

Deciphering the Structure of Gaseous Detonations by Numerical Simulation

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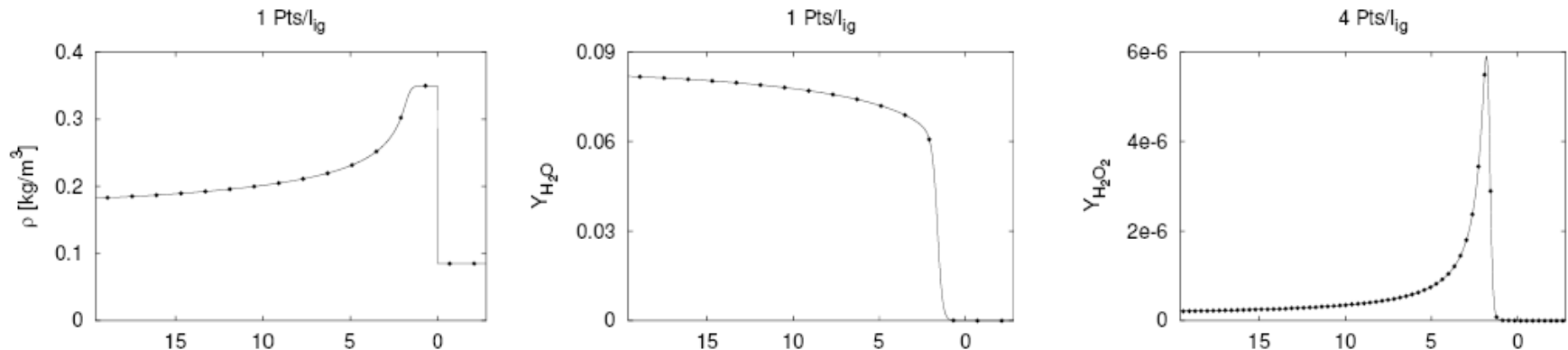
Work supported by DFG high priority research program “Analysis and Numerics of Conservation Laws”, grant Ba 840/3-3 and currently sponsored by the Office of Advanced Scientific Computing Research; U.S. Department of Energy (DOE) and was performed at the Oak Ridge National Laboratory, which is managed by UT-Battelle, LLC under Contract No. DE-AC05-00OR22725.

Outline of the talk

- Computational approach
 - *Introduction*
 - *Governing equations*
 - *Finite volume method*
 - *Upwind scheme, embedded boundaries*
 - *Dynamic mesh adaptation, parallelization*
 - *Shock-induced combustion examples*
- Mach reflection patterns in detonations
 - *Qualitative description*
 - *Shock polar analysis for triple points in detonations*
 - *Oblique shock relations for real gases*
 - *Reflection type transition criteria*
 - *Construction of a transition diagram for low-pressure $H_2:O_2:Ar$*
 - *Weak and strong structures in detonations in 2D pipe bends*
 - *Classification of observed Mach reflection structures*
- Short outlook: comparison 2D vs. 3D
- Conclusions

Difficulties in detonation simulations

1. Discontinuous solutions $\rightarrow$ high-resolution finite volume method with upwinding in all characteristic fields
2. Stiffness of reaction terms, $\Delta t_c \ll \Delta t \rightarrow$ Numerical decoupling of time operators with method of fractional steps and local time steps Δt_c
3. Extremely high spatial resolution in reaction zone necessary. Discretization of an exact ZND detonation:



minimal spatial resolution: 7 – 8 Pts/ $l_{ig} \rightarrow \Delta x \approx 0.2 - 0.175\text{mm}$

Uniform grids for typical geometries: $> 10^7$ Pts in 2D, $> 10^9$ Pts in 3D $\rightarrow$ Self-adaptive finite volume method (AMR)

4. Problem size even with AMR in 3D enormous $\rightarrow$ parallelization for massively parallel systems with distributed memory

Hydrodynamic equations

Euler equations for mixtures

$$\frac{\partial \rho}{\partial t} + \frac{\partial}{\partial x_k}(\rho u_k) = 0$$

$$\frac{\partial}{\partial t}(\rho u_i) + \frac{\partial}{\partial x_k}(\rho u_i u_k + \delta_{ik} p) = 0$$

$$\frac{\partial E}{\partial t} + \frac{\partial}{\partial x_k}(u_k(E + p)) = 0$$

$$\frac{\partial}{\partial t}(\rho Y_i) + \frac{\partial}{\partial x_k}(\rho Y_i u_k) = \dot{m}_i$$

Ideal gas law

$$p = \rho \mathcal{R} T \sum_{i=1}^N \frac{Y_i}{W_i}$$

Caloric equation

$$h = \sum_{i=1}^N h_i(T) Y_i \quad \text{with}$$

$$h_i(T) = h_i^0 + \int_{T^0}^T c_{pi}(T^*) dT^*$$

Implicit equation of state

$$\rho h - p - E + \frac{1}{2} \rho u_k u_k = 0$$

Chemical kinetics with Arrhenius law

$$\dot{m}_i = W_i \sum_{j=1}^M (\nu_{ji}^r - \nu_{ji}^f) [k_j^f \prod_{n=1}^N \left(\frac{\rho_n}{W_n}\right)^{\nu_{jn}^f} - k_j^r \prod_{n=1}^N \left(\frac{\rho_n}{W_n}\right)^{\nu_{jn}^r}]$$

Finite volume scheme

- Method of fractional steps

$$\mathcal{H}(\Delta t) : \quad \partial_t \mathbf{q} + \nabla \cdot \mathbf{f}(\mathbf{q}) = 0, \quad \text{IC: } \mathbf{Q}(t_m) \xrightarrow{\Delta t} \tilde{\mathbf{Q}}$$

$$\mathcal{S}(\Delta t) : \quad \partial_t \mathbf{q} = \mathbf{s}(\mathbf{q}), \quad \text{IC: } \tilde{\mathbf{Q}} \xrightarrow{\Delta t} \mathbf{Q}(t_m + \Delta t)$$

$$\text{1st-order: } \mathbf{Q}(t_m + \Delta t) = \mathcal{S}(\Delta t) \mathcal{H}(\Delta t)(\mathbf{Q}(t_m))$$

$$\text{2nd-order: } \mathbf{Q}(t_m + \Delta t) = \mathcal{S}(\tfrac{1}{2}\Delta t) \mathcal{H}(\Delta t) \mathcal{S}(\tfrac{1}{2}\Delta t)(\mathbf{Q}(t_m))$$

- Hydrodynamics

- Extension to 2d and 3d via dimensional splitting
- Linearized Riemann-solver of Roe-type for thermally perfect gas-mixtures (Grossman, Cinella, J. Comput. Phys. 85, 1990)
- Positivity-preserving
 - Switching to HLL for unphysical ρ, p
 - Fix for Y_i (Larrouturou, J. Comput. Phys. 95, 1991).
- Entropy fix modification to avoid carbuncle phenomenon (Sanders, Morano, Druguet, J. Comput. Phys. 145, 1998)
- 2nd-order MUSCL reconstruction

- Evaluation of T with Newton iteration / bisection

RD, G. Bader, Analysis and Numerics for Conservation Laws, p. 69-91, Springer, 2005

Reaction mechanism

- 4th-order semi-implicit Rosenbrock-Wanner ODE solver with stepsize adjustment
- Production rates evaluated with automatically generated F77 function (4x faster Chemkin)
- Jacobian approximated
- c_{pi} , h_i tabulated
- All subsequent computations for hydrogen-oxygen mechanism with 34 elementary reaction, 9 species O_2 , H_2 , H_2O , H , O , OH , HO_2 , H_2O_2 , Ar

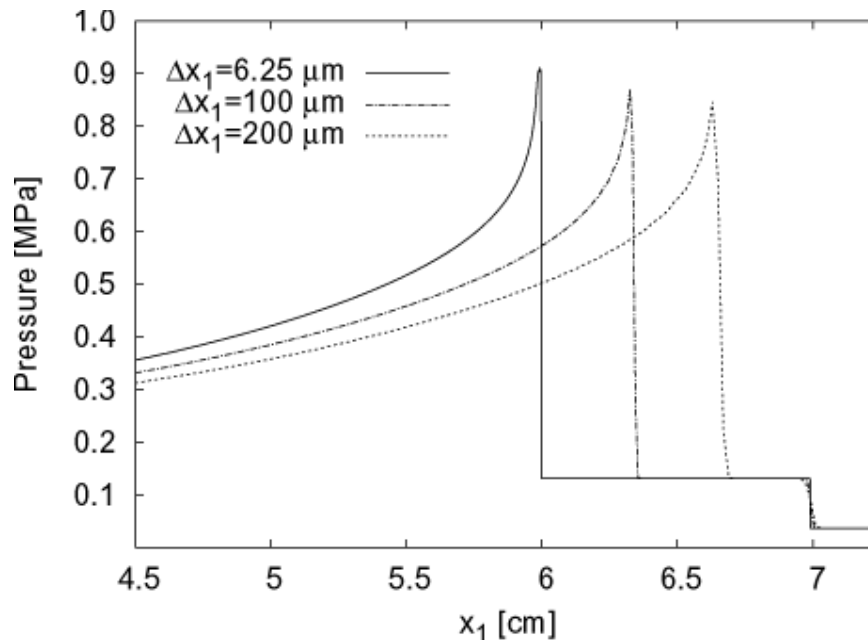
			A	β	E_{act}	
			[cm, mol, s]		[cal mol ⁻¹]	
1.	H + O ₂	→	O + OH	1.86×10^{14}	0.00	16790.
2.	O + OH	→	H + O ₂	1.48×10^{13}	0.00	680.
3.	H ₂ + O	→	H + OH	1.82×10^{10}	1.00	8900.
4.	H + OH	→	H ₂ + O	8.32×10^{09}	1.00	6950.
5.	H ₂ O + O	→	OH + OH	3.39×10^{13}	0.00	18350.
6.	OH + OH	→	H ₂ O + O	3.16×10^{12}	0.00	1100.
7.	H ₂ O + H	→	H ₂ + OH	9.55×10^{13}	0.00	20300.
8.	H ₂ + OH	→	H ₂ O + H	2.19×10^{13}	0.00	5150.
9.	H ₂ O ₂ + OH	→	H ₂ O + HO ₂	1.00×10^{13}	0.00	1800.
10.	H ₂ O + HO ₂	→	H ₂ O ₂ + OH	2.82×10^{13}	0.00	32790.
11.	HO ₂ + O	→	OH + O ₂	5.01×10^{13}	0.00	1000.
12.	OH + O ₂	→	HO ₂ + O	6.46×10^{13}	0.00	56160.
13.	HO ₂ + H	→	OH + OH	2.51×10^{14}	0.00	1900.
14.	OH + OH	→	HO ₂ + H	1.20×10^{13}	0.00	40100.
15.	HO ₂ + H	→	H ₂ + O ₂	2.51×10^{13}	0.00	700.
16.	H ₂ + O ₂	→	HO ₂ + H	5.50×10^{13}	0.00	57800.
17.	HO ₂ + OH	→	H ₂ O + O ₂	5.01×10^{13}	0.00	1000.
18.	H ₂ O + O ₂	→	HO ₂ + OH	6.31×10^{14}	0.00	73860.
19.	H ₂ O ₂ + O ₂	→	HO ₂ + HO ₂	3.98×10^{13}	0.00	42640.
20.	HO ₂ + HO ₂	→	H ₂ O ₂ + O ₂	1.00×10^{13}	0.00	1000.
21.	H ₂ O ₂ + H	→	HO ₂ + H ₂	1.70×10^{12}	0.00	3750.
22.	HO ₂ + H ₂	→	H ₂ O ₂ + H	7.24×10^{11}	0.00	18700.
23.	H ₂ O + M	→	H + OH + M	2.19×10^{16}	0.00	105000.
24.	H + OH + M	→	H ₂ O + M	1.41×10^{23}	-2.00	0.
25.	H + O ₂ + M	→	HO ₂ + M	1.66×10^{15}	0.00	-1000.
26.	HO ₂ + M	→	H + O ₂ + M	2.29×10^{15}	0.00	45900.
27.	H ₂ O ₂ + M	→	OH + OH + M	1.20×10^{17}	0.00	45500.
28.	OH + OH + M	→	H ₂ O ₂ + M	9.12×10^{14}	0.00	-5070.
29.	O + H + M	→	OH + M	1.00×10^{16}	0.00	0.
30.	OH + M	→	O + H + M	7.94×10^{19}	-1.00	103720.
31.	O ₂ + M	→	O + O + M	5.13×10^{15}	0.00	115000.
32.	O + O + M	→	O ₂ + M	4.68×10^{15}	-0.28	0.
33.	H ₂ + M	→	H + H + M	2.19×10^{14}	0.00	96000.
34.	H + H + M	→	H ₂ + M	3.02×10^{15}	0.00	0.

