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Lecture 7 Lattice Boltzmann methods

Course Block-structured Adaptive Finite Volume Methods in C++

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Adaptive lattice Boltzmann method

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Realistic aerodynamics computations

Vehicle geometries

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Fluid-structure coupling

Rigid body dynamics Validation simulations

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Wind turbine wake aerodynamics

Mexico benchmark Simulation of wind turbine wakes Wake interaction prediction

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Construction principles				

Is based on solving the Boltzmann equation with a simplified collision operator

$$\partial_t f + \mathbf{u} \cdot \nabla f = \omega (f^{eq} - f)$$

- $\operatorname{Kn} = I_f / L \ll 1$, where I_f is replaced with Δx
- Weak compressibility and small Mach number assumed
- Assume a simplified phase space

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Weak compressibility and small Mach number assumed

Assume a simplified phase space

Equation is approximated with a splitting approach.

1.) Transport step solves $\partial_t f_{\alpha} + \mathbf{e}_{\alpha} \cdot \nabla f_{\alpha} = 0$ Operator: \mathcal{T} : $\tilde{f}_{\alpha}(\mathbf{x} + \mathbf{e}_{\alpha}\Delta t, t + \Delta t) = f_{\alpha}(\mathbf{x}, t)$

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1.) Transport step solves
$$\partial_t f_\alpha + \mathbf{e}_\alpha \cdot \nabla f_\alpha = 0$$

Operator: \mathcal{T} : $\tilde{f}_\alpha(\mathbf{x} + \mathbf{e}_\alpha \Delta t, t + \Delta t) = f_\alpha(\mathbf{x}, t)$
 $\rho(\mathbf{x}, t) = \sum_{\alpha=0}^8 f_\alpha(\mathbf{x}, t), \quad \rho(\mathbf{x}, t) u_i(\mathbf{x}, t) = \sum_{\alpha=0}^8 \mathbf{e}_{\alpha i} f_\alpha(\mathbf{x}, t)$

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Discrete velocities:

 $\mathbf{e}_0=(0,0), \mathbf{e}_1=(1,0)c, \mathbf{e}_2=(-1,0)c, \mathbf{e}_3=(0,1)c, \mathbf{e}_4=(1,1)c, ...$

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Discrete velocities:

$$\begin{split} \mathbf{e}_0 &= (0,0), \mathbf{e}_1 = (1,0)c, \mathbf{e}_2 = (-1,0)c, \mathbf{e}_3 = (0,1)c, \mathbf{e}_4 = (1,1)c, ... \\ c &= \frac{\Delta x}{\Delta t}, \text{ Physical speed of sound: } c_s = \frac{c}{\sqrt{3}} \end{split}$$

Approximation of Boltzmann equation

Is based on solving the Boltzmann equation with a simplified collision operator

$$\partial_t f + \mathbf{u} \cdot \nabla f = \omega (f^{eq} - f)$$

- $\text{Kn} = l_f / L \ll 1$, where l_f is replaced with Δx
- Weak compressibility and small Mach number assumed
- Assume a simplified phase space

Equation is approximated with a splitting approach.

1.) Transport step solves
$$\partial_t f_{\alpha} + \mathbf{e}_{\alpha} \cdot \nabla f_{\alpha} = 0$$

Operator: \mathcal{T} : $\tilde{f}_{\alpha}(\mathbf{x} + \mathbf{e}_{\alpha}\Delta t, t + \Delta t) = f_{\alpha}(\mathbf{x}, t)$
 $\rho(\mathbf{x}, t) = \sum_{\alpha=0}^{18} f_{\alpha}(\mathbf{x}, t), \quad \rho(\mathbf{x}, t) u_i(\mathbf{x}, t) = \sum_{\alpha=0}^{18} \mathbf{e}_{\alpha i} f_{\alpha}(\mathbf{x}, t)$



Discrete velocities:

$$\mathbf{e}_{\alpha} = \begin{cases} 0, & \alpha = 0, \\ (\pm 1, 0, 0)c, (0, \pm 1, 0)c, (0, 0, \pm 1)c, & \alpha = 1, \dots, 6, \\ (\pm 1, \pm 1, 0)c, (\pm 1, 0, \pm 1)c, (0, \pm 1, \pm 1)c, & \alpha = 7, \dots, 18, \end{cases}$$

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Construction principles				

2.) Collision step solves $\partial_t f_{\alpha} = \omega (f_{\alpha}^{eq} - f_{\alpha})$ Operator C:

$$f_{\alpha}(\cdot,t+\Delta t)=\tilde{f}_{\alpha}(\cdot,t+\Delta t)+\omega_{L}\Delta t\left(\tilde{f}_{\alpha}^{eq}(\cdot,t+\Delta t)-\tilde{f}_{\alpha}(\cdot,t+\Delta t)\right)$$

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2.) Collision step solves $\partial_t f_{\alpha} = \omega (f_{\alpha}^{eq} - f_{\alpha})$ Operator C:

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with equilibrium function

$$f_{\alpha}^{eq}(\rho,\mathbf{u}) = \rho t_{\alpha} \left[1 + \frac{3\mathbf{e}_{\alpha}\mathbf{u}}{c^2} + \frac{9(\mathbf{e}_{\alpha}\mathbf{u})^2}{2c^4} - \frac{3u^2}{2c^2} \right]$$

with $t_{\alpha} = \frac{1}{9} \left\{ 4, 1, 1, 1, \frac{1}{4}, \frac{1}{4}, 1, \frac{1}{4}, \frac{1}{4} \right\}$ Pressure $\delta p = \sum_{\alpha} f_{\alpha}^{eq} c_s^2 = \rho c_s^2$. Dev. stress $\Sigma_{ij} = \left(1 - \frac{\omega_L \Delta t}{2}\right) \sum_{\alpha} \mathbf{e}_{\alpha i} \mathbf{e}_{\alpha j} (f_{\alpha}^{eq} - f_{\alpha})$

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with $t_{\alpha} = \frac{1}{9} \left\{ 3, \frac{1}{2}, \frac{1}{2}, \frac{1}{2}, \frac{1}{2}, \frac{1}{2}, \frac{1}{2}, \frac{1}{2}, \frac{1}{4}, \frac$

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$$f_{\alpha}^{eq}(\rho,\mathbf{u}) = \rho t_{\alpha} \left[1 + \frac{3\mathbf{e}_{\alpha}\mathbf{u}}{c^2} + \frac{9(\mathbf{e}_{\alpha}\mathbf{u})^2}{2c^4} - \frac{3\mathbf{u}^2}{2c^2} \right]$$

with $t_{\alpha} = \frac{1}{9} \left\{ 3, \frac{1}{2}, \frac{1}{2}, \frac{1}{2}, \frac{1}{2}, \frac{1}{2}, \frac{1}{2}, \frac{1}{2}, \frac{1}{4}, \frac$

Using the third-order equilibrium function

$$f_{\alpha}^{eq}(\rho,\mathbf{u}) = \rho t_{\alpha} \left[1 + \frac{3\mathbf{e}_{\alpha}\mathbf{u}}{c^2} + \frac{9(\mathbf{e}_{\alpha}\mathbf{u})^2}{2c^4} - \frac{3\mathbf{u}^2}{2c^2} + \frac{\mathbf{e}_{\alpha}\mathbf{u}}{3c^2} \left(\frac{9(\mathbf{e}_{\alpha}\mathbf{u})^2}{2c^4} - \frac{3\mathbf{u}^2}{2c^2} \right) \right]$$

allows higher flow velocities.

