Delta x = 12.5$ mm



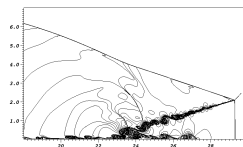
$\Delta x = 3.125$ mm



$\Delta x_e = 22.8$ mm

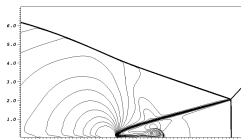


$\Delta x_e = 11.4$ mm

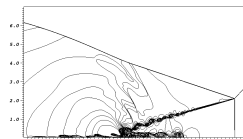


$\Delta x_e = 2.849$ mm

2nd order MUSCL scheme
with Van Leer FVS, dimensional
splitting



$\Delta x = 12.5$ mm



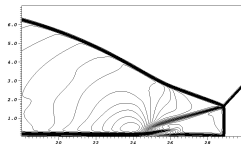
$\Delta x = 3.125$ mm

Shock reflection: solution for Navier-Stokes equations

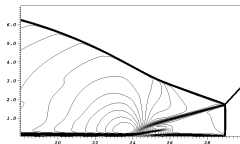
- ▶ No-slip boundary conditions enforced
- ▶ Conservative 2nd order centered differences to approximate stress tensor and heat flow

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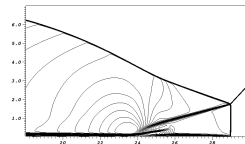
- ▶ No-slip boundary conditions enforced
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$\Delta x = 50$ mm



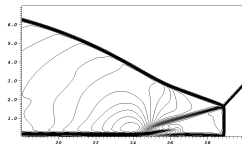
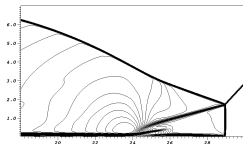
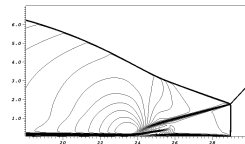
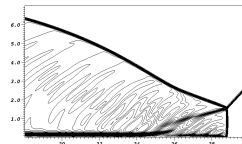
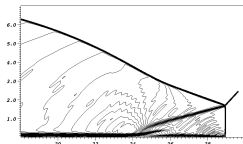
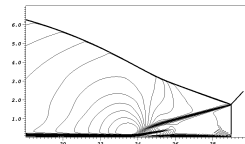
$\Delta x = 25$ mm



$\Delta x = 12.5$ mm, SAMR

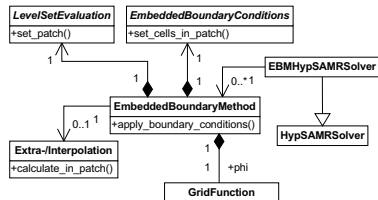
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 $\Delta x = 50$ mm $\Delta x = 25$ mm $\Delta x = 12.5$ mm, SAMR $\Delta x_e = 45.6$ mm $\Delta x_e = 22.8$ mm $\Delta x_e = 11.4$ mm, SAMR

Embedded boundary method

- ▶ Multiple independent EmbeddedBoundaryMethod objects possible
- ▶ Specialization of GFM boundary conditions



Embedded boundary method

- Multiple independent EmbeddedBoundaryMethod objects possible
- Specialization of GFM boundary conditions
- The generic embedded boundary method is implemented in **GhostFluidMethod<VectorType, dim >** and has a **GFMLevelSet<DataType, dim >** and **GFMBoundary<VectorType, dim >** object.

`code/amroc/doc/html/amr/classGhostFluidMethod.html` `code/amroc/doc/html/amr/classGFMLevelSet.html`

`code/amroc/doc/html/amr/classGFMBoundary_3_01VectorType_00_012_01_4.html`

- Multiple **GhostFluidMethod<VectorType, dim >** can be registered with **AMRGFMSolver<VectorType, FixupType, FlagType, dim >** and are called as part of the boundary condition setting routine.

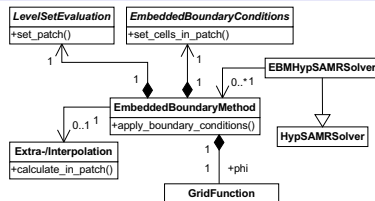
`code/amroc/doc/html/amr/classAMRGFMSolver.html`

- Interface classes to **GFMLevelSet<DataType, dim >** and **GFMBoundary<VectorType, dim >** are available `amroc/amr/F77Interfaces` and in `amroc/amr/Interfaces`

`code/amroc/doc/html/amr/classF77GFMBoundary.html` `code/amroc/doc/html/amr/classSchemeGFMBoundary.html` make the approach available to all current solvers

- For instance `code/amroc/doc/html/clp/ClpStdGFMProblem_8h.html` or

`code/amroc/doc/html/weno/WENOStdGFMProblem_8h.html` in `Problem.h` uses `AMRGFMSolver<>`



Outline

Complex geometry

- Boundary aligned meshes
- Cartesian techniques
- Implicit geometry representation
- Accuracy / verification
- Implementation

Combustion

- Equations
- Upwind schemes for combustion
- Tests with one-step chemistry
- Shock-induced combustion with real chemistry