Third body efficiencies: $f(O_2) = 0.40$, $f(H_2O) = 6.50$

C. K. Westbrook. Chemical kinetics of hydrocarbon oxidation in gaseous detonations. *J. Combustion and Flame*, 46:191–210, 1982.

Detonation ignition in a shock tube

- Shock-induced detonation ignition of $\text{H}_2 : \text{O}_2 : \text{Ar}/2 : 1 : 7$ in a 1d shock tube closed at the left end, domain simulated: 12 cm
- Insufficient resolution leads to inaccurate results
- Reflected shock is captured by the FV scheme correctly at all resolutions, but the detonation is resolution-dependent



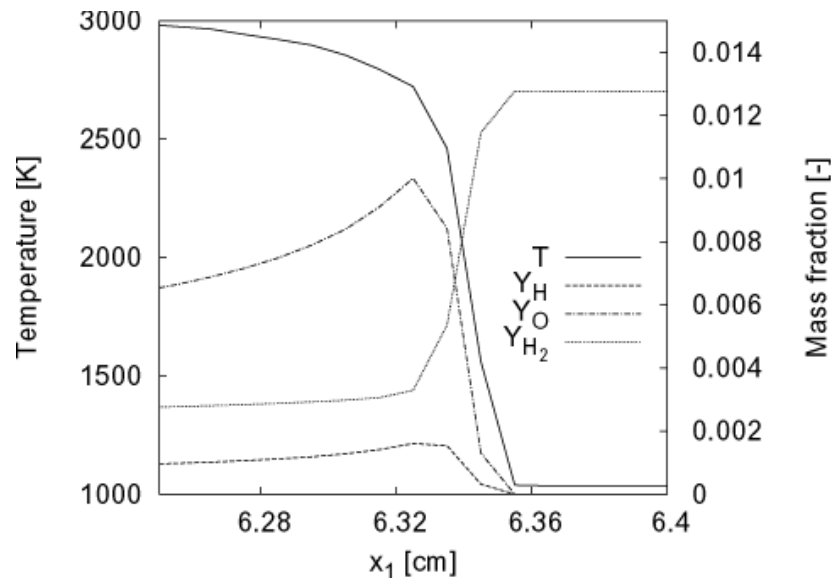
$\Delta x_1 = 100 \mu\text{m}$

$\Delta x_1 = 6.25 \mu\text{m}$

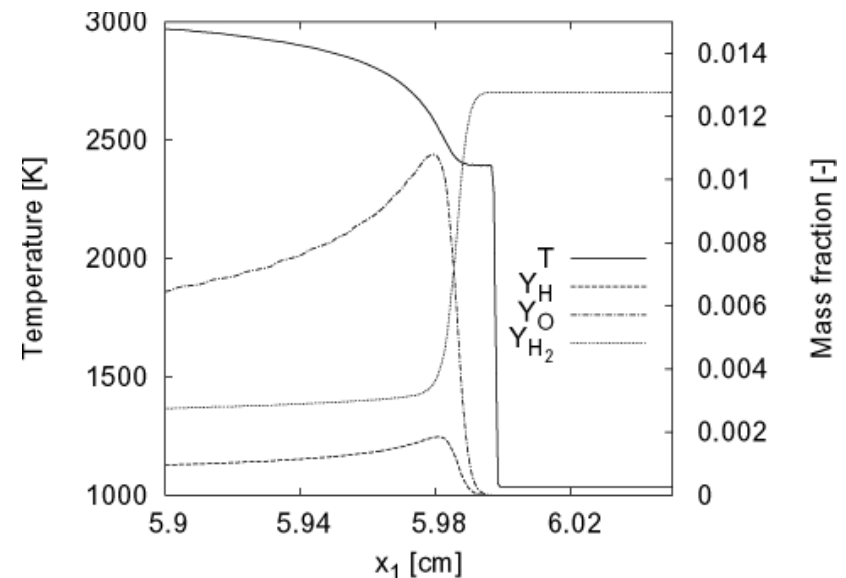
Analysis of configuration in Oran et al., J. Combust. Flame, 48:135–148, 1982

Detonation ignition in a shock tube

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Structured AMR for hyperbolic problems

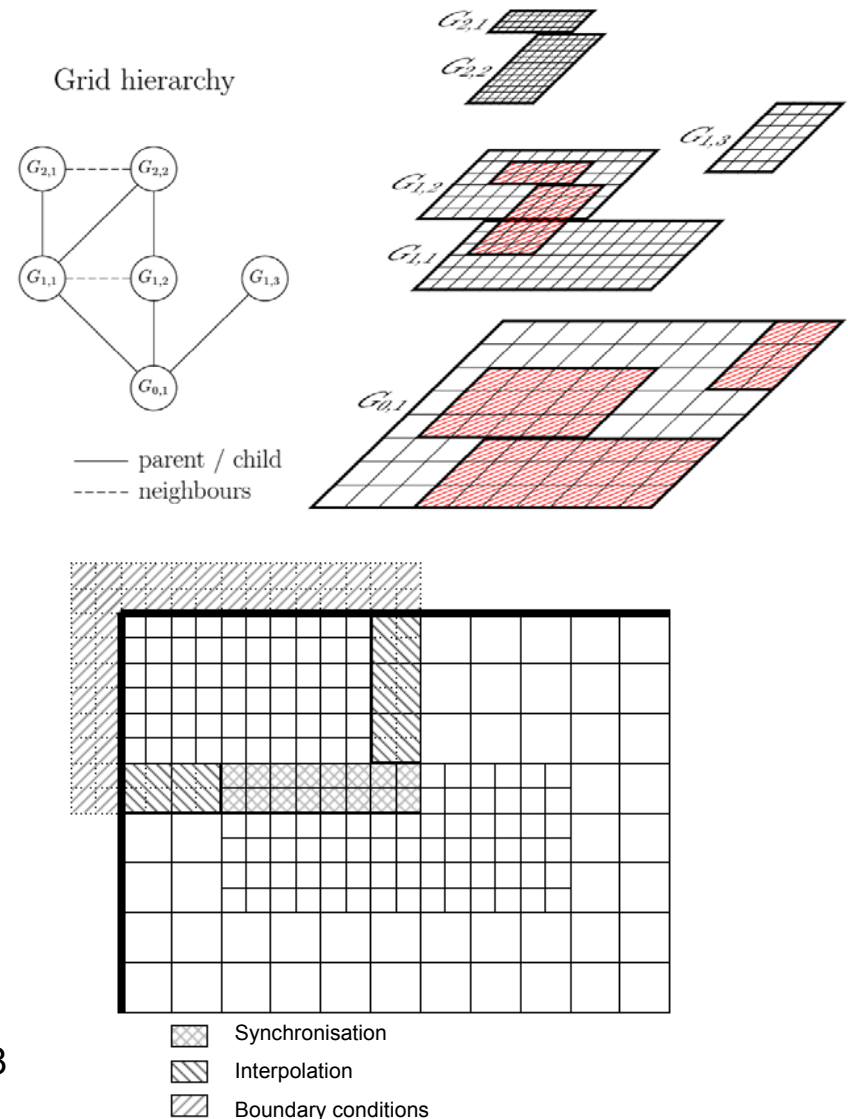
- Refined subgrids overlay coarser ones
- Computational decoupling of subgrids by using ghost cells
- Refinement in space *and* time
- Block-based data structures
- Cells without mark are refined
- Cluster-algorithm necessary
- Efficient cache-reuse / vectorization possible
- Discretization

$$Q_{jk}^{n+1} = Q_{jk}^n - \frac{\Delta t}{\Delta x_1} \left[F_{j+\frac{1}{2},k}^1 - F_{j-\frac{1}{2},k}^1 \right] - \frac{\Delta t}{\Delta x_2} \left[F_{j,k+\frac{1}{2}}^2 - F_{j,k-\frac{1}{2}}^2 \right]$$

is applied patch-wise

→ Inherently parallel approach

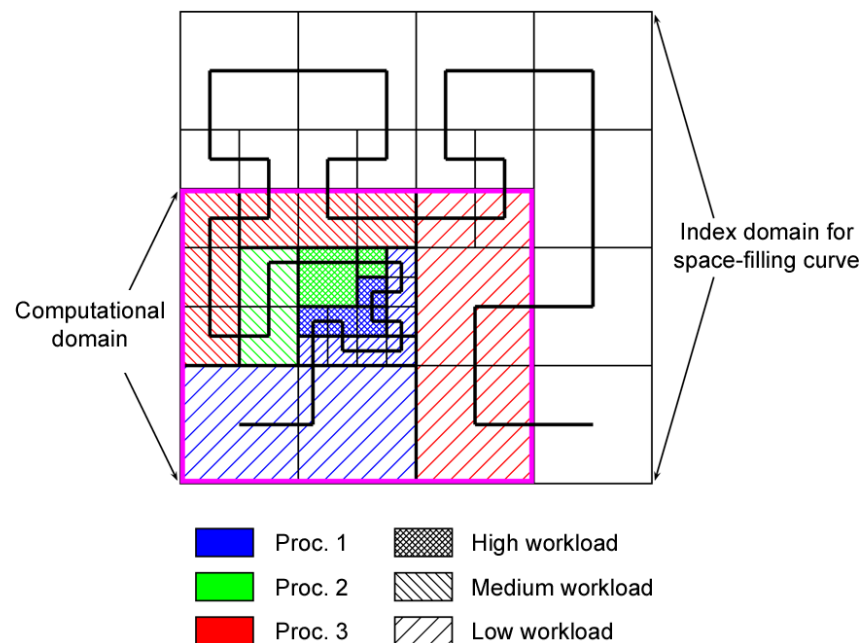
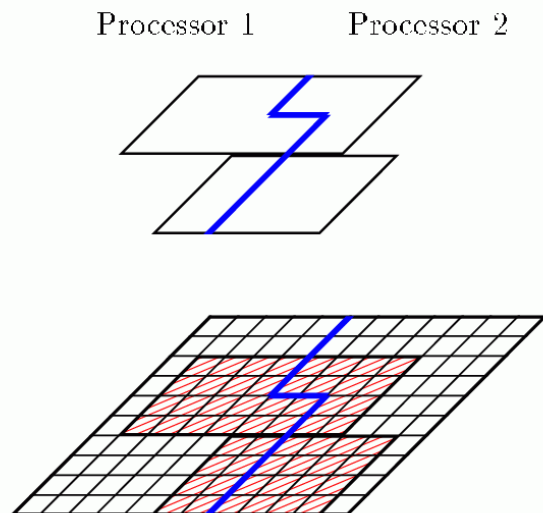
M. Berger and P. Colella, J. Comput. Phys. 82, 1988