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Relation to Navier-Stokes equations

Inserting a Chapman-Enskog expansion, that is,

$$f_{\alpha} = f_{\alpha}(0) + \epsilon f_{\alpha}(1) + \epsilon^2 f_{\alpha}(2) + ...$$

and using

$$\frac{\partial}{\partial t} = \epsilon \frac{\partial}{\partial t_1} + \epsilon^2 \frac{\partial}{\partial t_2} + ..., \qquad \nabla = \epsilon \nabla_1 + \epsilon^2 \nabla_2 + ...$$

into the LBM and summing over α one can show that the continuity and moment equations are recoverd to $O(\epsilon^2)$ [Hou et al., 1996]

$$\partial_t \rho + \nabla \cdot (\rho \mathbf{u}) = \mathbf{0}$$

 $\partial_t \mathbf{u} + \mathbf{u} \cdot \nabla \mathbf{u} = -\nabla \rho + \nu \nabla^2 \mathbf{u}$

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Relation to Navier-Stokes equations

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$$\partial_t \mathbf{u} + \mathbf{u} \cdot \nabla \mathbf{u} = -\nabla \rho + \nu \nabla^2 \mathbf{u}$$

Kinematic viscosity and collision time are connected by

$$\nu = \frac{1}{3} \left(\frac{\tau_L}{\Delta t} - \frac{1}{2} \right) c \Delta x$$

from which one gets with $\sqrt{3}c_{s}=\frac{\Delta x}{\Delta t}$ [Hähnel, 2004]

$$\omega_L = \tau_L^{-1} = \frac{c_s^2}{\nu + \Delta t c_s^2/2}$$

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Construction principles				

Pursue a large-eddy simulation approach with \bar{f}_{α} and \bar{f}_{α}^{eq} , i.e.

1.)
$$\overline{f}_{\alpha}(\mathbf{x} + \mathbf{e}_{\alpha}\Delta t, t + \Delta t) = \overline{f}_{\alpha}(\mathbf{x}, t)$$

2.) $\overline{f}_{\alpha}(\cdot, t + \Delta t) = \widetilde{\overline{f}}_{\alpha}(\cdot, t + \Delta t) + \frac{1}{\tau^{\star}}\Delta t \left(\widetilde{\overline{f}}_{\alpha}^{eq}(\cdot, t + \Delta t) - \widetilde{\overline{f}}_{\alpha}(\cdot, t + \Delta t) \right)$

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Pursue a large-eddy simulation approach with \overline{f}_{α} and $\overline{f}_{\alpha}^{eq}$, i.e. 1.) $\tilde{f}_{\alpha}(\mathbf{x} + \mathbf{e}_{\alpha}\Delta t, t + \Delta t) = \overline{f}_{\alpha}(\mathbf{x}, t)$ 2.) $\overline{f}_{\alpha}(\cdot, t + \Delta t) = \tilde{f}_{\alpha}(\cdot, t + \Delta t) + \frac{1}{\tau^{*}}\Delta t \left(\tilde{f}_{\alpha}^{eq}(\cdot, t + \Delta t) - \tilde{f}_{\alpha}(\cdot, t + \Delta t) \right)$ Effective viscosity: $\nu^{*} = \nu + \nu_{t} = \frac{1}{3} \left(\frac{\tau_{L}^{*}}{\Delta t} - \frac{1}{2} \right) c\Delta x$ with $\tau_{L}^{*} = \tau_{L} + \tau_{t}$

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$$ar{m{S}} = \sqrt{2\sum_{i,j}m{m{S}}_{ij}m{m{S}}_{ij}}$$

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Pursue a large-eddy simulation approach with \overline{f}_{α} and $\overline{f}_{\alpha}^{eq}$, i.e. 1.) $\tilde{f}_{\alpha}(\mathbf{x} + \mathbf{e}_{\alpha}\Delta t, t + \Delta t) = \overline{f}_{\alpha}(\mathbf{x}, t)$ 2.) $\overline{f}_{\alpha}(\cdot, t + \Delta t) = \tilde{f}_{\alpha}(\cdot, t + \Delta t) + \frac{1}{\tau^{*}}\Delta t \left(\tilde{f}_{\alpha}^{eq}(\cdot, t + \Delta t) - \tilde{f}_{\alpha}(\cdot, t + \Delta t) \right)$ Effective viscosity: $\nu^{*} = \nu + \nu_{t} = \frac{1}{3} \left(\frac{\tau_{L}^{*}}{\Delta t} - \frac{1}{2} \right) c\Delta x$ with $\tau_{L}^{*} = \tau_{L} + \tau_{t}$ Use Smagorinsky model to evaluate ν_{t} , e.g., $\nu_{t} = (C_{sm}\Delta x)^{2}\overline{S}$, where

$$ar{m{S}} = \sqrt{2\sum_{i,j}m{m{S}}_{ij}m{m{S}}_{jj}}$$

The filtered strain rate tensor $\mathbf{\bar{S}}_{ij} = (\partial_j \bar{u}_i + \partial_i \bar{u}_j)/2$ can be computed as a second moment as

$$\mathbf{\bar{S}}_{ij} = \frac{\Sigma_{ij}}{2\rho c_s^2 \tau_L^{\star} \left(1 - \frac{\omega_L \Delta t}{2}\right)} = \frac{1}{2\rho c_s^2 \tau_L^{\star}} \sum_{\alpha} \mathbf{e}_{\alpha i} \mathbf{e}_{\alpha j} (\bar{f}_{\alpha}^{eq} - \bar{f}_{\alpha})$$

Adaptive lattice Boltzmann method A	lerodynamics cases	Fluid-structure coupling	Wind turbine wake aerodynamics	Reference
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Construction principles				

Pursue a large-eddy simulation approach with $ar{f}_{lpha}$ and $ar{f}_{lpha}^{eq}$, i.e.

1.)
$$f_{\alpha}(\mathbf{x} + \mathbf{e}_{\alpha}\Delta t, t + \Delta t) = f_{\alpha}(\mathbf{x}, t)$$

2.) $\overline{f}_{\alpha}(\cdot, t + \Delta t) = \tilde{\overline{f}}_{\alpha}(\cdot, t + \Delta t) + \frac{1}{\tau^{\star}}\Delta t \left(\tilde{\overline{f}}_{\alpha}^{eq}(\cdot, t + \Delta t) - \tilde{\overline{f}}_{\alpha}(\cdot, t + \Delta t)\right)$

Effective viscosity: $\nu^{\star} = \nu + \nu_t = \frac{1}{3} \left(\frac{\tau_L^{\star}}{\Delta t} - \frac{1}{2} \right) c \Delta x$ with $\tau_L^{\star} = \tau_L + \tau_t$