Governing equations for premixed combustion

Euler equations with reaction terms

$$\begin{aligned}\frac{\partial \rho_i}{\partial t} + \frac{\partial}{\partial x_n}(\rho_i u_n) &= \dot{\omega}_i, \quad i = 1, \dots, K \\ \frac{\partial}{\partial t}(\rho u_k) + \frac{\partial}{\partial x_n}(\rho u_k u_n + \delta_{kn} p) &= 0, \quad k = 1, \dots, d \\ \frac{\partial}{\partial t}(\rho E) + \frac{\partial}{\partial x_n}(u_n(\rho E + p)) &= 0\end{aligned}$$

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Ideal gas law and Dalton's law for gas-mixtures

$$p(\rho_1, \dots, \rho_K, T) = \sum_{i=1}^K p_i = \sum_{i=1}^K \rho_i \frac{\mathcal{R}}{W_i} T = \rho \frac{\mathcal{R}}{W} T \quad \text{with} \quad \sum_{i=1}^K \rho_i = \rho, Y_i = \frac{\rho_i}{\rho}$$

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Caloric equation

$$h(Y_1, \dots, Y_K, T) = \sum_{i=1}^K Y_i h_i(T) \quad \text{with} \quad h_i(T) = h_i^0 + \int_0^T c_{pi}(s) ds$$

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Computation of $T = T(\rho_1, \dots, \rho_K, e)$ from implicit equation

$$\sum_{i=1}^K \rho_i h_i(T) - \mathcal{R} T \sum_{i=1}^K \frac{\rho_i}{W_i} - \rho e = 0$$

for *thermally perfect* gases with $\gamma_i(T) = c_{pi}(T)/c_{vi}(T)$

Chemistry

Arrhenius-Kinetics:

$$\dot{\omega}_i = \sum_{j=1}^M (\nu_{ji}^r - \nu_{ji}^f) \left[k_j^f \prod_{n=1}^K \left(\frac{\rho_n}{W_n} \right)^{\nu_{jn}^f} - k_j^r \prod_{n=1}^K \left(\frac{\rho_n}{W_n} \right)^{\nu_{jn}^r} \right] \quad i = 1, \dots, K$$

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- ▶ Parsing of mechanisms with Chemkin-II
- ▶ Evaluation of $\dot{\omega}_i$ with automatically generated optimized Fortran-77 functions in the line of Chemkin-II

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- ▶ Parsing of mechanisms with Chemkin-II
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Integration of reaction rates: ODE integration in $\mathcal{S}^{(\cdot)}$ for Euler equations with chemical reaction

- ▶ Standard implicit or semi-implicit ODE-solver subcycles within each cell
- ▶ ρ_i , e , u_k remain unchanged!

$$\partial_t \rho_i = W_i \dot{\omega}_i(\rho_1, \dots, \rho_K, T) \quad i = 1, \dots, K$$

Use Newton or bisection method to compute T iteratively.

Hyperbolicity of the homogeneous equations

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with

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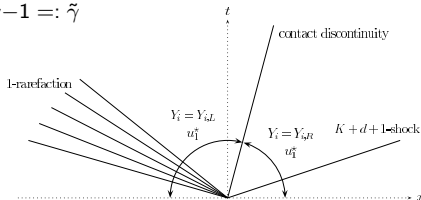
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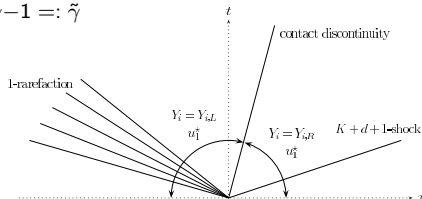
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Empirical argument: $\partial_t Y_i + u_1 \partial_x Y_i = 0$

Steger-Warming flux vector splitting

$$\mathbf{f}(\mathbf{q})^+ = \mathbf{R}\mathbf{\Lambda}^+\mathbf{R}^{-1}\mathbf{q} \quad \mathbf{f}(\mathbf{q})^- = \mathbf{R}\mathbf{\Lambda}^-\mathbf{R}^{-1}\mathbf{q}$$

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$$\delta^\pm = \lambda_2^\pm = \dots = \lambda_{K+d}^\pm,$$

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$$\xi^\pm = \lambda_1^\pm - \lambda_{K+d+1}^\pm.$$

The corresponding stability condition is

$$C_{CFL}^{SW} := \max_{j \in \mathbb{Z}} (|u_{1,j}| + a_j) \leq 1$$

[Grossmann and Cinella, 1990, Larroturou and Fezoui, 1989, Liu and Vinokur, 1989]

Steger-Warming flux vector splitting

$$\mathbf{f}(\mathbf{q})^+ = \mathbf{R}\mathbf{\Lambda}^+\mathbf{R}^{-1}\mathbf{q}$$