Parallelization strategy

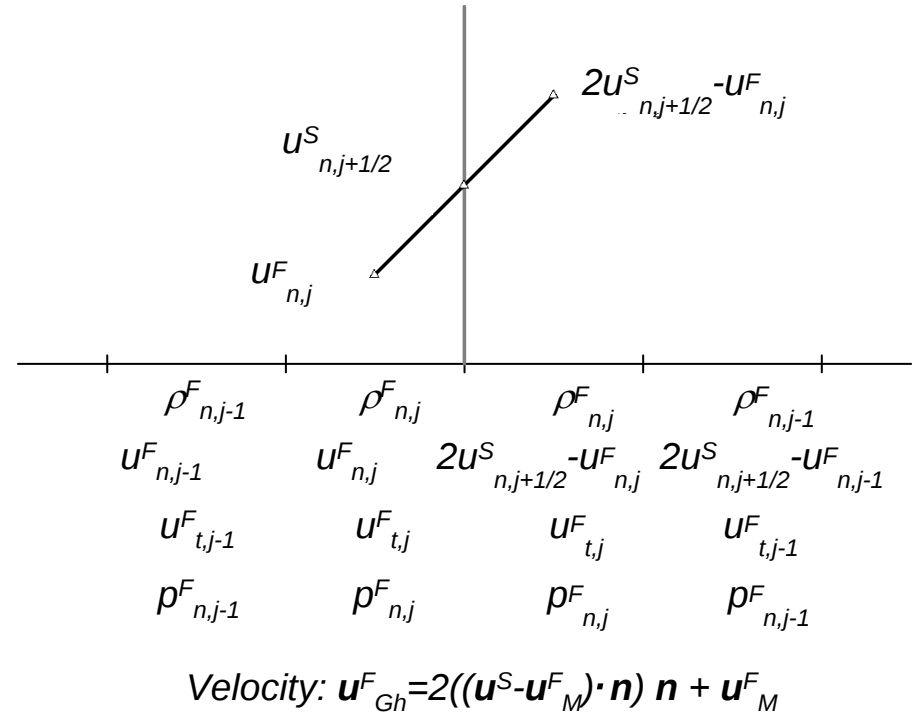
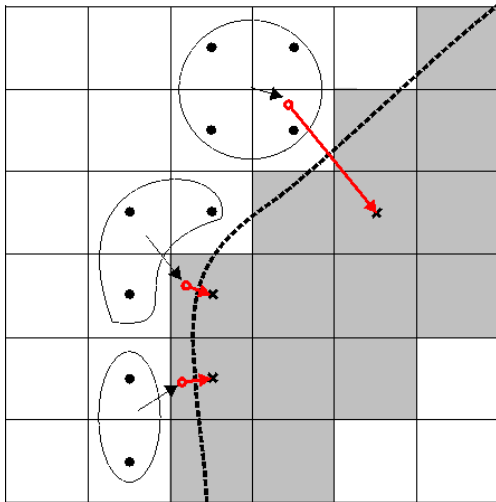
- Data of all levels resides on same node → Interpolation and averaging remain strictly local
- Only parallel operations to be considered:
 - *Parallel synchronization as part of ghost cell setting*
 - *Load-balanced repartitioning of data blocks as part of **Regrid(1)***
 - *Application of flux correction terms on coarse-grid cells*
- Partitioning at root level with generalized Hilbert space-filling curve defined in AMR index coordinate system by M. Parashar

$$\mathcal{W}(\Omega) = \sum_{l=0}^{l_{\max}} \left[\mathcal{N}_l(G_l \cap \Omega) \prod_{\nu=0}^l r_{\nu} \right]$$



Embedded boundary method

- Incorporate complex moving boundary/interfaces into a Cartesian solver (extension of work by R.Fedkiw and T.Aslam)
- Implicit boundary representation via distance function ϕ , normal $n = \nabla\phi / |\nabla\phi|$
- Treat an interface as a moving rigid wall
- Method diffuses boundary and is therefore not conservative
- Construction of values in embedded boundary cells by interpolation / extrapolation

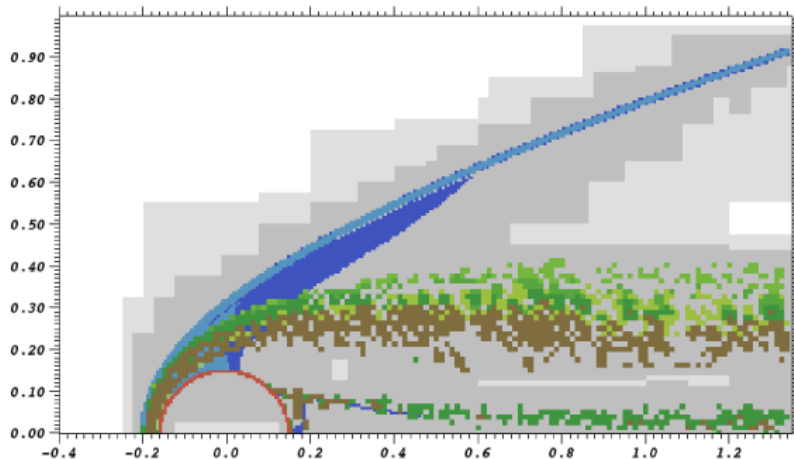


- Higher resolution at embedded boundary required than with first-order unstructured scheme
- Appropriate level-set-based refinement criteria are available to cure deficiencies

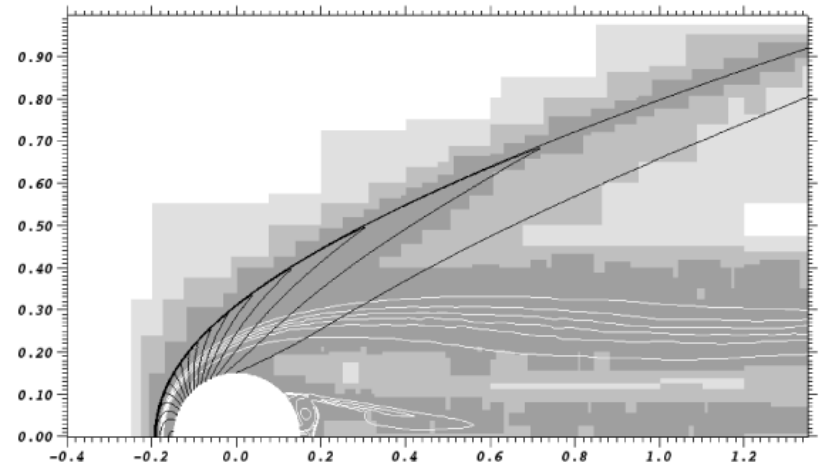
RD, Computers & Structures, 87: 769-783, 2009

Shock-induced combustion around a sphere

- Spherical projectile of radius 1.5 mm travels with constant velocity $v_f = 2170.6$ m/s through $\text{H}_2 : \text{O}_2 : \text{Ar}/2:1:7$ at 6.67 kPa and $T=298$ K
- Cylindrical symmetric simulation on AMR base mesh of 70x40 cells
- Comparison in quasi-steady state at $t=350 \mu\text{s}$



Refinement indicators on level 3

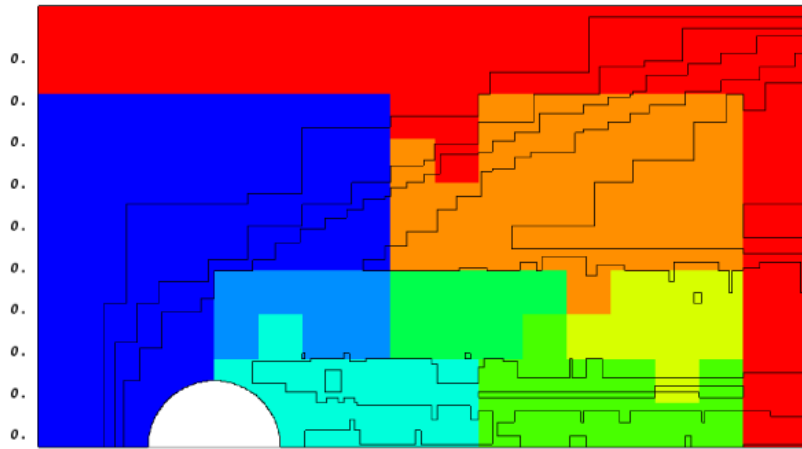


4-level with factors 2,2,4 (~ 19 Pts/ l_{ig})

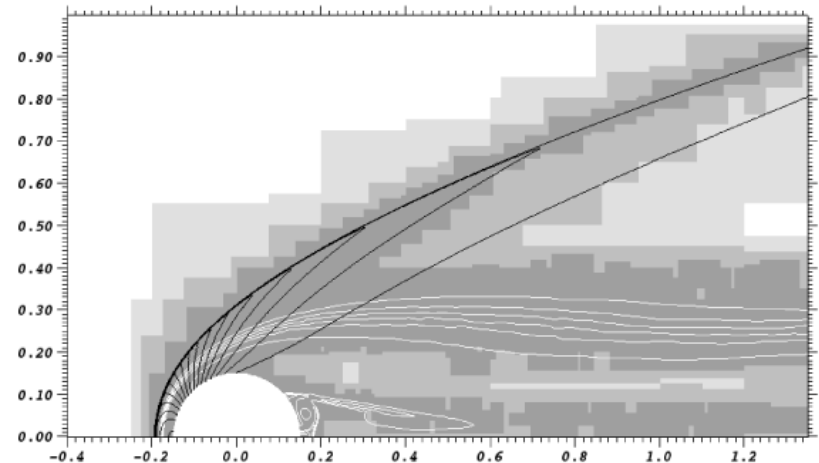
Setup from P. Hung, PhD thesis, GalCIT, 2003