Use Smagorinsky model to evaluate ν_t , e.g., $\nu_t = (C_{sm}\Delta x)^2 \bar{S}$, where

$$ar{m{S}} = \sqrt{2\sum_{i,j}m{m{S}}_{ij}m{m{S}}_{jj}}$$

The filtered strain rate tensor $\mathbf{\bar{S}}_{ij} = (\partial_j \bar{u}_i + \partial_i \bar{u}_j)/2$ can be computed as a second moment as

$$\bar{\mathbf{S}}_{ij} = \frac{\sum_{ij}}{2\rho c_s^2 \tau_L^* \left(1 - \frac{\omega_L \Delta t}{2}\right)} = \frac{1}{2\rho c_s^2 \tau_L^*} \sum_{\alpha} \mathbf{e}_{\alpha i} \mathbf{e}_{\alpha j} (\bar{f}_{\alpha}^{eq} - \bar{f}_{\alpha})$$

 τ_t can be obtained as [Yu, 2004, Hou et al., 1996]

$$\tau_t = \frac{1}{2} \left(\sqrt{\tau_L^2 + 18\sqrt{2}(\rho_0 c^2)^{-1} C_{sm}^2 \Delta x \bar{S}} - \tau_L \right)$$

Initial and boundary conditions

• Initial conditions are constructed as $f_{\alpha}^{eq}(\rho(t=0), \mathbf{u}(t=0))$



Initial and boundary conditions

• Initial conditions are constructed as $f^{eq}_{\alpha}(\rho(t=0), \mathbf{u}(t=0))$



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Initial and boundary conditions

• Initial conditions are constructed as $f^{eq}_{\alpha}(\rho(t=0), \mathbf{u}(t=0))$



- Outlet basically copies all distributions (zero gradient)
- Inlet and pressure boundary conditions use f^{eq}_α
- Embedded boundary conditions use ghost cell construction as before, then use $f^{eq}_{\alpha}(\rho', \mathbf{u}')$ to construct distributions in embedded ghost cells

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Adaptive mesh refinement for LBM				

1. Complete update on coarse grid: $f_{\alpha}^{C,n+1} := CT(f_{\alpha}^{C,n})$



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Adaptive mesh refinement for LBM				
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- 1. Complete update on coarse grid: $f_{\alpha}^{C,n+1} := CT(f_{\alpha}^{C,n})$
- 2. Interpolate $f_{\alpha,in}^{C,n}$ onto $f_{\alpha,in}^{f,n}$ to fill fine halos. Set physical boundary conditions.



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$$f^{f,n}_{\alpha,ir}$$

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Adaptive mesh refinement for LBM				

- 1. Complete update on coarse grid: $f_{\alpha}^{\mathcal{C},n+1} := \mathcal{CT}(f_{\alpha}^{\mathcal{C},n})$
- 2. Interpolate $f_{\alpha,in}^{C,n}$ onto $f_{\alpha,in}^{f,n}$ to fill fine halos. Set physical boundary conditions.
- 3. $\tilde{f}^{f,n}_{\alpha} := \mathcal{T}(f^{f,n}_{\alpha})$ on whole fine mesh. $f^{f,n+1/2}_{\alpha} := \mathcal{C}(\tilde{f}^{f,n}_{\alpha})$ in interior.



$$\tilde{f}^{f,n}_{lpha,in}$$

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$$\tilde{\mathbf{f}}^{f,n+1/2}_{\alpha,\mathrm{in}}$$

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- 2. Interpolate $f_{\alpha,in}^{C,n}$ onto $f_{\alpha,in}^{f,n}$ to fill fine halos. Set physical boundary conditions.
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 $\tilde{f}^{f,n+1/2}_{lpha,in}$

 $f_{\alpha,out}^{f,n}$

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- 2. Interpolate $f_{\alpha,in}^{C,n}$ onto $f_{\alpha,in}^{f,n}$ to fill fine halos. Set physical boundary conditions.
- 3. $\tilde{f}_{\alpha}^{f,n} := \mathcal{T}(f_{\alpha}^{f,n})$ on whole fine mesh. $f_{\alpha}^{f,n+1/2} := \mathcal{C}(\tilde{f}_{\alpha}^{f,n})$ in interior.
- 4. $\tilde{f}^{f,n+1/2}_{\alpha} := \mathcal{T}(f^{f,n+1/2}_{\alpha})$ on whole fine mesh. $f^{f,n+1}_{\alpha} := \mathcal{C}(\tilde{f}^{f,n+1/2}_{\alpha})$ in interior.



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 $\tilde{f}^{f,n}_{\alpha out}$

Adaptive lattice Boltzmann method	Aerodynamics cases	Fluid-structure coupling	Wind turbine wake aerodynamics	References
000000000				
Adaptive mesh refinement for LBM				

- 1. Complete update on coarse grid: $f_{\alpha}^{C,n+1} := CT(f_{\alpha}^{C,n})$
- 2. Interpolate $f_{\alpha,in}^{C,n}$ onto $f_{\alpha,in}^{f,n}$ to fill fine halos. Set physical boundary conditions.
- 3. $\tilde{f}_{\alpha}^{f,n} := \mathcal{T}(f_{\alpha}^{f,n})$ on whole fine mesh. $f_{\alpha}^{f,n+1/2} := \mathcal{C}(\tilde{f}_{\alpha}^{f,n})$ in interior.
- 4. $\tilde{f}^{f,n+1/2}_{\alpha} := \mathcal{T}(f^{f,n+1/2}_{\alpha})$ on whole fine mesh. $f^{f,n+1}_{\alpha} := \mathcal{C}(\tilde{f}^{f,n+1/2}_{\alpha})$ in interior.



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 $\tilde{f}^{f,n+1/2}_{lpha,out}$
Adaptive lattice Boltzmann method	Aerodynamics cases	Fluid-structure coupling	Wind turbine wake aerodynamics	References
000000000				
Adaptive mesh refinement for LBM				

- 1. Complete update on coarse grid: $f_{\alpha}^{C,n+1} := CT(f_{\alpha}^{C,n})$
- 2. Interpolate $f_{\alpha,in}^{C,n}$ onto $f_{\alpha,in}^{f,n}$ to fill fine halos. Set physical boundary conditions.
- 3. $\tilde{f}^{f,n}_{\alpha} := \mathcal{T}(f^{f,n}_{\alpha})$ on whole fine mesh. $f^{f,n+1/2}_{\alpha} := \mathcal{C}(\tilde{f}^{f,n}_{\alpha})$ in interior.
- 4. $\tilde{f}^{f,n+1/2}_{\alpha} := \mathcal{T}(f^{f,n+1/2}_{\alpha})$ on whole fine mesh. $f^{f,n+1}_{\alpha} := \mathcal{C}(\tilde{f}^{f,n+1/2}_{\alpha})$ in interior.

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$$\tilde{\textit{f}}_{\alpha,out}^{f,n+1/2}, \tilde{\textit{f}}_{\alpha,in}^{f,n+1/2}$$

Adaptive lattice Boltzmann method	Aerodynamics cases	Fluid-structure coupling	Wind turbine wake aerodynamics	References
000000000				
Adaptive mesh refinement for LBM				

- 1. Complete update on coarse grid: $f_{\alpha}^{C,n+1} := CT(f_{\alpha}^{C,n})$
- 2. Interpolate $f_{\alpha,in}^{C,n}$ onto $f_{\alpha,in}^{f,n}$ to fill fine halos. Set physical boundary conditions.
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- 4. $\tilde{f}^{f,n+1/2}_{\alpha} := \mathcal{T}(f^{f,n+1/2}_{\alpha})$ on whole fine mesh. $f^{f,n+1}_{\alpha} := \mathcal{C}(\tilde{f}^{f,n+1/2}_{\alpha})$ in interior.