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$$\mathbf{f}^\pm(\mathbf{q}) = \mathbf{A}^\pm(\mathbf{q})\mathbf{q} = \frac{\rho}{2\gamma} \begin{pmatrix} Y_1\tau^\pm \\ \vdots \\ Y_K\tau^\pm \\ u_1\tau^\pm - a\xi^\pm \\ u_2\tau^\pm \\ \vdots \\ u_d\tau^\pm \\ H\tau^\pm - u_1a\xi^\pm - 2\delta^\pm a^2 \end{pmatrix} \quad \text{with}$$

$$\delta^\pm = \lambda_2^\pm = \dots = \lambda_{K+d}^\pm,$$

$$\tau^\pm = \lambda_1^\pm + 2\delta^\pm(\gamma - 1) + \lambda_{K+d+1}^\pm,$$

$$\xi^\pm = \lambda_1^\pm - \lambda_{K+d+1}^\pm.$$

The corresponding stability condition is

$$C_{CFL}^{SW} := \max_{j \in \mathbb{Z}} (|u_{1,j}| + a_j) \leq 1$$

[Grossmann and Cinella, 1990, Larrouturou and Fezoui, 1989, Liu and Vinokur, 1989]

[code/amroc/doc/html/clp/rp1eurhokswg.f_fsourc.html](http://code.amroc/doc/html/clp/rp1eurhokswg.f_fsourc.html)

Van Leer flux vector splitting

$$\mathbf{f}^{\pm}(\mathbf{q}) = \pm \frac{\rho}{4a} (u_1 \pm a)^2 \begin{bmatrix} Y_1 \\ \vdots \\ Y_K \\ u_1 - (u_1 \mp 2a)/\gamma \\ u_2 \\ \vdots \\ u_d \\ H - \zeta(u_1 \mp a)^2 \end{bmatrix} \quad \text{with}$$

$$H = h + \frac{\mathbf{u}^2}{2} ,$$

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$$\mathbf{f}(\mathbf{q})^+ = \mathbf{f}(\mathbf{q}), \quad \mathbf{f}(\mathbf{q})^- = 0 \quad \text{if } u_1 \geq a$$

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Stability condition

$$C_{CFL}^{VL} := \max_{j \in \mathbb{Z}} [(|u_{1,j}| + a_j) \Pi_j] \frac{\Delta t}{\Delta x} \leq 1$$

$$\text{with } \Pi_j = \begin{cases} \frac{\gamma_j + 3}{2\gamma_j + u_{1,j}(3 - \gamma_j)/a_j} & \text{if } |u_{1,j}| < a_j, \\ 1 & \text{otherwise.} \end{cases}$$

[Shuen et al., 1990, Liu and Vinokur, 1989, Larrouturou and Fezoui, 1989, Grossmann and Cinella, 1990]

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Roe solver

Roe averages $\hat{\rho}$, $\hat{u}$, $\hat{v}$, $\hat{H}$, $\hat{W}$, $\hat{T}$, $\hat{h}_i$, $\hat{e}_i$, $\hat{Y}_i$

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$$\Delta \mathbf{f} := \mathbf{f}(\mathbf{q}_R) - \mathbf{f}(\mathbf{q}_L) = \sum_{m=1}^M a_m \lambda_m(\hat{\mathbf{q}}) \mathbf{r}_m(\hat{\mathbf{q}}) \quad \text{with} \quad \Delta \mathbf{q} := \mathbf{q}_R - \mathbf{q}_L = \sum_{m=1}^M a_m \mathbf{r}_m(\hat{\mathbf{q}}) .$$

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With matrix of right eigenvectors

$$\mathbf{R}_1(\mathbf{q}) = \begin{bmatrix} Y_1 & 1 & 0 & \dots & 0 & 0 & 0 & Y_1 \\ & 0 & & & & & & \\ \vdots & \vdots & & \ddots & \vdots & \vdots & \vdots & \vdots \\ & & & & 0 & & & \\ Y_K & 0 & \dots & 0 & 1 & 0 & 0 & Y_K \\ u_1 - a & u_1 & \dots & & u_1 & 0 & 0 & u_1 + a \\ u_2 & u_2 & \dots & & u_2 & 1 & 0 & u_2 \\ u_3 & u_3 & \dots & & u_3 & 0 & 1 & u_3 \\ H - u_1 a & \mathbf{u}^2 - \frac{\phi_1}{\bar{\gamma}} & \dots & & \mathbf{u}^2 - \frac{\phi_K}{\bar{\gamma}} & u_2 & u_3 & H + u_1 a \end{bmatrix}$$

and evaluating $\mathbf{R}^{-1}(\hat{\mathbf{q}})\Delta \mathbf{q}$ one gets the characteristic wave strengths eventually as

$$a_1, a_{K+d+1} = \frac{\Delta p \mp \hat{\rho} \hat{a} \Delta u_1}{2 \hat{a}^2}, \quad a_{1+i} = \Delta \rho_i - \hat{Y}_i \frac{\Delta p}{\hat{a}^2}, \quad a_{K+n} = \hat{\rho} \Delta u_n.$$

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[code/amroc/doc/html/clp/rpleurhok.f_fsourc.html](http://code.amroc/doc/html/clp/rpleurhok.f_fsourc.html)

Roe solver - fixes

Mass fraction positivity: Calculate numerical fluxes of partial densities by [Larrouturou, 1991]

$$\mathbf{F}_i^*(\mathbf{q}_L, \mathbf{q}_R) = \mathbf{F}_\rho(\mathbf{q}_L, \mathbf{q}_R) \cdot \begin{cases} Y_{i,L} , & \mathbf{F}_\rho(\mathbf{q}_L, \mathbf{q}_R) > 0 , \\ Y_{i,R} , & \mathbf{F}_\rho(\mathbf{q}_L, \mathbf{q}_R) < 0 . \end{cases}$$

to ensure positivity of Y_i .