Shock-induced combustion around a sphere

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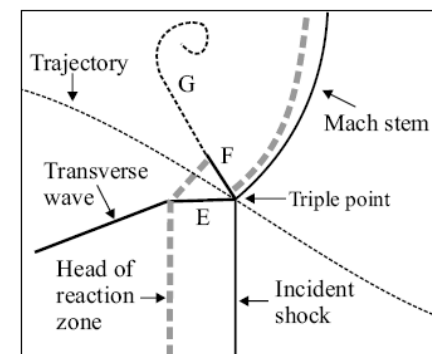
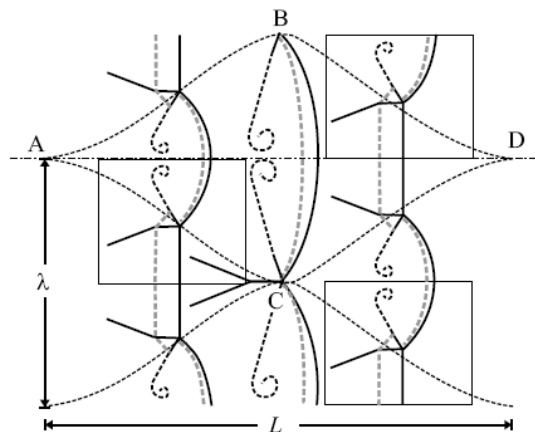
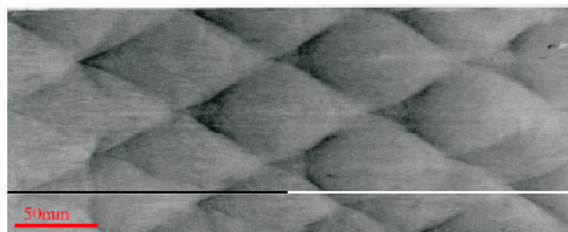
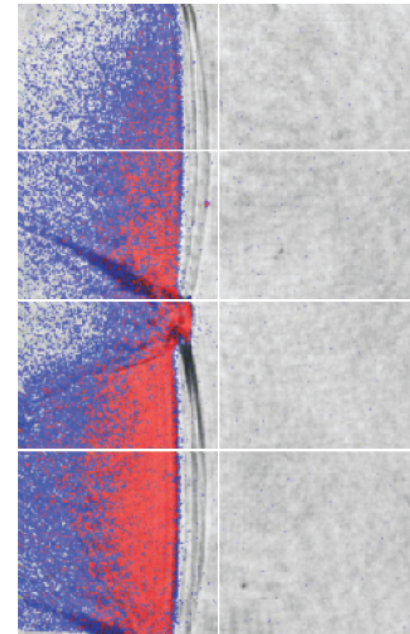
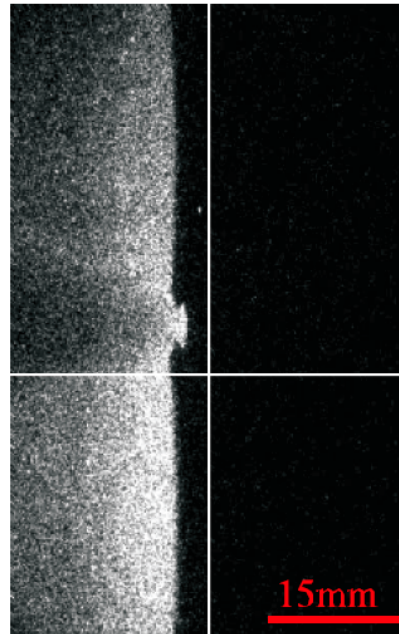
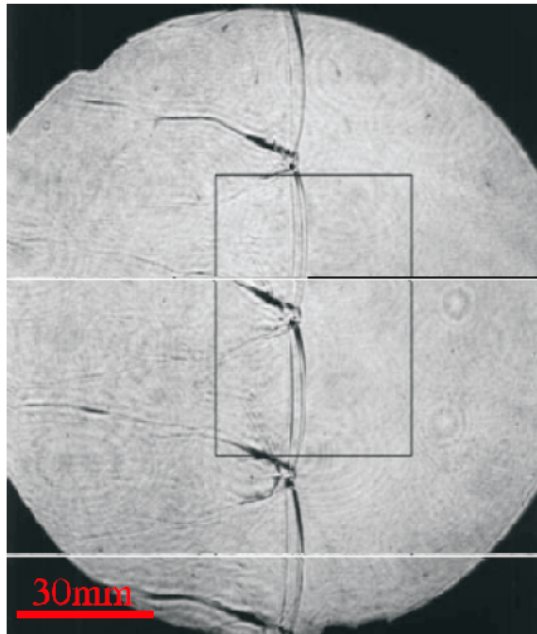
Distribution to 8 processors



4-level with factors 2,2,4 (~ 19 Pts/ l_{ig})

Setup from P. Hung, PhD thesis, GalCIT, 2003

Transverse detonation structure - Regular instability



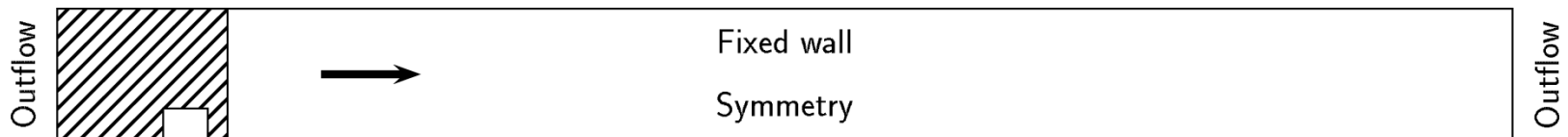
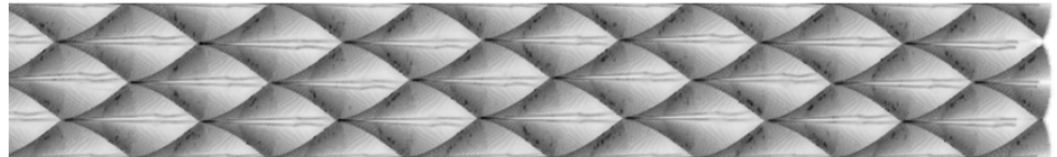
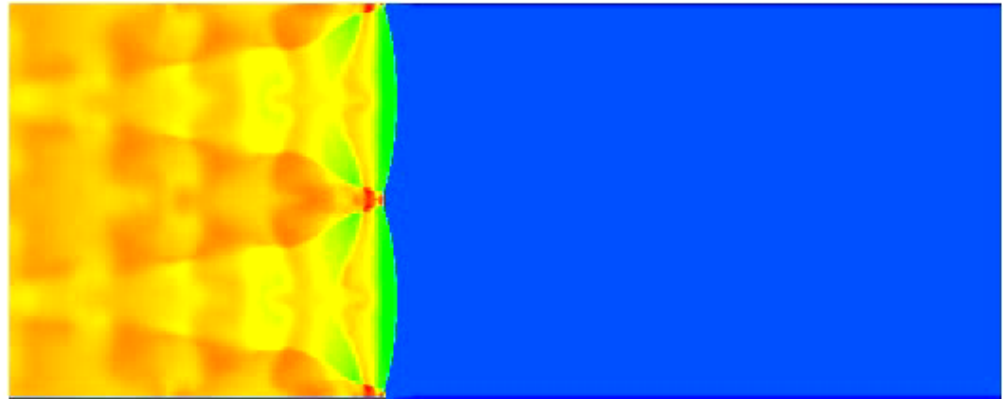
E: Reflected shock. F: Slip line. G: Diffusive extension of slip line.

Simulation of regular cellular structures

- Regular Chapman-Jouguet detonation for $H_2 : O_2 : Ar/2 : 1 : 7$ at $T_0 = 298K$ and $p_0 = 10$ kPa, cell width 1.6 cm.
- Perturb 1d solution with unreacted high-pressure pocket behind front
- Triple point trajectories by tracking $\max|\omega|$ on auxiliary mesh

$$\omega = \frac{\partial u_2}{\partial x_1} - \frac{\partial u_1}{\partial x_2}$$

- Adaptation criteria:
 - Scaled gradients of ρ and p
 - Error estimation in Y_i by Richardson extrapolation
- 67.6 Pts within induction length. 4 additional refinement levels (2,2,2,4).
- Similar configuration as E. Oran et al., J. Combustion and Flame 113, 1998.

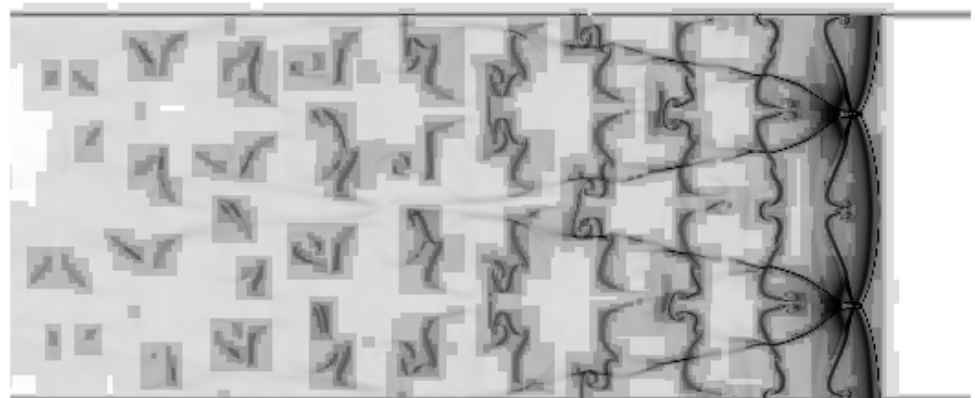
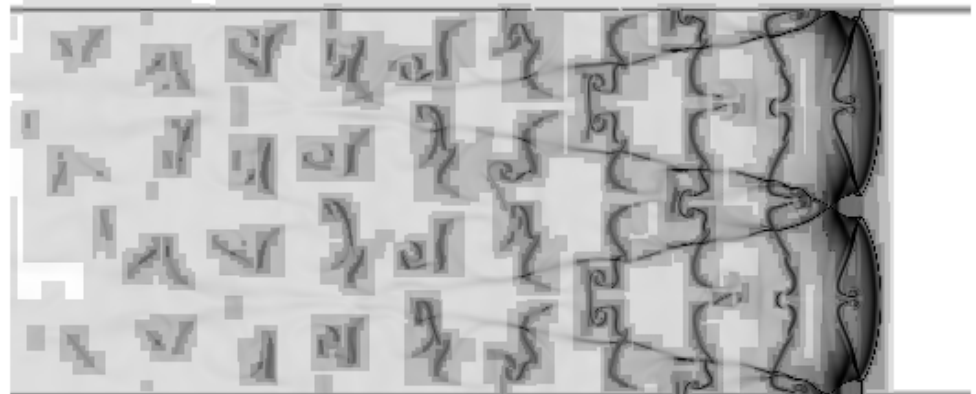