5. Average $\tilde{f}_{\alpha,out}^{f,n+1/2}$ (inner halo layer), $\tilde{f}_{\alpha,out}^{f,n}$ (outer halo layer) to obtain $\tilde{f}_{\alpha,out}^{C,n}$.

Adaptive lattice Boltzmann method	Aerodynamics cases	Fluid-structure coupling	Wind turbine wake aerodynamics	References
000000000				
Adaptive mesh refinement for LBM				

- 1. Complete update on coarse grid: $f_{\alpha}^{C,n+1} := CT(f_{\alpha}^{C,n})$
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- 4. $\tilde{f}^{f,n+1/2}_{\alpha} := \mathcal{T}(f^{f,n+1/2}_{\alpha})$ on whole fine mesh. $f^{f,n+1}_{\alpha} := \mathcal{C}(\tilde{f}^{f,n+1/2}_{\alpha})$ in interior.



5. Average $\tilde{f}_{\alpha,out}^{f,n+1/2}$ (inner halo layer), $\tilde{f}_{\alpha,out}^{f,n}$ (outer halo layer) to obtain $\tilde{f}_{\alpha,out}^{C,n}$.

Adaptive lattice Boltzmann method	Aerodynamics cases	Fluid-structure coupling	Wind turbine wake aerodynamics	References
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Adaptive mesh refinement for LBM				

- 1. Complete update on coarse grid: $f_{\alpha}^{C,n+1} := CT(f_{\alpha}^{C,n})$
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- 4. $\tilde{f}^{f,n+1/2}_{\alpha} := \mathcal{T}(f^{f,n+1/2}_{\alpha})$ on whole fine mesh. $f^{f,n+1}_{\alpha} := \mathcal{C}(\tilde{f}^{f,n+1/2}_{\alpha})$ in interior.



- 5. Average $\tilde{f}_{\alpha,out}^{f,n+1/2}$ (inner halo layer), $\tilde{f}_{\alpha,out}^{f,n}$ (outer halo layer) to obtain $\tilde{f}_{\alpha,out}^{C,n}$.
- 6. Revert transport into halos: $\bar{f}_{\alpha,out}^{C,n} := \mathcal{T}^{-1}(\tilde{f}_{\alpha,out}^{C,n})$

Adaptive lattice Boltzmann method	Aerodynamics cases	Fluid-structure coupling	Wind turbine wake aerodynamics	References
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Adaptive mesh refinement for LBM				

- 1. Complete update on coarse grid: $f_{\alpha}^{C,n+1} := CT(f_{\alpha}^{C,n})$
- 2. Interpolate $f_{\alpha,in}^{C,n}$ onto $f_{\alpha,in}^{f,n}$ to fill fine halos. Set physical boundary conditions.
- 3. $\tilde{f}_{\alpha}^{f,n} := \mathcal{T}(f_{\alpha}^{f,n})$ on whole fine mesh. $f_{\alpha}^{f,n+1/2} := \mathcal{C}(\tilde{f}_{\alpha}^{f,n})$ in interior.
- 4. $\tilde{f}^{f,n+1/2}_{\alpha} := \mathcal{T}(f^{f,n+1/2}_{\alpha})$ on whole fine mesh. $f^{f,n+1}_{\alpha} := \mathcal{C}(\tilde{f}^{f,n+1/2}_{\alpha})$ in interior.



- 5. Average $\tilde{f}_{\alpha,out}^{f,n+1/2}$ (inner halo layer), $\tilde{f}_{\alpha,out}^{f,n}$ (outer halo layer) to obtain $\tilde{f}_{\alpha,out}^{C,n}$.
- 6. Revert transport into halos: $\bar{f}_{\alpha,out}^{C,n} := \mathcal{T}^{-1}(\tilde{f}_{\alpha,out}^{C,n})$
- 7. Parallel synchronization of $f_{\alpha}^{C,n}, \bar{f}_{\alpha,out}^{C,n}$

Adaptive lattice Boltzmann method	Aerodynamics cases	Fluid-structure coupling	Wind turbine wake aerodynamics	References
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Adaptive mesh refinement for LBM				

- 1. Complete update on coarse grid: $f_{\alpha}^{C,n+1} := CT(f_{\alpha}^{C,n})$
- 2. Interpolate $f_{\alpha,in}^{C,n}$ onto $f_{\alpha,in}^{f,n}$ to fill fine halos. Set physical boundary conditions.
- 3. $\tilde{f}_{\alpha}^{f,n} := \mathcal{T}(f_{\alpha}^{f,n})$ on whole fine mesh. $f_{\alpha}^{f,n+1/2} := \mathcal{C}(\tilde{f}_{\alpha}^{f,n})$ in interior.
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- 5. Average $\tilde{f}_{\alpha,out}^{f,n+1/2}$ (inner halo layer), $\tilde{f}_{\alpha,out}^{f,n}$ (outer halo layer) to obtain $\tilde{f}_{\alpha,out}^{C,n}$.
- 6. Revert transport into halos: $\bar{f}_{\alpha,out}^{C,n} := \mathcal{T}^{-1}(\tilde{f}_{\alpha,out}^{C,n})$
- 7. Parallel synchronization of $f_{\alpha}^{C,n}, \overline{f}_{\alpha,out}^{C,n}$
- 8. Cell-wise update where correction is needed: $f_{\alpha}^{C,n+1} := CT(f_{\alpha}^{C,n}, \bar{f}_{\alpha,out}^{C,n})$

Adaptive lattice Boltzmann method	Aerodynamics cases	Fluid-structure coupling	Wind turbine wake aerodynamics	References
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Adaptive mesh refinement for LBM				

- 1. Complete update on coarse grid: $f_{\alpha}^{C,n+1} := CT(f_{\alpha}^{C,n})$
- 2. Interpolate $f_{\alpha,in}^{C,n}$ onto $f_{\alpha,in}^{f,n}$ to fill fine halos. Set physical boundary conditions.
- 3. $\tilde{f}_{\alpha}^{f,n} := \mathcal{T}(f_{\alpha}^{f,n})$ on whole fine mesh. $f_{\alpha}^{f,n+1/2} := \mathcal{C}(\tilde{f}_{\alpha}^{f,n})$ in interior.
- 4. $\tilde{f}^{f,n+1/2}_{\alpha} := \mathcal{T}(f^{f,n+1/2}_{\alpha})$ on whole fine mesh. $f^{f,n+1}_{\alpha} := \mathcal{C}(\tilde{f}^{f,n+1/2}_{\alpha})$ in interior.



- 5. Average $\tilde{f}_{\alpha,out}^{f,n+1/2}$ (inner halo layer), $\tilde{f}_{\alpha,out}^{f,n}$ (outer halo layer) to obtain $\tilde{f}_{\alpha,out}^{C,n}$.
- 6. Revert transport into halos: $\overline{f}_{\alpha,out}^{C,n} := \mathcal{T}^{-1}(\widetilde{f}_{\alpha,out}^{C,n})$
- 7. Parallel synchronization of $f_{\alpha}^{C,n}, \bar{f}_{\alpha,out}^{C,n}$
- 8. Cell-wise update where correction is needed: $f_{\alpha}^{C,n+1} := CT(f_{\alpha}^{C,n}, \bar{f}_{\alpha,out}^{C,n})$

Algorithm equivalent to [Chen et al., 2006].