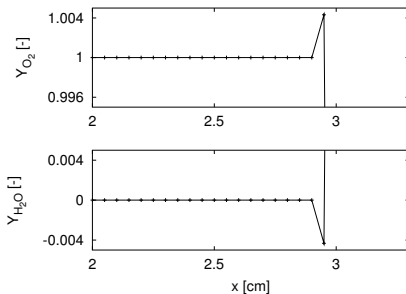
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Example: Mass fraction for a typical Riemann problem after 1 time step with the Roe solver



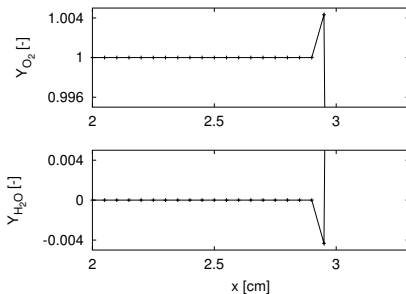
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Example: Mass fraction for a typical Riemann problem after 1 time step with the Roe solver



Further: Switch from Roe to HLL scheme near vacuum state to avoid unphysical values.

Roe solver - entropy and carbuncle fix

Entropy corrections

Roe solver - entropy and carbuncle fix

Entropy corrections [Harten, 1983]

$$1. \quad |\tilde{s}_\ell| = \begin{cases} |s_\ell| & \text{if } |s_\ell| \geq 2\eta \\ \frac{|s_\ell^2|}{4\eta} + \eta & \text{otherwise} \end{cases}$$

$$\eta = \frac{1}{2} \max_\ell \{ |s_\ell(\mathbf{q}_R) - s_\ell(\mathbf{q}_L)| \}$$

Roe solver - entropy and carbuncle fix

Entropy corrections [Harten, 1983]

[Harten and Hyman, 1983]

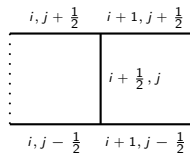
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2. Replace $|s_\ell|$ by $|\tilde{s}_\ell|$ only if
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Roe solver - entropy and carbuncle fix

Entropy corrections [Harten, 1983]
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2D modification of entropy correction (here named EF 3)
[Sanders et al., 1998]:



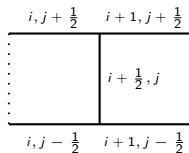
$$\tilde{\eta}_{i+1/2,j} = \max \{ \eta_{i+1/2,j}, \eta_{i,j-1/2}, \eta_{i,j+1/2}, \eta_{i+1,j-1/2}, \eta_{i+1,j+1/2} \}$$

Roe solver - entropy and carbuncle fix

Entropy corrections [Harten, 1983]
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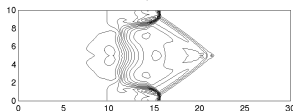
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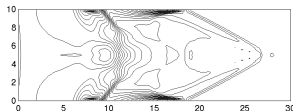


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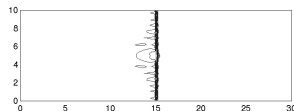
Roe + EF 1.



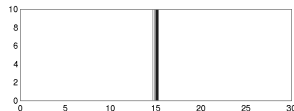
Roe + EF 2.



Exact Riemann solver



SW FVS, VL FVS, HLL, Roe + EF 3



Carbuncle phenomenon

- [Quirk, 1994b]
- Test from [Deiterding, 2003]

code/amroc/doc/html/apps/
clawpack_2applications_2euler_
_znd_22d_2Carbuncle_2src_
2Problem_8h_source.html

Riemann solver for combustion

(S1) Calculate standard Roe-averages $\hat{\rho}$, $\hat{u}_n$, $\hat{H}$, $\hat{Y}_i$, $\hat{T}$.

(S2) Compute $\hat{\gamma} := \hat{c}_p / \hat{c}_v$ with $\hat{c}_{\{p/v\}i} = \frac{1}{T_R - T_L} \int_{T_L}^{T_R} c_{\{p,v\}i}(\tau) d\tau$.

(S3) Calculate $\hat{\phi}_i := (\hat{\gamma} - 1) \left(\frac{\hat{u}^2}{2} - \hat{h}_i \right) + \hat{\gamma} R_i \hat{T}$ with standard Roe-averages $\hat{e}_i$ or $\hat{h}_i$.

(S4) Calculate $\hat{a} := \left(\sum_{i=1}^K \hat{Y}_i \hat{\phi}_i - (\hat{\gamma} - 1) \hat{u}^2 + (\hat{\gamma} - 1) \hat{H} \right)^{1/2}$.

(S5) Use $\Delta \mathbf{q} = \mathbf{q}_R - \mathbf{q}_L$ and Δp to compute the wave strengths a_m .

(S6) Calculate $\mathcal{W}_1 = a_1 \hat{\mathbf{r}}_1$, $\mathcal{W}_2 = \sum_{\iota=2}^{K+d} a_\iota \hat{\mathbf{r}}_\iota$, $\mathcal{W}_3 = a_{K+d+1} \hat{\mathbf{r}}_{K+d+1}$.