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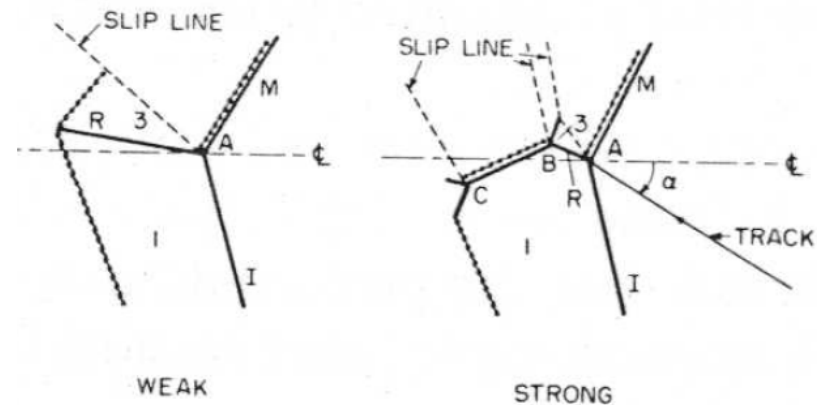
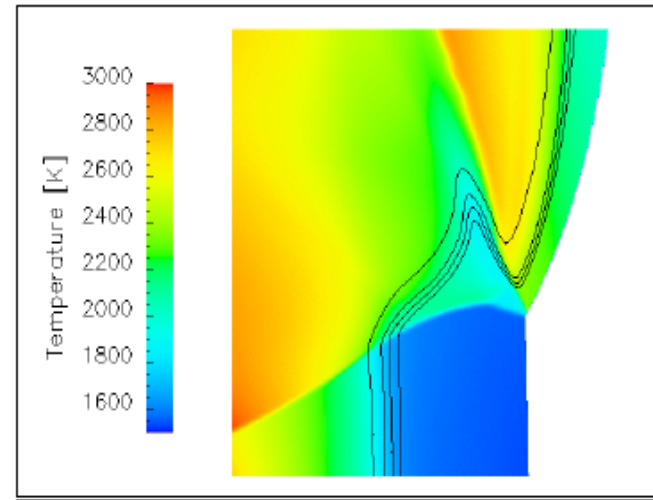
- Regular Chapman-Jouguet detonation for $\text{H}_2 : \text{O}_2 : \text{Ar}/2 : 1 : 7$ at $T_0 = 298\text{K}$ and $p_0 = 10\text{ kPa}$, cell width 1.6 cm.
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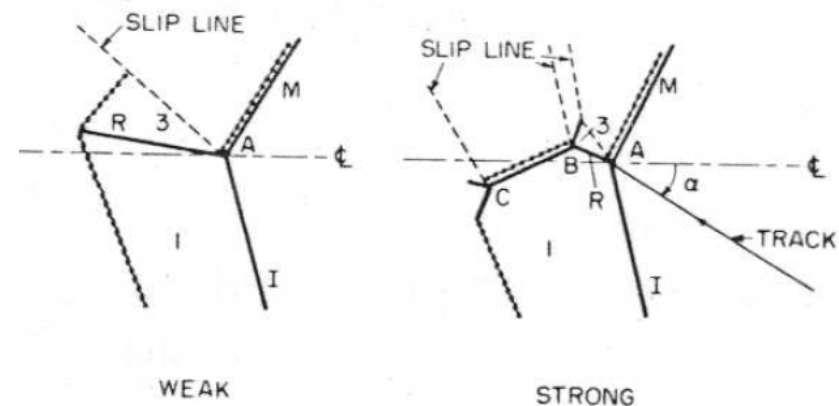
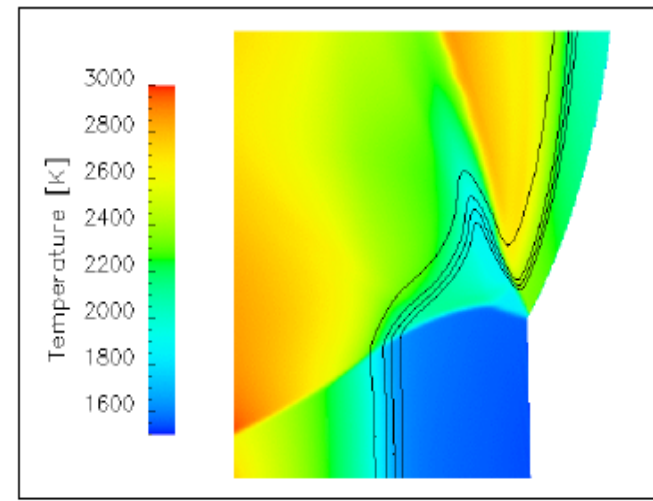
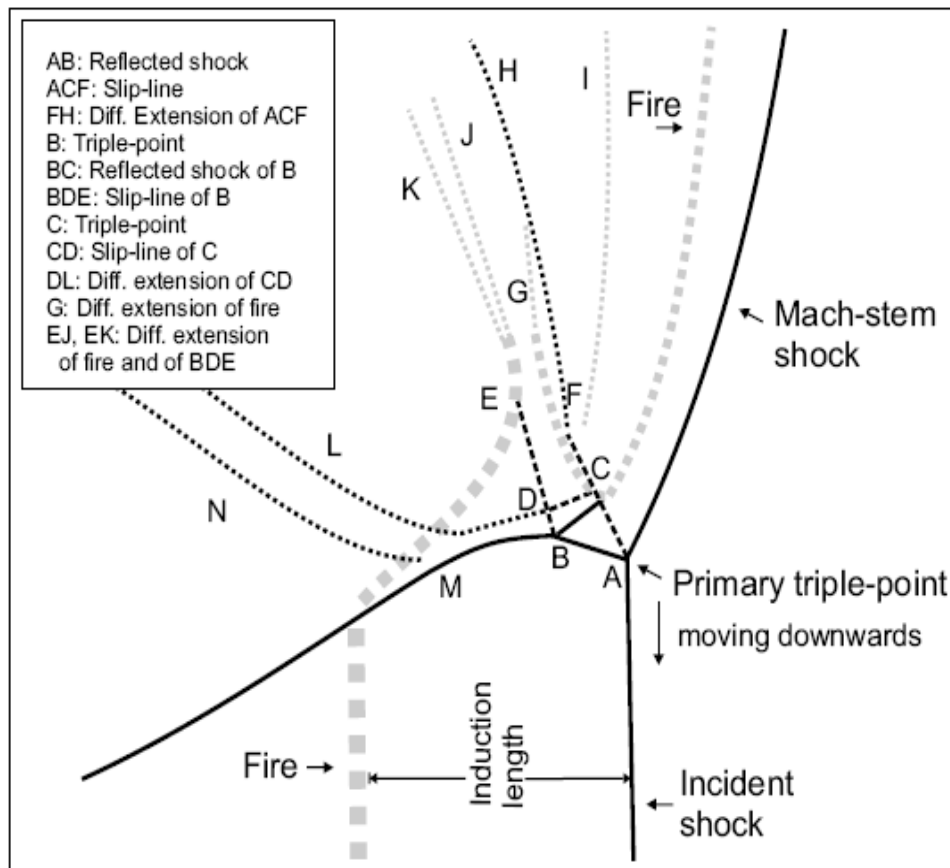
Flow Around a Triple-point



RD, Parallel adaptive simulation of multi-dimensional detonation structures, PhD thesis, BTU Cottbus, 2003

See also: Hu et al., The structure and evolution of a two-dimensional $H_2/O_2/Ar$ cellular detonation, Shock Waves, 2004.

Flow Around a Triple-point



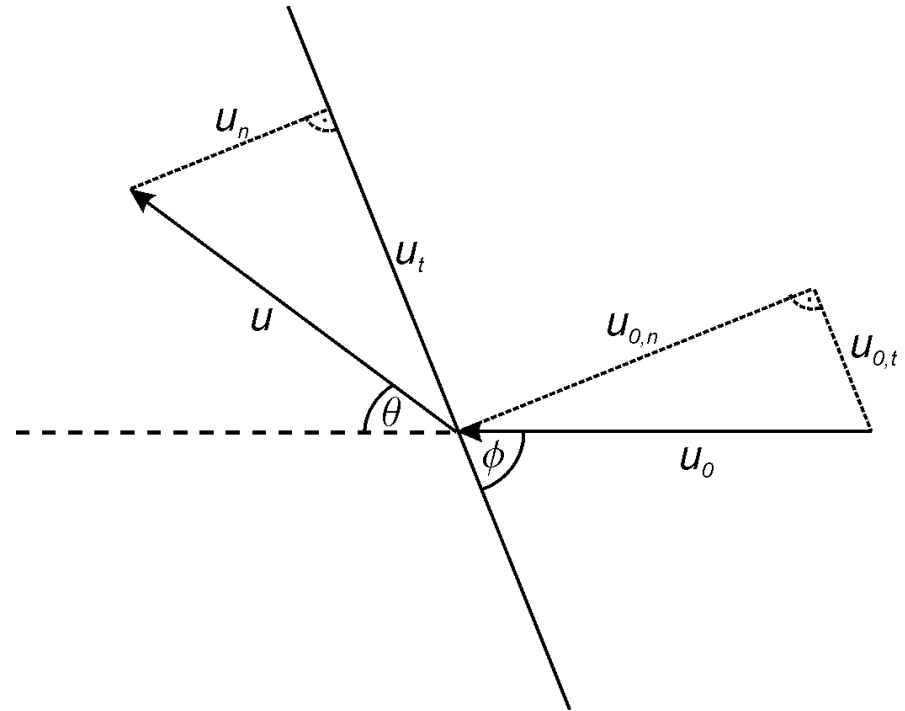
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Oblique shock relations

Apply Rankine-Hugoniot condition $\sigma(\mathbf{q} - \mathbf{q}_0) = \omega \cdot (\mathbf{f}(\mathbf{q}) - \mathbf{f}(\mathbf{q}_0))$
to steady shock wave with $\sigma=0$
described by the 2D Euler equations:

$$\begin{aligned}\rho_0 u_{0,n} &= \rho u_n \\ p_0 + \rho_0 u_{0,n}^2 &= p + \rho u_n^2 \\ u_{0,t} &= u_t \\ h_0 + \frac{1}{2} u_{0,n}^2 &= h + \frac{1}{2} u_n^2\end{aligned}$$



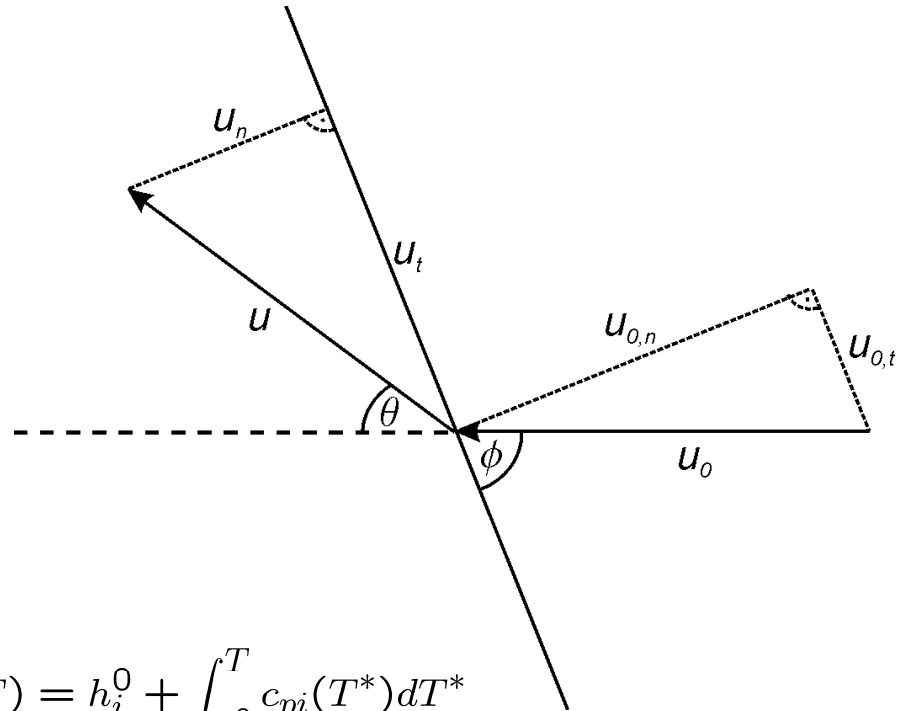
Commonly used:

$$\begin{aligned}\rho_0 u_0 \sin \phi &= \rho u \sin(\phi - \theta) \\ p_0 + \rho_0 u_0^2 \sin^2 \phi &= p + \rho u^2 \sin^2(\phi - \theta) \\ \rho_0 \tan \phi &= \rho \tan(\phi - \theta) \\ h_0 + \frac{1}{2} u_0^2 \sin^2 \phi &= h + \frac{1}{2} u^2 \sin^2(\phi - \theta)\end{aligned}$$