Adaptive lattice Boltzmann method	Aerodynamics cases	Fluid-structure coupling	Wind turbine wake aerodynamics	References
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Implementation				
Classes				

Directory <code>amroc/lbm</code> contains the lattice Boltzmann integrator that is in C++ throughout and also is built on the classes in <code>amroc/amr/Interfaces</code>.

Several SchemeType-classes are already provided: LBMD1Q3<DataType
 , LBMD2Q9<DataType >, LBMD3Q19<DataType >,
 LBMD2Q9Thermal<DataType >, LBMD3Q19Thermal<DataType > included a large number of boundary conditions.

code/amroc/doc/html/lbm/classLBMD1Q3.html code/amroc/doc/html/lbm/classLBMD2Q9.html

code/amroc/doc/html/lbm/classLBMD3Q19Thermal.html

Using function within LBMD?D?, the special coarse-fine correction is implemented in LBMFixup<LBMType, FixupType, dim>

code/amroc/doc/html/lbm/classLBMFixup.html

LBMIntegrator<LBMType, dim >, LBMGFMBoundary<LBMType, dim >, etc. interface to the generic classes in amroc/amr/Interfaces

code/amroc/doc/html/amr/classSchemeGFMBoundary.html

Problem.h: Specific simulation is defined in Problem.h only. Predefined classes specified in LBMStdProblem.h, LBMStdGFMProblem.h and LBMProblem.h.

code/amroc/doc/html/lbm/LBMProblem_8h_source.html code/amroc/doc/html/lbm/LBMStdProblem_8h.html

code/amroc/doc/html/lbm/LBMStdGFMProblem_8h.html

Flow over 2D cylinder, $d = 2 \,\mathrm{cm}$

- Air with $\nu = 1.61 \cdot 10^{-5} \text{ m}^2/\text{s},$ $\rho = 1.205 \text{ kg/m}^3$
- ▶ Domain size [-8d, 24d] × [-8d, 8d]
- Dynamic refinement based on velocity. Last level to refine structure further.
- Inflow from left. Characteristic boundary conditions [Schlaffer, 2013] elsewhere.



- Base lattice 320×160 , 3 additional levels with factors $r_l = 2, 4, 4$.
- Resolution: \sim 320 points in diameter d
- Computation of C_D on 400 equidistant points along circle and averaged over time. Comparison above with [Henderson, 1995].

Adaptive lattice Boltzmann method	Aerodynamics cases	Fluid-structure coupling	Wind turbine wake aerodynamics	Referenc
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Verification				

Flow over cylinder in 2d - Re = 300, $u = 0.2415 \,\text{m/s}$

Isolines on refinement and distribution to processors



Mesh adaptation with LBM:

- 1. Level-wise evaluation of $\omega_L^l = \frac{c_s^2}{\nu + \Delta t_l c_s^2/2}$
- 2. Exchange of distributions streaming across refinement interfaces

 $\verb|code/amroc/doc/html/apps/lbm_2applications_2Navier-Stokes_22d_2CylinderDrag_2src_2Problem_8h_source.html/apps/lbm_2applications_2Navier-Stokes_22d_2CylinderDrag_2src_2Problem_8h_source.html/apps/lbm_2applications_2Navier-Stokes_22d_2CylinderDrag_2src_2Problem_8h_source.html/apps/lbm_2applications_2Navier-Stokes_22d_2CylinderDrag_2src_2Problem_8h_source.html/apps/lbm_2applications_2Navier-Stokes_22d_2CylinderDrag_2src_2Problem_8h_source.html/apps/lbm_2applications_2Navier-Stokes_22d_2CylinderDrag_2src_2Problem_8h_source.html/apps/lbm_2applications_2Navier-Stokes_2Apps/lbm_2applications_2Navier-Stokes_2Apps/lbm_2applications_2Navier-Stokes_2Apps/lbm_2apps/$

Flow over cylinder in 2d - Re = 300, $u = 0.2415 \,\text{m/s}$

Isolines on refinement and distribution to processors



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Flow over cylinder in 2d - Re = 300, $u = 0.2415 \,\text{m/s}$

Isolines on refinement and distribution to processors



Mesh adaptation with LBM:

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Flow over cylinder in 2d - Re = 300, $u = 0.2415 \,\text{m/s}$

Isolines on refinement and distribution to processors



Mesh adaptation with LBM:

- 1. Level-wise evaluation of $\omega_L^l = \frac{c_s^2}{\nu + \Delta t_l c_s^2/2}$
- 2. Exchange of distributions streaming across refinement interfaces

Flow over cylinder in 2d - Re = 300, $u = 0.2415 \,\text{m/s}$

Isolines on refinement and distribution to processors



Mesh adaptation with LBM:

- 1. Level-wise evaluation of $\omega_L^{\prime} = \frac{c_s^2}{\nu + \Delta t_l c_s^2/2}$
- 2. Exchange of distributions streaming across refinement interfaces

Flow over cylinder in 2d - Re = 300, $u = 0.2415 \,\text{m/s}$

Isolines on refinement and distribution to processors



Mesh adaptation with LBM:

- 1. Level-wise evaluation of $\omega_L^l = \frac{c_s^2}{\nu + \Delta t_l c_s^2/2}$
- 2. Exchange of distributions streaming across refinement interfaces

Adaptive lattice Boltzmann method	Aerodynamics cases		Wind turbine wake aerodynamics	References
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Outline

Adaptive lattice Boltzmann method

Construction principles Adaptive mesh refinement for LBM Implementation Verification

Realistic aerodynamics computations Vehicle geometries

Fluid-structure coupling

Rigid body dynamics Validation simulations

Wind turbine wake aerodynamics

Mexico benchmark Simulation of wind turbine wakes Wake interaction prediction



- Inflow 40 m/s. LES model active. Characteristic boundary conditions.
- To t = 0.5 s (~ 4 characteristic lengths) with 31,416 time steps on finest level in ~ 37 h on 200 cores (7389 h CPU). Channel: 15 m × 5 m × 3.3 m

Adaptive lattice Boltzmann method	Aerodynamics cases	Fluid-structure coupling	Wind turbine wake aerodynamics	References
	00000			
Vehicle geometries				

Mesh adaptation



code/amroc/doc/html/apps/lbm_2applications_2Navier-Stokes_23d_2VehicleOnGround_2src_2Problem_8h_source.html

Adaptive lattice Boltzmann method	Aerodynamics cases	Fluid-structure coupling	Wind turbine wake aerodynamics	References
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Vehicle geometries				
Mesh adaptatio	n Used refinement blo	ocks and levels (indicate	ed by color)	

- SAMR base grid $600 \times 200 \times 132$ cells, $r_{1,2,3} = 2$ yielding finest resolution of $\Delta x = 3.125$ mm
- Adaptation based on level set and scaled gradient of magnitude of vorticity vector
- > 236M cells vs. 8.1 billion (uniform) at $t = 0.4075 \,\mathrm{s}$

Refinement at $t = 0.4075 \,\mathrm{s}$

Level	Grids	Cells
0	11,605	15,840,000
1	11,513	23,646,984
2	31,382	144,447,872
3	21,221	52,388,336