(S7) Evaluate $s_1 = \hat{u}_1 - \hat{a}$, $s_2 = \hat{u}_1$, $s_3 = \hat{u}_1 + \hat{a}$.

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- (S8) Evaluate $\rho_{L/R}^*$, $u_{1,L/R}^*$, $e_{L/R}^*$, $a_{L/R}^*$ from $\mathbf{q}_L^* = \mathbf{q}_L + \mathcal{W}_1$ and $\mathbf{q}_R^* = \mathbf{q}_R - \mathcal{W}_3$.
- (S9) If $\rho_{L/R}^* \leq 0$ or $e_{L/R}^* \leq 0$ use $\mathbf{F}_{HLL}(\mathbf{q}_L, \mathbf{q}_R)$ and go to (S12).

Riemann solver for combustion

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- (S10) Entropy correction: Evaluate $|\tilde{s}_\iota|$.

$$\mathbf{F}_{Roe}(\mathbf{q}_L, \mathbf{q}_R) = \frac{1}{2} (\mathbf{f}(\mathbf{q}_L) + \mathbf{f}(\mathbf{q}_R) - \sum_{\iota=1}^3 |\tilde{s}_\iota| \mathcal{W}_\iota)$$

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- (S8) Evaluate $\rho_{L/R}^*$, $u_{1,L/R}^*$, $e_{L/R}^*$, $a_{L/R}^*$ from $\mathbf{q}_L^* = \mathbf{q}_L + \mathcal{W}_1$ and $\mathbf{q}_R^* = \mathbf{q}_R - \mathcal{W}_3$.
- (S9) If $\rho_{L/R}^* \leq 0$ or $e_{L/R}^* \leq 0$ use $\mathbf{F}_{HLL}(\mathbf{q}_L, \mathbf{q}_R)$ and go to (S12).
- (S10) Entropy correction: Evaluate $|\tilde{s}_l|$.

$$\mathbf{F}_{Roe}(\mathbf{q}_L, \mathbf{q}_R) = \frac{1}{2} (\mathbf{f}(\mathbf{q}_L) + \mathbf{f}(\mathbf{q}_R) - \sum_{l=1}^3 |\tilde{s}_l| \mathcal{W}_l)$$
- (S11) Positivity correction: Replace $\mathbf{F}_i$ by

$$\mathbf{F}_i^* = \mathbf{F}_\rho \cdot \begin{cases} \mathbf{Y}_i^l, & \mathbf{F}_\rho \geq 0, \\ \mathbf{Y}_i^r, & \mathbf{F}_\rho < 0. \end{cases}$$
- (S12) Evaluate maximal signal speed by $S = \max(|s_1|, |s_3|)$.

Riemann solver for combustion

(S1) Calculate standard Roe-averages $\hat{\rho}$, $\hat{u}_n$, $\hat{H}$, $\hat{Y}_i$, $\hat{T}$.

(S2) Compute $\hat{\gamma} := \hat{c}_p / \hat{c}_v$ with $\hat{c}_{\{p,v\}i} = \frac{1}{T_R - T_L} \int_{T_L}^{T_R} c_{\{p,v\}i}(\tau) d\tau$.

(S3) Calculate $\hat{\phi}_i := (\hat{\gamma} - 1) \left(\frac{\hat{u}^2}{2} - \hat{h}_i \right) + \hat{\gamma} R_i \hat{T}$ with standard Roe-averages $\hat{e}_i$ or $\hat{h}_i$.

(S4) Calculate $\hat{a} := \left(\sum_{i=1}^K \hat{Y}_i \hat{\phi}_i - (\hat{\gamma} - 1) \hat{u}^2 + (\hat{\gamma} - 1) \hat{H} \right)^{1/2}$.

(S5) Use $\Delta \mathbf{q} = \mathbf{q}_R - \mathbf{q}_L$ and Δp to compute the wave strengths a_m .

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`code/amroc/doc/html/clp/rpleurhokefix.f_fsourc.html`
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ZND Detonation Model with Simplified Chemistry

Assume a stationary 1D detonation with irreversible reaction



with energy release $q_0 = -\Delta h^0$ and $k^f(T) = K \exp(-E_A/\mathcal{R}T)$.

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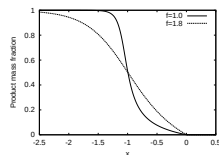
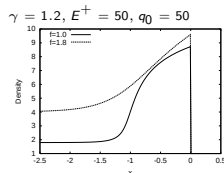
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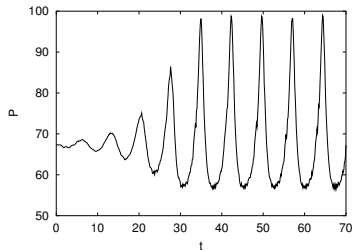