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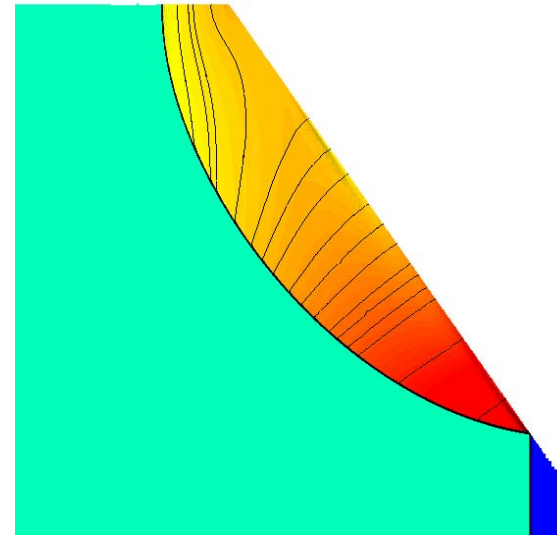
For thermally perfect mixtures with $h_i(T) = h_i^0 + \int_{T_0}^T c_{pi}(T^*) dT^*$
one solves

$$f(T) := \frac{RT_0}{u_{0,n}} + u_{0,n} - \frac{RT}{u_n} + u_n = 0 \quad \text{with} \quad u_n = \sqrt{u_{0,n}^2 - 2 \int_{T_0}^T c_p(\nu) d\nu}$$

numerically

Triple point configurations & Transition criteria

- Irregular reflection (IR) to regular reflection (RR): $M_B^T=1$ with $M_B^T>1$ for RR

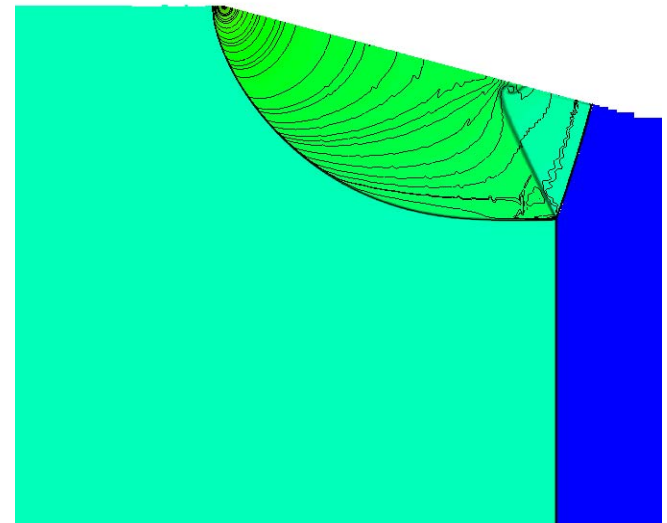
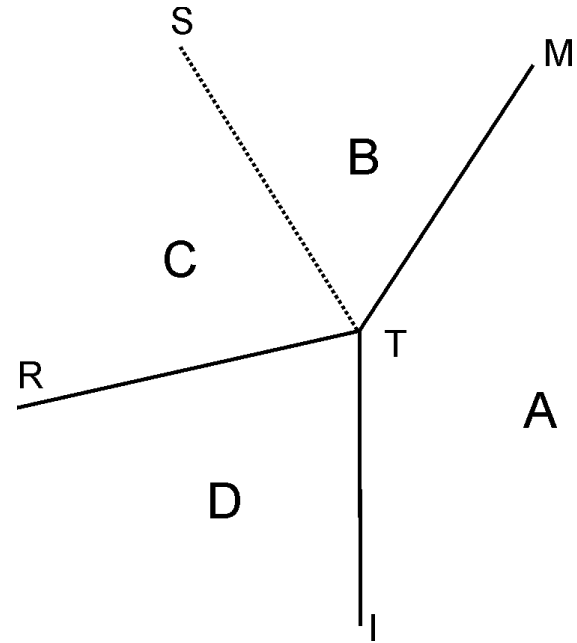


Triple point configurations & Transition criteria

- Irregular reflection (IR) to regular reflection (RR): $M_B^T=1$ with $M_B^T>1$ for RR

In the IR regime:

- von Neumann reflection (NR) to Mach reflection (MR): $M_D^T=1$ with $M_D^T>1$ for MR

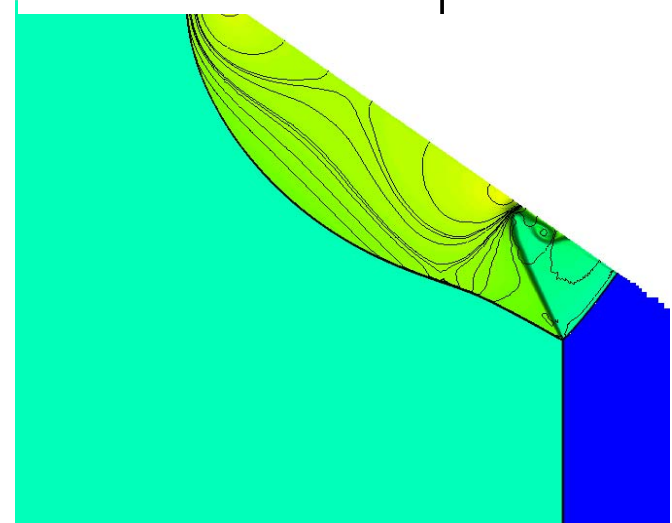
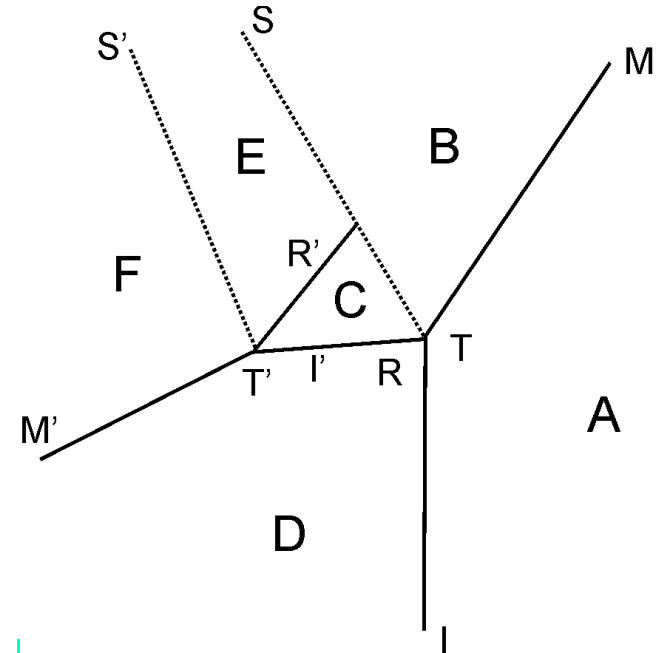


Triple point configurations & Transition criteria

- Irregular reflection (IR) to regular reflection (RR): $M_B^T=1$ with $M_B^T>1$ for RR

In the IR regime:

- von Neumann reflection (NR) to Mach reflection (MR): $M_D^T=1$ with $M_D^T>1$ for MR
- Single Mach reflection (SMR) to transitional or double Mach reflection (TMR/DMR): $M_C^T=1$ with $M_C^T>1$ for TMR/DMR



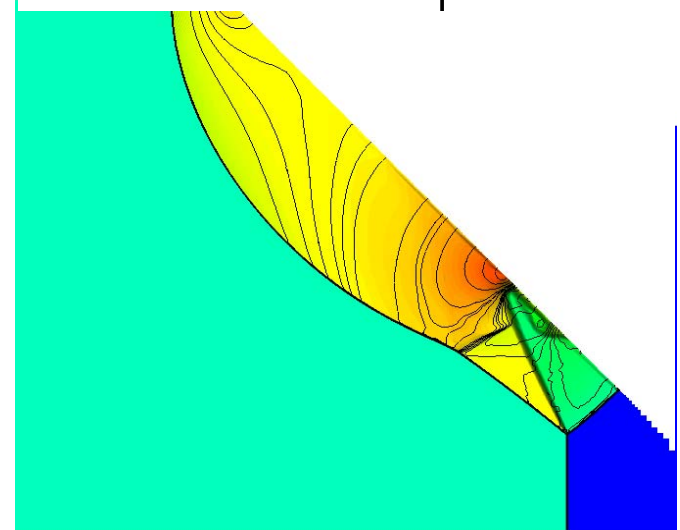
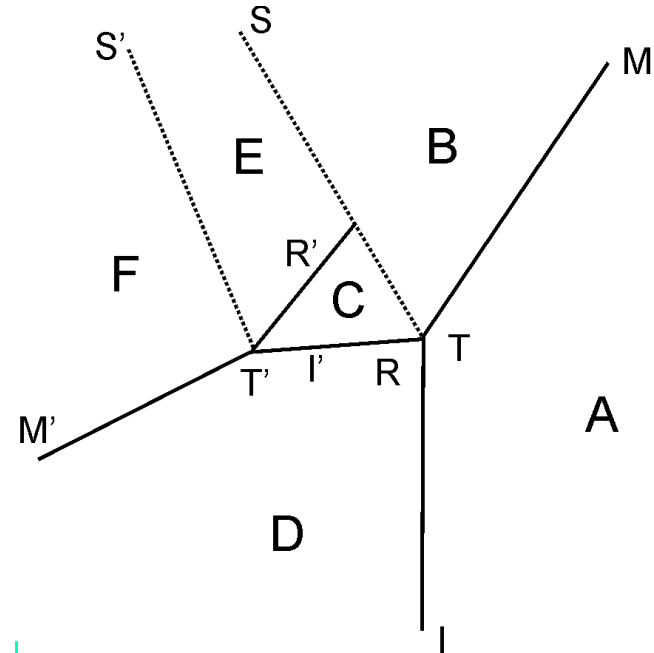
Triple point configurations & Transition criteria

- Irregular reflection (IR) to regular reflection (RR): $M_B^T=1$ with $M_B^T>1$ for RR

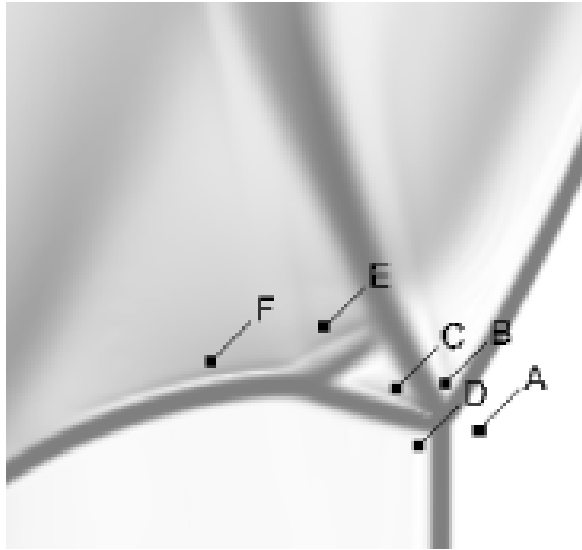
In the IR regime:

- von Neumann reflection (NR) to Mach reflection (MR): $M_D^T=1$ with $M_D^T>1$ for MR
- Single Mach reflection (SMR) to transitional or double Mach reflection (TMR/DMR): $M_C^T=1$ with $M_C^T>1$ for TMR/DMR
- Transitional (TMR) to double Mach reflection (DMR): $M_C^{T'}=1$ with $M_C^{T'}>1$ for DMR

G. Ben-Dor, Shock wave reflection phenomena, 2nd ed., Springer, 2007.

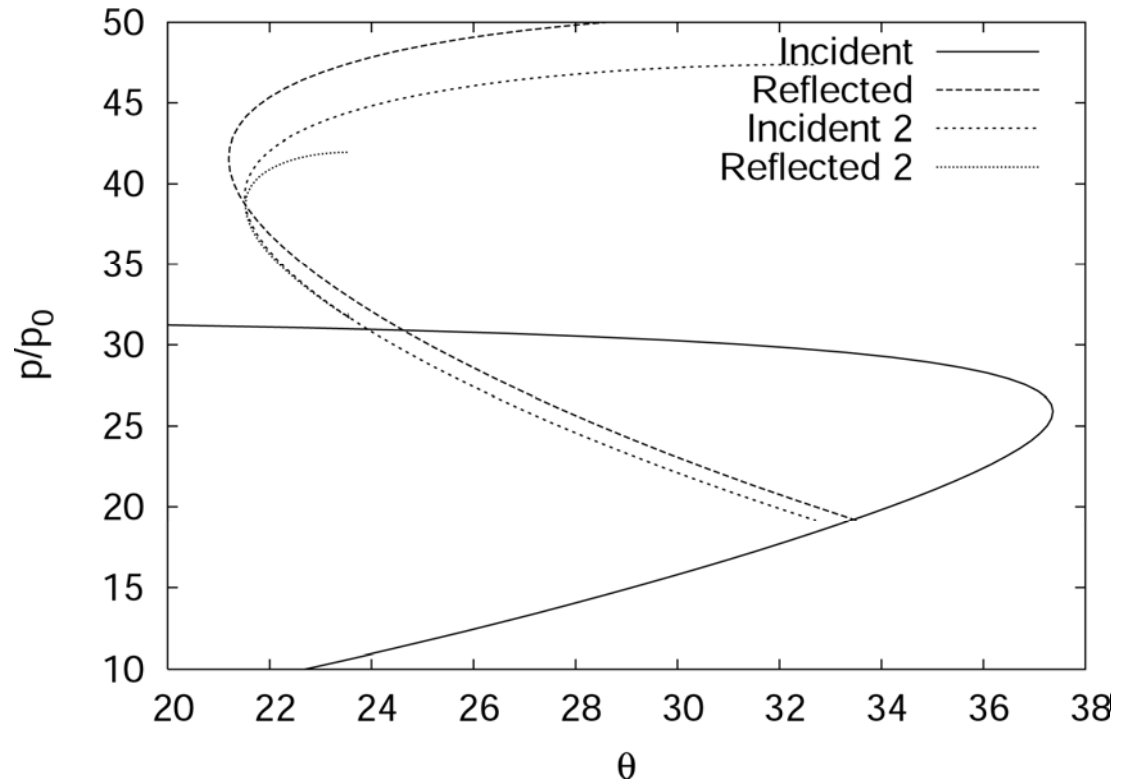
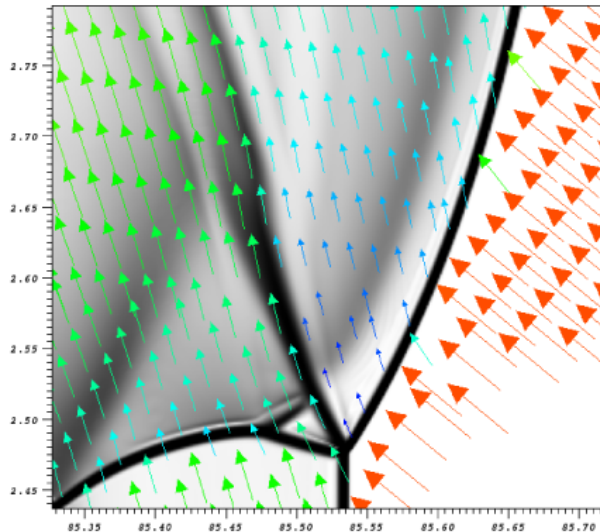


Shock polar analysis for a DMR



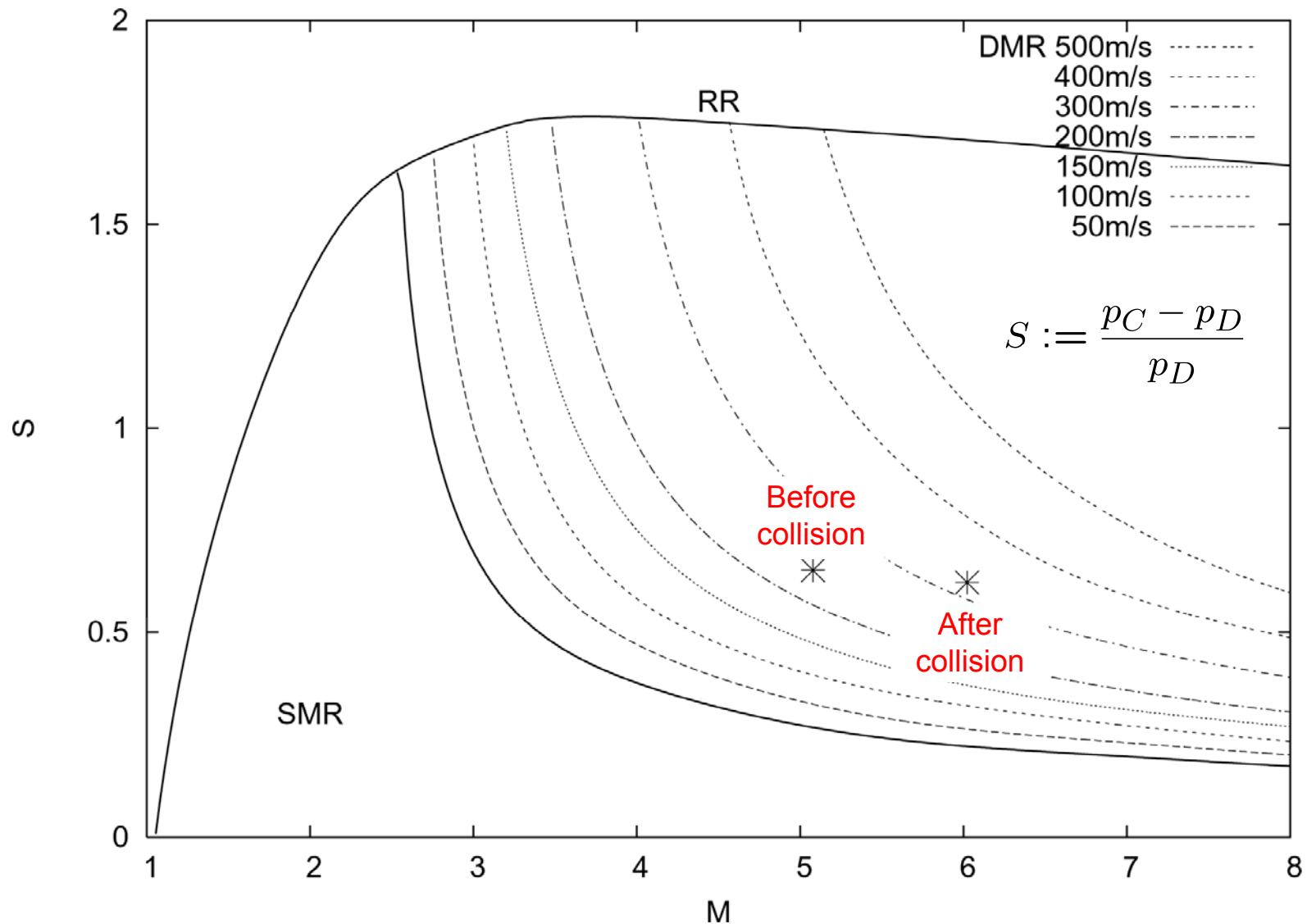
	p/p_0	r/r_0	u [m/s]	α	γ	M	u' [m/s]	α'	M'
A	1.00	1.00	1775	-39.5	1.557	5.078	1718	-40.2	4.916
B	31.45	4.17	447	-64.4	1.487	0.477	407	-70.3	0.433
C	31.69	5.32	965	-64.4	1.500	1.153	923	-67.0	1.103
D	19.17	3.84	1178	-73.2	1.509	1.533	1144	-75.7	1.488
E	35.61	5.72	901	-67.9	1.497	1.053	862	-70.8	1.007
F	40.61	6.09	777	-70.0	1.494	0.880	740	-73.6	0.838

$a_t \approx 60$ m/s used.



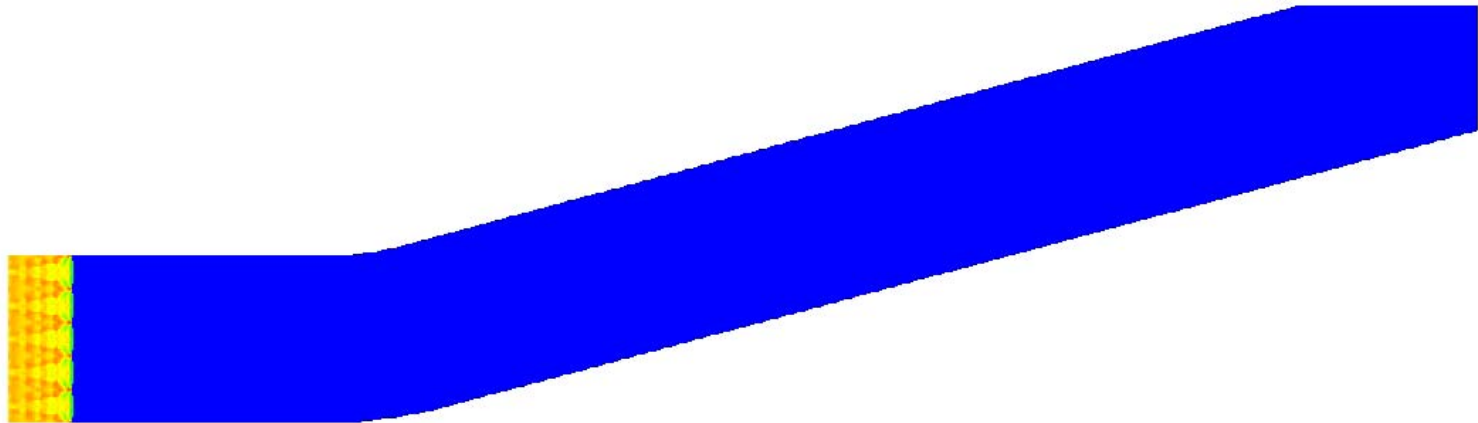
Reflection types depending on transverse wave strength

$\text{H}_2 : \text{O}_2 : \text{Ar}/2 : 1 : 7$ at $T_0 = 298\text{K}$ and $p_0 = 10\text{ kPa}$