 $\verb|code/amroc/doc/html/apps/lbm_2applications_2Navier-Stokes_23d_2VehicleOnGround_2src_2Problem_8h_source.html/apps/lbm_2applications_2Navier-Stokes_23d_2VehicleOnGround_2src_2Problem_8h_source.html/apps/lbm_2applications_2Navier-Stokes_23d_2VehicleOnGround_2src_2Problem_8h_source.html/apps/lbm_2applications_2Navier-Stokes_23d_2VehicleOnGround_2src_2Problem_8h_source.html/apps/lbm_2applications_2Navier-Stokes_23d_2VehicleOnGround_2src_2Problem_8h_source.html/apps/lbm_2applications_2Navier-Stokes_23d_2VehicleOnGround_2src_2Problem_8h_source.html/apps/lbm_2applications_2Navier-Stokes_23d_2VehicleOnGround_2src_2Problem_8h_source.html/apps/lbm_2applications_2Navier-Stokes_23d_2VehicleOnGround_2src_2Problem_8h_source.html/apps/lbm_2applications_2Navier-Stokes_2Apps/lbm_2applications_2Navier-Stokes_2Apps/lbm_2apps/lbm_2$

Adaptive lattice Boltzmann method 00000000 Vehicle geometries Aerodynamics cases

Fluid-structure coupling

Vind turbine wake aerodynamic: 00000000000000000 References 00

Next Generation Train (NGT)

- 1:25 train model of 74,670 triangles
- Wind tunnel: air at room temperature with 33.48 m/s, $\mathrm{Re} = 250,000$, yaw angle 30°
- Comparison between LBM (fluid air) and incompressible OpenFOAM solvers





Adaptive lattice Boltzmann method 00000000 Vehicle geometries Aerodynamics cases

Fluid-structure coupling 0000 Vind turbine wake aerodynamics

References 00

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Averaged vorticity LBM-LES Averaged vorticity OpenFOAM-LES code/amroc/doc/html/apps/lbm_2applications_2Navier-Stokes_23d_2NGT2_2src_2Problem_8h_source.html Aerodynamics cases

Fluid-structure coupling 0000 Wind turbine wake aerodynamic: 0000000000000000 References 00

NGT model

- LBM-AMR computation with 5 additional levels, factor 2 refinement (uniform: 120.4e9 cells)
- Dynamic AMR until $T_c = 34$, then static for $\sim 12T_C$ to obtain average coefficients
- OpenFOAM simulations by M. Fragner (DLR)

Simulation	Mesh	CFX	CFY	CMX
Wind tunnel	-	-0.06	-5.28	-3.46
DDES	low	-0.40	-5.45	-3.61
Σ only	low	0.10	-0.04	-0.05
LES	high	-0.45	-6.07	-4.14
DDES	high	-0.43	-5.72	-3.77
LBM - p only	-	-0.30	-5.09	-3.46



Aerodynamics cases

Fluid-structure coupling 0000 Wind turbine wake aerodynamic: 00000000000000000 References 00

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LBM - p only	-	-0.30	-5.09	-3.46



	LBM	DDES(I)	LES	DDES(h)
Cells	147M	34.1M	219M	219M
y ⁺	43	3.2	1.7	1.7
x^{+}, z^{+}	43	313	140	140
Δx wake [mm]	0.936	3.0	1.5	1.5
Runtime $[T_C]$	34	35.7	10.3	9.2
Processors	200	80	280	280
CPU [h]	34,680	49,732	194,483	164,472
$T_C/\Delta t$	1790	1325	1695	1695
CPU [h]/ T_C /1M cells	5.61	39.75	86.4	81.36

Aerodynamics cases

Fluid-structure coupling 0000 Wind turbine wake aerodynamic: 00000000000000000 References 00

NGT model

- LBM-AMR computation with 5 additional levels, factor 2 refinement (uniform: 120.4e9 cells)
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Adaptive LBM code 16x faster than OpenFOAM with PISO algorithm on static mesh! Adaptive lattice Boltzmann method 00000000 Vehicle geometries Aerodynamics cases 0000● Fluid-structure coupling

Wind turbine wake aerodynamics

Strong scalability test (1:25 train)

- Computation is restarted from disk checkpoint at t = 0.526408 s from 96 core run.
- Time for initial re-partitioning removed from benchmark.
- 200 coarse level time steps computed.
- Regridding and re-partitioning every 2nd level-0 step.
- Computation starts with 51.8M cells (I3: 10.2M, I2: 15.3M, I1: 21.5M, I0= 4.8M) vs. 19.66 billion (uniform).
- Portions for parallel communication quite considerable (4 ghost cells still used).



		· • · · · / • • - -			-		
Cores	48	96	192	288	384	576	768
Time per step	132.43s	69.79s	37.47s	27.12s	21.91s	17.45s	15.15s
Par. Efficiency	100.0	94.88	88.36	81.40	75.56	63.24	54.63
LBM Update	5.91	5.61	5.38	4.92	4.50	3.73	3.19
Regridding	15.44	12.02	11.38	10.92	10.02	8.94	8.24
Partitioning	4.16	2.43	1.16	1.02	1.04	1.16	1.34
Interpolation	3.76	3.53	3.33	3.05	2.83	2.37	2.06
Sync Boundaries	54.71	59.35	59.73	56.95	54.54	52.01	51.19
Sync Fixup	9.10	10.41	12.25	16.62	20.77	26.17	28.87
Level set	0.78	0.93	1.21	1.37	1.45	1.48	1.47
Interp./Extrap.	3.87	3.81	3.76	3.49	3.26	2.75	2.39
Misc	2.27	1.91	1.79	1.67	1.58	1.38	1.25

Time in % spent in main operations

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Rigid body dynamics				
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Adaptive lattice Boltzmann method	Aerodynamics cases	Fluid-structure coupling	Wind turbine wake aerodynamics	References

Motion solver

Based on the Newton-Euler method solution of dynamics equation of kinetic chains [Tsai, 1999]

$$\begin{pmatrix} \mathbf{F} \\ \boldsymbol{\tau}_{\mathrm{P}} \end{pmatrix} = \begin{pmatrix} m\mathbf{1} & -m[\mathbf{c}]^{\times} \\ m[\mathbf{c}]^{\times}\mathbf{I}_{\mathrm{cm}} & -m[\mathbf{c}]^{\times}[\mathbf{c}]^{\times} \end{pmatrix} \begin{pmatrix} \mathbf{a}_{\mathrm{P}} \\ \boldsymbol{\alpha} \end{pmatrix} + \begin{pmatrix} m[\boldsymbol{\omega}]^{\times}[\boldsymbol{\omega}]^{\times}\mathbf{c} \\ [\boldsymbol{\omega}]^{\times}(\mathbf{I}_{\mathrm{cm}} - m[\mathbf{c}]^{\times}[\mathbf{c}]^{\times}) \boldsymbol{\omega} \end{pmatrix}.$$

$$\begin{split} m &= \text{mass of the body, } 1 = \text{the } 4 \times 4 \text{ homogeneous identity matrix,} \\ \mathbf{a}_p &= \text{acceleration of link frame with origin at } \mathbf{p} \text{ in the preceding link's frame,} \\ \mathbf{I}_{\rm cm} &= \text{moment of inertia about the center of mass,} \\ \boldsymbol{\omega} &= \text{angular velocity of the body,} \\ \boldsymbol{\alpha} &= \text{angular acceleration of the body,} \end{split}$$

 ${\bf c}$ is the location of the body's center of mass,

and $[\mathbf{c}]^{\times}$, $[\boldsymbol{\omega}]^{\times}$ denote skew-symmetric cross product matrices.