Unstable ZND Detonation

$$\gamma = 1.2, E^+ = 50, q_0 = 50, f = 1.6, CFL = 0.9$$

quasi-stationary

moving

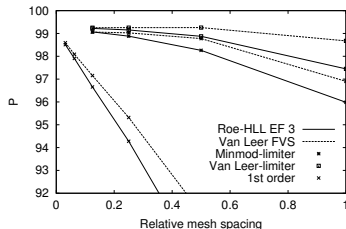
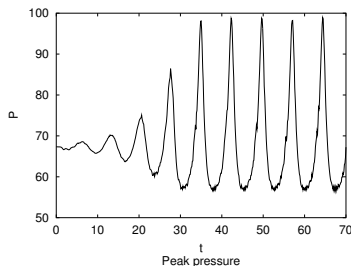


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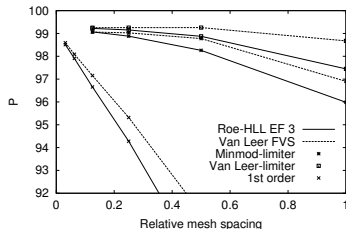
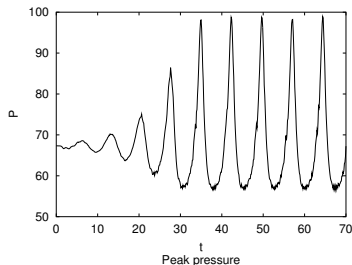
moving



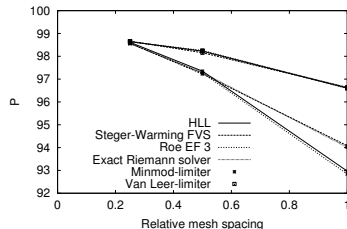
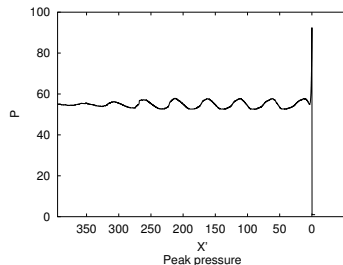
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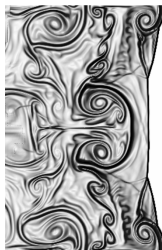
moving



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Comparison of FV Schemes: MUSCL, Van Albada-limiter

$$\gamma = 1.2, E = 50, q_0 = 50, f = 3.0, 40 \text{ Pts}/L_{1/2}$$

 T  Z 

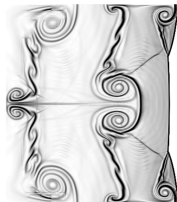
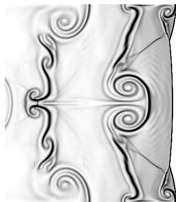
Van Leer FVS, DimSplit

Roe-HLL EF 3, DimSplit

Roe-HLL EF 3,
WaveProp

Comparison of FV Schemes - II

$$\gamma = 1.2, E = 10, q_0 = 50, f = 1.2, 40 \text{ Pts}/L_{1/2}$$

 T  Z 

Van Leer FVS, DimSplit

Roe-HLL EF 3*-H,
DimSplitRoe-HLL EF 3*-H,
WaveProp

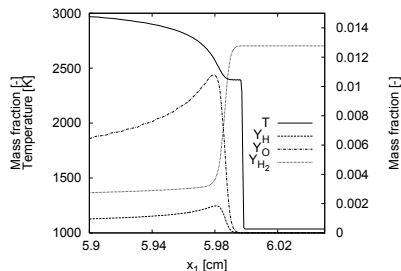
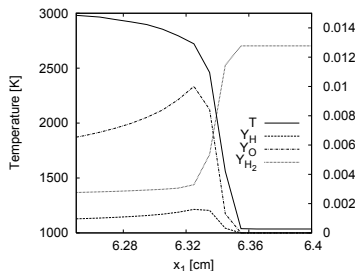
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Detonations - motivation for SAMR

- ▶ Extremely high spatial resolution in reaction zone necessary.

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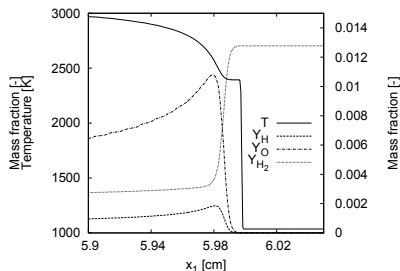
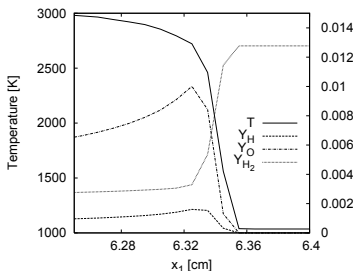
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- ▶ Minimal spatial resolution: $7 - 8 \text{ Pts}/l_{ig} \rightarrow \Delta x_1 \approx 0.2 - 0.175 \text{ mm}$



Approximation of $\text{H}_2 : \text{O}_2$ detonation at $\sim 1.5 \text{ Pts}/l_{ig}$ (left) and $\sim 24 \text{ Pts}/l_{ig}$ (right)

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- ▶ Uniform grids for typical geometries: $> 10^7 \text{ Pts}$ in 2D, $> 10^9 \text{ Pts}$ in 3D $\rightarrow$ Self-adaptive finite volume method (AMR)



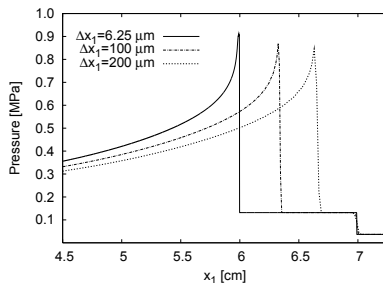
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Detonation ignition in a shock tube