Rigid body dynamics				
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Adaptive lattice Boltzmann method	Aerodynamics cases	Fluid-structure coupling	Wind turbine wake aerodynamics	References

Motion solver

Based on the Newton-Euler method solution of dynamics equation of kinetic chains [Tsai, 1999]

$$\begin{pmatrix} \mathbf{F} \\ \boldsymbol{\tau}_{\mathrm{P}} \end{pmatrix} = \begin{pmatrix} m\mathbf{1} & -m[\mathbf{c}]^{\times} \\ m[\mathbf{c}]^{\times}\mathbf{I}_{\mathrm{cm}} & -m[\mathbf{c}]^{\times}[\mathbf{c}]^{\times} \end{pmatrix} \begin{pmatrix} \mathbf{a}_{\mathrm{P}} \\ \boldsymbol{\alpha} \end{pmatrix} + \begin{pmatrix} m[\boldsymbol{\omega}]^{\times}[\boldsymbol{\omega}]^{\times}\mathbf{c} \\ [\boldsymbol{\omega}]^{\times}(\mathbf{I}_{\mathrm{cm}} - m[\mathbf{c}]^{\times}[\mathbf{c}]^{\times}) \boldsymbol{\omega} \end{pmatrix}.$$

m = mass of the body, 1 = the 4×4 homogeneous identity matrix, $\mathbf{a}_p =$ acceleration of link frame with origin at \mathbf{p} in the preceding link's frame, $\mathbf{l}_{cm} =$ moment of inertia about the center of mass, $\boldsymbol{\omega} =$ angular velocity of the body, $\boldsymbol{\alpha} =$ angular acceleration of the body, \mathbf{c} is the location of the body's center of mass,

and $[\mathbf{c}]^{ imes}$, $[oldsymbol{\omega}]^{ imes}$ denote skew-symmetric cross product matrices.

Here, we additionally define the total force and torque acting on a body,

 $\mathbf{F} = (\mathbf{F}_{\textit{FSI}} + \mathbf{F}_{\textit{prescribed}}) \cdot \boldsymbol{\mathcal{C}}_{\textit{xyz}}$ and

 $\tau = (\tau_{FSI} + \tau_{prescribed}) \cdot \mathcal{C}_{\alpha\beta\gamma}$ respectively.

Where C_{xyz} and $C_{\alpha\beta\gamma}$ are the translational and rotational constraints, respectively.

Two-segment hinged wing

Configuration by [Toomey and Eldredge, 2008]. Manufactured bodies in tank filled with water. Prescribed translation and rotation

$$X_t(t) = rac{A_0}{2} rac{G_t(ft)}{max \ Gt} C(ft), \quad lpha_1(t) = -eta rac{G_r(ft)}{max \ Gr}$$

with $G_r(t) = tanh[\sigma_r cos(2\pi t + \Phi)],$

$$G_t(t) = \int_t tanh[\sigma_t cos(2\pi t')]dt'$$



$A_0(cm)$	7.1
c (cm)	5.1
d (cm)	0.25
$\rho_b (\mathrm{kg/m^3})$	5080
f (Hz)	0.15

Two-segment hinged wing

Configuration by [Toomey and Eldredge, 2008]. Manufactured bodies in tank filled with water. Prescribed translation and rotation

$$X_t(t) = \frac{A_0}{2} \frac{G_t(ft)}{\max Gt} C(ft), \quad \alpha_1(t) = -\beta \frac{G_r(ft)}{\max Gr}$$

with $G_r(t) = tanh[\sigma_r cos(2\pi t + \Phi)],$

$$G_t(t) = \int_t tanh[\sigma_t cos(2\pi t')]dt'$$

- 7 cases constructed by varying σ_r , σ_t , Φ
- ► Rotational Reynolds number $\operatorname{Re}_{r} = 2\pi\beta\sigma_{r}fc^{2}/(\tanh(\sigma_{r})\nu)$ varied between 2200 and 7200 in experiments
- [Toomey and Eldredge, 2008] reference simulations with a viscous particle method are for $Re_r = \{100, 500\}$



$A_0(cm)$	7.1
c (cm)	5.1
d (cm)	0.25
$\rho_b (\mathrm{kg/m^3})$	5080
f (Hz)	0.15

Case 1 - $\sigma_r = \sigma_t = 0.628$, $\Phi = 0$, $\operatorname{Re}_r = 100$

- Quiescent water $\rho_f = 997 \text{ kg/m}^3$ $c_s = 1497 \text{ m/s}$
- No-slip boundaries in y, periodic in x-direction
- ▶ Base level: 100 × 100 for [-0.5, 0.5] × [-0.5, 0.5] domain
- 4 additional levels with factors 2,2,2,4
- Coupling to rigid body motion solver on 4th level

Right: computed vorticity field (enlarged)



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Validation simulations				

Quantitative comparison

- Evaluate normalized force $F_{x,y} = 2F_{x,y}^* / (\rho_f^2 c^3)$ and moment $M = 2M^* / (\rho_f f^2 c^4)$ over 3 periods
- [Wood and Deiterding, 2015] Used finest spatial resolution $\Delta x/c = 0.0122$ [Toomey and Eldredge, 2008]: $\Delta x/c = 0.013$ (Re_r = 100), $\Delta x/c = 0.0032$ (Re_r = 500)
- Temporal resolution ~ 113 and ~ 28 times finer

Relative difference in mean force and moment							
	$Re_r = 100$				$Re_r = 500$		
Case	\bar{F}_{x}	\bar{F}_{y}	Ā	\bar{F}_{x}	\bar{F}_{y}	Ā	
1	-2.59	3.33	-3.85	3.33	5.45	-3.75	
2	2.47	0.74	2.55	2.35	3.83	-4.29	
3	1.27	0.45	0.72	2.31	4.65	-3.43	
4	4.86	4.28	3.54	3.51	2.37	-2.32	
5	4.83	0.47	0.25	4.34	4.39	-2.67	
6	2.10	3.19	1.52	3.00	1.82	-3.96	
7	1.41	0.99	3.28	4.31	2.32	-3.07	

	Relativ	e differenc	e in peak	force and 1	noment	
	$Re_r = 100$			$Re_r = 500$		
Case	$ F_x _{\infty}$	$ F_y _{\infty}$	$ M _{\infty}$	$ F_x _{\infty}$	$ F_y _{\infty}$	$ M _{\infty}$
1	4.40	5.07	-3.66	4.40	3.98	-4.17
2	4.46	2.42	2.62	2.72	4.33	-2.34
3	4.20	3.20	4.80	3.32	2.68	-4.59
4	4.67	2.22	3.71	0.18	2.51	-2.85
5	3.57	3.37	1.26	4.09	4.97	-3.63
6	2.04	3.08	1.52	3.92	2.08	-4.44
7	2.20	1.91	2.26	3.29	3.79	-4.40

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Validation simulations				

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Mexico experimental turbine -0° inflow



- Setup and measurements by Energy Research Centre of the Netherlands (ECN) and the Technical University of Denmark (DTU) [Schepers and Boorsma, 2012]
- ▶ Inflow velocity $14.93 \,\mathrm{m/s}$ in wind tunnel of $9.5 \,\mathrm{m} \times 9.5 \,\mathrm{m}$ cross section.
- ▶ Rotor diameter D = 4.5w m. Prescribed motion with 424.5 rpm: tip speed 100 m/s, Re_r ≈ 75839 TSR 6.70
- Simulation with three additional levels with $r_l = \{2, 2, 4\}$. Resolution of rotor and tower $\Delta x = 1.6$ cm
- 149.5 h on 120 cores Intel-Xeon (17490 h CPU) for 10 s

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- Simulation with three additional levels with $r_l = \{2, 2, 4\}$. Resolution of rotor and tower $\Delta x = 1.6$ cm
- $\blacktriangleright\,$ 149.5 h on 120 cores Intel-Xeon (17490 h CPU) for 10 s
- Blade loads: F_x : Ref = 1516.76 N, cur. = 1632.71 N (7.6%)
- T_x: Ref = 284.60 Nm, cur. = 307.87 Nm (8.1%)

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Mexico benchmark				

Comparison along transects -0° inflow



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Mexico benchmark				

Comparison along transects -0° inflow


Mexico experimental turbine – 30° yaw



- ▶ Load collected as average during $t \in [5, 10]$ on blade 1 as it passes through $\theta = 0^{\circ}$ (pointing vertically upwards), 35 rotations
- Blade loads: F_x : Ref = 13.66 N, cur. = 14.8 N (8.3%)
- T_x: Ref = 7.72 Nm, cur. = 8.36 Nm (8.3%)

Level 0: 768,000 cells Level 1: 1,524,826 cells Level 2: 6,832,602 cells Level 3: 3,019,205 cells

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Mexico benchmark				

Comparison along transects – 30° yaw



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Mexico benchmark				

Comparison along transects – 30° yaw



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Mexico benchmark				

Normalized %-error along transects

	yaw	0	0	30)°
	transect	in	out	in	out
	u _x	6.416	7.663	5.742	6.410
Axial	\mathbf{u}_{y}	3.400	4.061	3.043	3.373
	u _z	3.073	3.678	2.752	3.068
		up	down	up	down
	u _x	6.556	7.325	7.093	6.655
Radial	\mathbf{u}_{y}	3.409	3.809	3.684	3.466
	u _z	3.242	3.659	3.511	3.294

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Comparison of normal and tangential forces on sections of blade 1 when $\theta_x = 0^\circ$ (pointing vertically upward) in aligned (left) and yaw 30°

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Simulation of wind turbine wakes

- $\blacktriangleright\,$ Geometry from realistic Vestas V27 turbine. Rotor diameter 27 $\rm m,\,tower$ height $\sim35\,\rm m.\,$ Ground considered.
- Prescribed motion of rotor with 15 rpm. Inflow velocity 7 m/s.
- Simulation domain $200 \text{ m} \times 100 \text{ m} \times 100 \text{ m}$.
- Base mesh 400 \times 200 \times 200 cells with refinement factors 2,2,4. Resolution of rotor and tower $\Delta x=3.125\,{\rm cm}.$
- 141,344 highest level iterations to $t_e = 30 \text{ s}$ computed.



Aerodynamics ca 00000 Fluid-structure coupling

Wind turbine wake aerodynamics

References 00

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Aerodynamics ca 00000 Fluid-structure coupling

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Adaptive lattice Boltzmann method Aerodynamics c 000000000 00000 Simulation of wind turbine wakes Fluid-structure coupling

Wind turbine wake aerodynamics

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Adaptive lattice Boltzmann method Aerodynamics 000000000 00000 Simulation of wind turbine wakes

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Adaptive lattice Boltzmann method	Aerodynamics cases	Fluid-structure coupling	Wind turbine wake aerodynamics	References
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Simulation of wind turbine wakes				

Adaptive refinement

Time=6.65651 sec

Dynamic evolution of refinement blocks (indicated by color). code/doc/html/capps/motion-amroc_2WindTurbine_Terrain_2src_2FluidProblem_8h_source.html, code/doc/html/capps/motion-amroc_2WindTurbine_Terrain_2src_2SolidProblem_8h_source.html, code/doc/html/capps/Terrain_2src_2Terrain_8h_source.html Adaptive lattice Boltzmann method Wake interaction prediction

Simulation of the SWIFT array

- \blacktriangleright Three Vestas V27 turbines. 225 kW power generation at wind speeds 14 to $25 \,\mathrm{m/s}$ (then cut-off)
- ► Prescribed motion of rotor with 33 and 43 rpm. Inflow velocity 8 and 25 m/s
- TSR: 5.84 and 2.43, $Re_r \approx 919,700$ and 1,208,000 ►
- ► Simulation domain $448 \text{ m} \times 240 \text{ m} \times 100 \text{ m}$
- Base mesh $448 \times 240 \times 100$ cells with ► refinement factors 2,2,4. Resolution of rotor and tower $\Delta x = 6.25 \,\mathrm{cm}$
- > 94,224 highest level iterations to $t_e = 40 \, \mathrm{s}$ computed, then statistics are gathered for 10s [Deiterding and Wood, 2015]







- On 288 cores Intel Xeon-Ivybride 10 s in 38.5 h (11,090 h CPU)
- Only levels 0 and 1 used for iso-surface visualization
- At *t_e* approximately 140M cells used vs. 44 billion (factor 315)
- Only levels 0 and 1 used for iso-surface visualization

Level	Grids	Cells
0	3,234	10,752,000
1	11,921	21,020,256
2	66,974	102,918,568
3	896	5,116,992

Adaptive lattice Boltzmann method	Aerodynamics cases	Fluid-structure coupling	Wind turbine wake aerodynamics	Reference
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Wake interaction prediction				

Vorticity generation - u = 25 m/s, 43 rpm



• Refinement of wake up to level 2 ($\Delta x = 25 \text{ cm}$).

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Wake interaction prediction				

Vorticity generation - u = 8 m/s, 33 rpm



- Refinement of wake up to level 2 ($\Delta x = 25 \text{ cm}$).
- Vortex break-up before 2nd turbine is reached.

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Wake interaction prediction				

Vorticity development - $u=8\,\mathrm{m/s}$, 33 rpm



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Refinement $u = 8 \,\mathrm{m/s}$, $33 \,\mathrm{rpm}$



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Wake interaction prediction



Adaptive lattice Boltzmann method Wake interaction prediction

Wind turbine wake aerodynamics

Refinement u = 8 m/s, 33 rpm


































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Wake interaction prediction				

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Mean point values

- Turbines located at (0,0,0), (135, 0, 0), (-5.65, 80.80, 0)
- Lines of 13 sensors with $\Delta y = 5 \,\mathrm{m}, z = 37 \,\mathrm{m}$ (approx. center of rotor)
- u and p measured over [40 s, 50 s] (1472 level-0 time steps) and averaged





Velocity deficits larger for higher TSR.

Adaptive lattice Boltzmann method	Aerodynamics cases	Fluid-structure coupling	Wind turbine wake aerodynamics	References
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Wake interaction prediction				

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- Velocity deficits larger for higher TSR.
- Velocity deficit before 2nd turbine more homogenous.

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Wake interaction prediction				

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