- ▶ Shock-induced detonation ignition of $\text{H}_2 : \text{O}_2 : \text{Ar}$ mixture at molar ratios 2:1:7 in closed 1d shock tube
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Detonation ignition in a shock tube

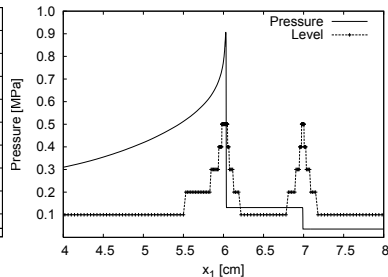
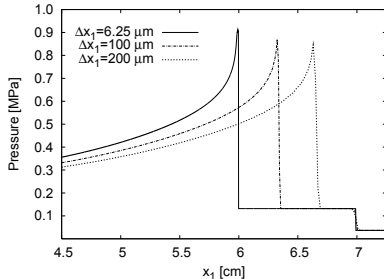
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Left: Comparison of pressure distribution $t = 170 \mu\text{s}$ after shock reflection.

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Left: Comparison of pressure distribution $t = 170 \mu\text{s}$ after shock reflection. Right: Domains of refinement levels

Non-equilibrium mechanism for hydrogen-oxygen combustion

				A	β	E_{act}
				[cm, mol, s]		[cal mol ⁻¹]
1.	H + O ₂	→	O + OH	1.86×10^{14}	0.00	16790.
2.	O + OH	→	H + O ₂	1.48×10^{13}	0.00	680.
3.	H ₂ + O	→	H + OH	1.82×10^{10}	1.00	8900.
4.	H + OH	→	H ₂ + O	8.32×10^{09}	1.00	6950.
5.	H ₂ O + O	→	OH + OH	3.39×10^{13}	0.00	18350.
6.	OH + OH	→	H ₂ O + O	3.16×10^{12}	0.00	1100.
7.	H ₂ O + H	→	H ₂ + OH	9.55×10^{13}	0.00	20300.
8.	H ₂ + OH	→	H ₂ O + H	2.19×10^{13}	0.00	5150.
9.	H ₂ O ₂ + OH	→	H ₂ O + HO ₂	1.00×10^{13}	0.00	1800.
10.	H ₂ O + HO ₂	→	H ₂ O ₂ + OH	2.82×10^{13}	0.00	32790.
...	...		...	...	...	...
30.	OH + M	→	O + H + M	7.94×10^{19}	-1.00	103720.
31.	O ₂ + M	→	O + O + M	5.13×10^{15}	0.00	115000.
32.	O + O + M	→	O ₂ + M	4.68×10^{15}	-0.28	0.
33.	H ₂ + M	→	H + H + M	2.19×10^{14}	0.00	96000.
34.	H + H + M	→	H ₂ + M	3.02×10^{15}	0.00	0.

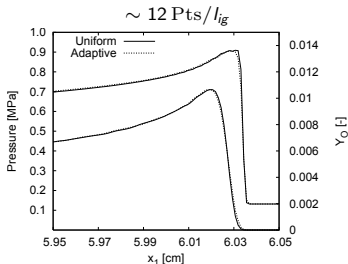
Third body efficiencies: $f(\text{O}_2) = 0.40$, $f(\text{H}_2\text{O}) = 6.50$

[Westbrook, 1982]

Detonation ignition in 1d - adaptive vs. uniform

Uniformly refined vs. dynamic adaptive simulations (Intel Xeon 3.4 GHz CPU)

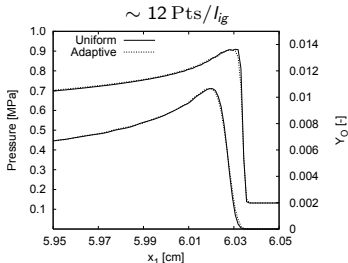
$\Delta x_1 [\mu\text{m}]$	Uniform			Adaptive			
	Cells	$t_m [\mu\text{s}]$	Time [s]	l_{max}	r_l	$t_m [\mu\text{s}]$	Time [s]
400	300	166.1	31				
200	600	172.6	90	2	2	172.6	99
100	1200	175.5	277	3	2,2	175.8	167
50	2400	176.9	858	4	2,2,2	177.3	287
25	4800	177.8	2713	4	2,2,4	177.9	393
12.5	9600	178.3	9472	5	2,2,2,4	178.3	696
6.25	19200	178.6	35712	5	2,2,4,4	178.6	1370



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Refinement criteria:

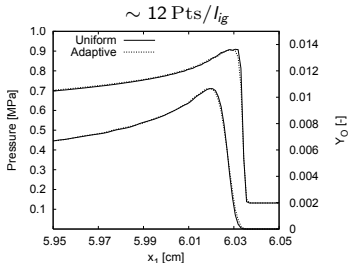
Y_i	$S_{Y_i} \cdot 10^{-4}$	$\eta_{Y_i}^r \cdot 10^{-3}$
O ₂	10.0	2.0
H ₂ O	7.8	8.0
H	0.16	5.0
O	1.0	5.0
OH	1.8	5.0
H ₂	1.3	2.0

$$\epsilon_\rho = 0.07 \text{ kg m}^{-3}, \epsilon_p = 50 \text{ kPa}$$

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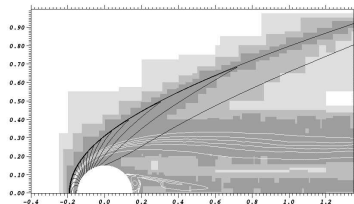
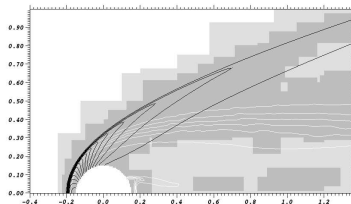
code/amroc/doc/html/apps/clawpack_2applications_2euler_chem_21d_2ReflecDet_2src_2Problem_8h_source.html

Shock-induced combustion around a sphere

- ▶ Spherical projectile of radius 1.5 mm travels with constant velocity $v_I = 2170.6 \text{ m/s}$ through $\text{H}_2 : \text{O}_2 : \text{Ar}$ mixture (molar ratios 2:1:7) at 6.67 kPa and $T = 298 \text{ K}$
- ▶ Cylindrical symmetric simulation on AMR base mesh of 70×40 cells

Shock-induced combustion around a sphere

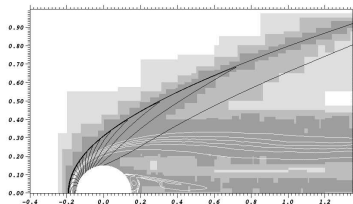
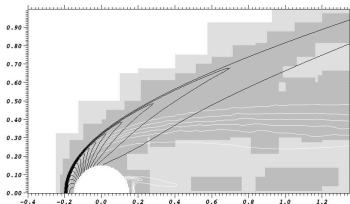
- ▶ Spherical projectile of radius 1.5 mm travels with constant velocity $v_I = 2170.6 \text{ m/s}$ through $\text{H}_2 : \text{O}_2 : \text{Ar}$ mixture (molar ratios 2:1:7) at 6.67 kPa and $T = 298 \text{ K}$
- ▶ Cylindrical symmetric simulation on AMR base mesh of 70×40 cells
- ▶ Comparison of 3-level computation with refinement factors 2,2 ($\sim 5 \text{ Pts}/l_{ig}$) and a 4-level computation with refinement factors 2,2,4 ($\sim 19 \text{ Pts}/l_{ig}$) at $t = 350 \mu\text{s}$



Iso-contours of p (black) and Y_{H_2} (white) on refinement domains for 3-level (left) and 4-level computation (right)

Shock-induced combustion around a sphere

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- ▶ Higher resolved computation captures combustion zone visibly better and at slightly different position (see below)

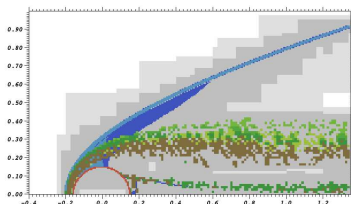


Iso-contours of p (black) and Y_{H_2} (white) on refinement domains for 3-level (left) and 4-level computation (right)

code/amroc/doc/html/apps/clawpack_2applications_2euler_chem_22d_2Sphere_2src_2Problem_8h_source.html

Combustion around a sphere - adaptation

Refinement indicators on $l = 2$ at $t = 350 \mu\text{s}$.
 Blue: ϵ_ρ , light blue: ϵ_p , green shades: $\eta_{Y_i}^r$,
 red: embedded boundary



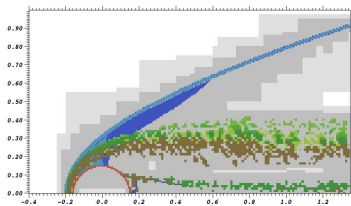
Refinement criteria:

Y_i	$S_{Y_i} \cdot 10^{-4}$	$\eta_{Y_i}^r \cdot 10^{-4}$
O ₂	10.0	4.0
H ₂ O	5.8	3.0
H	0.2	10.0
O	1.4	10.0
OH	2.3	10.0
H ₂	1.3	4.0

$$\epsilon_\rho = 0.02 \text{ kg m}^{-3}, \epsilon_p = 16 \text{ kPa}$$

Combustion around a sphere - adaptation

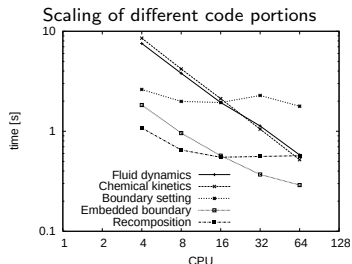
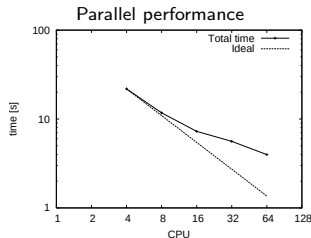
Refinement indicators on $l = 2$ at $t = 350 \mu\text{s}$.
 Blue: ϵ_ρ , light blue: ϵ_p , green shades: $\eta_{Y_i}^r$,
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