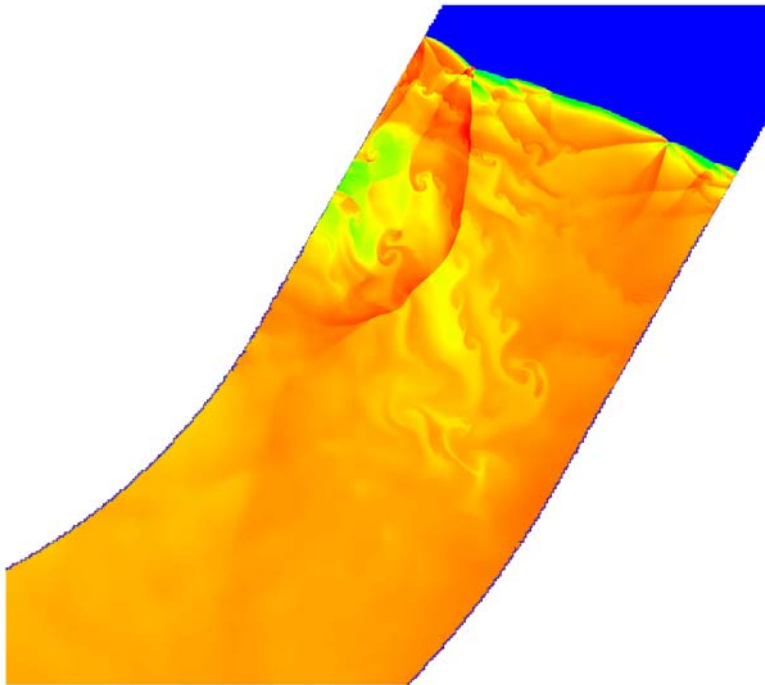
Transient conditions: Propagation through smooth pipe bends

- Regular Chapman-Jouguet detonation for $\text{H}_2 : \text{O}_2 : \text{Ar}/2 : 1 : 7$ at $T_0 = 298\text{K}$ and $p_0 = 10\text{ kPa}$, cell width 1.6 cm, tube width of 5 detonation cells (8 cm)
- Pipe bend with same radius. Angle: 15° , 30° , 45° , 60°
- 56.2 Pts within induction length. 4 additional refinement levels (2,2,2,4)
- Adaptive computations use $\approx 7 \cdot 10^6$ cells ($\approx 5 \cdot 10^6$ on highest level) instead of $1.2 \cdot 10^9$ cells (uniform grid)
- $\sim 70,000\text{h}$ CPU each on 128 CPUs Pentium-4 2.2GHz



Color plot of temperature, 15° pipe bend

Dynamic mesh adaptation (60° , $56.2 \text{ Pts}/l_{ig}$)

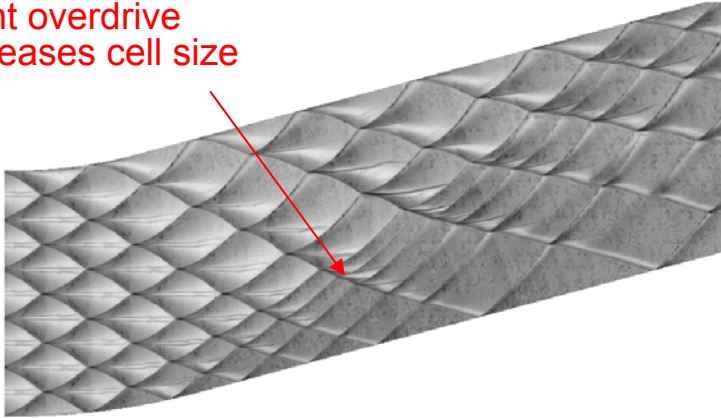


$\sim 170 \mu\text{s}$

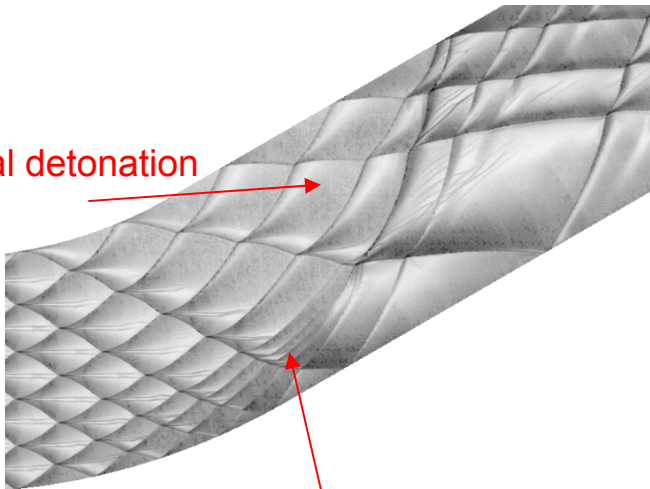
Time after entering bend

Enlarged triple point tracks (56.2 Pts/ l_{ig})

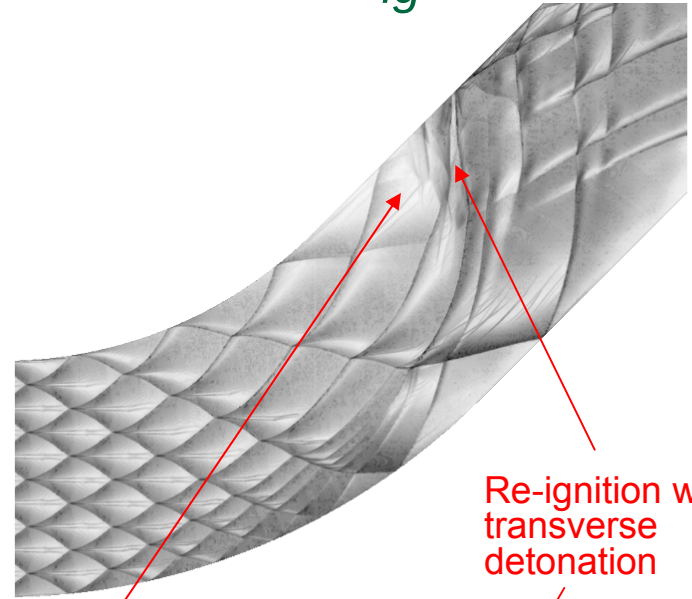
Slight overdrive
decreases cell size



Marginal detonation

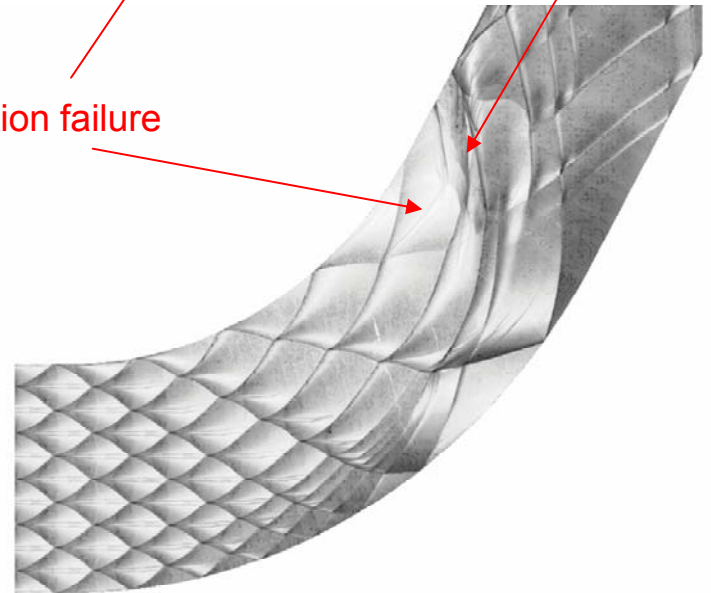


Triple point compression,
structure disappears

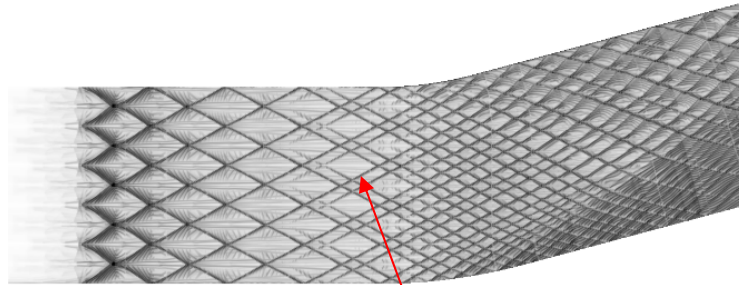


Re-ignition with
transverse
detonation

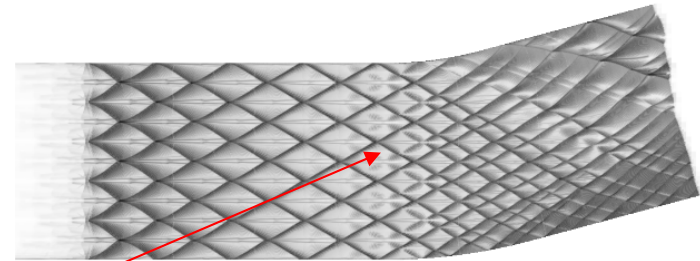
Detonation failure



Resolution comparison (15°)– triple point tracks

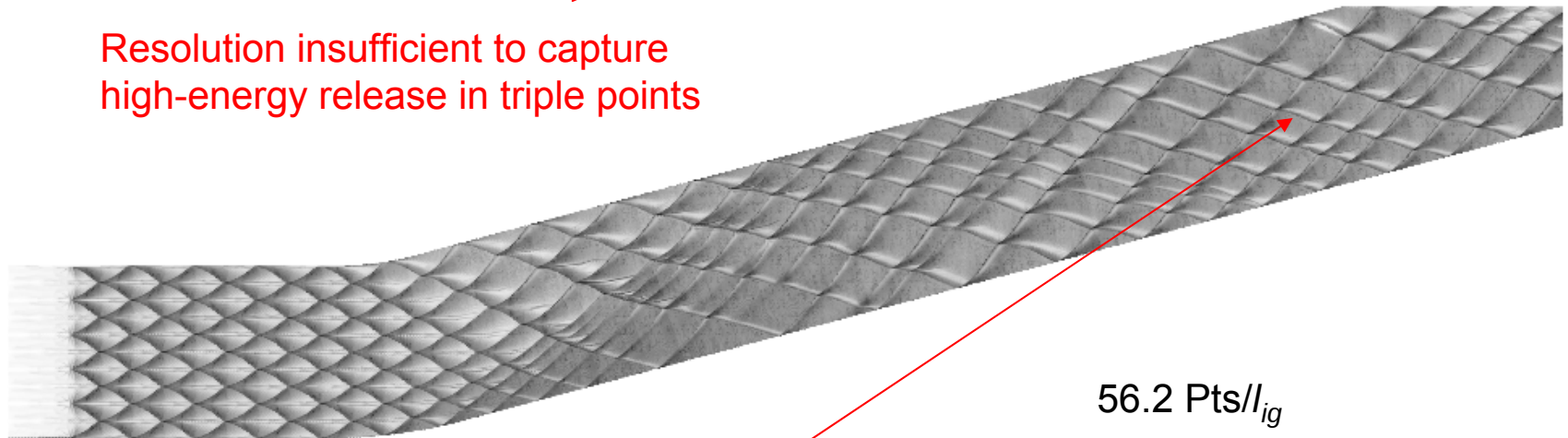


14.05 Pts/ l_{ig}



28.1 Pts/ l_{ig}

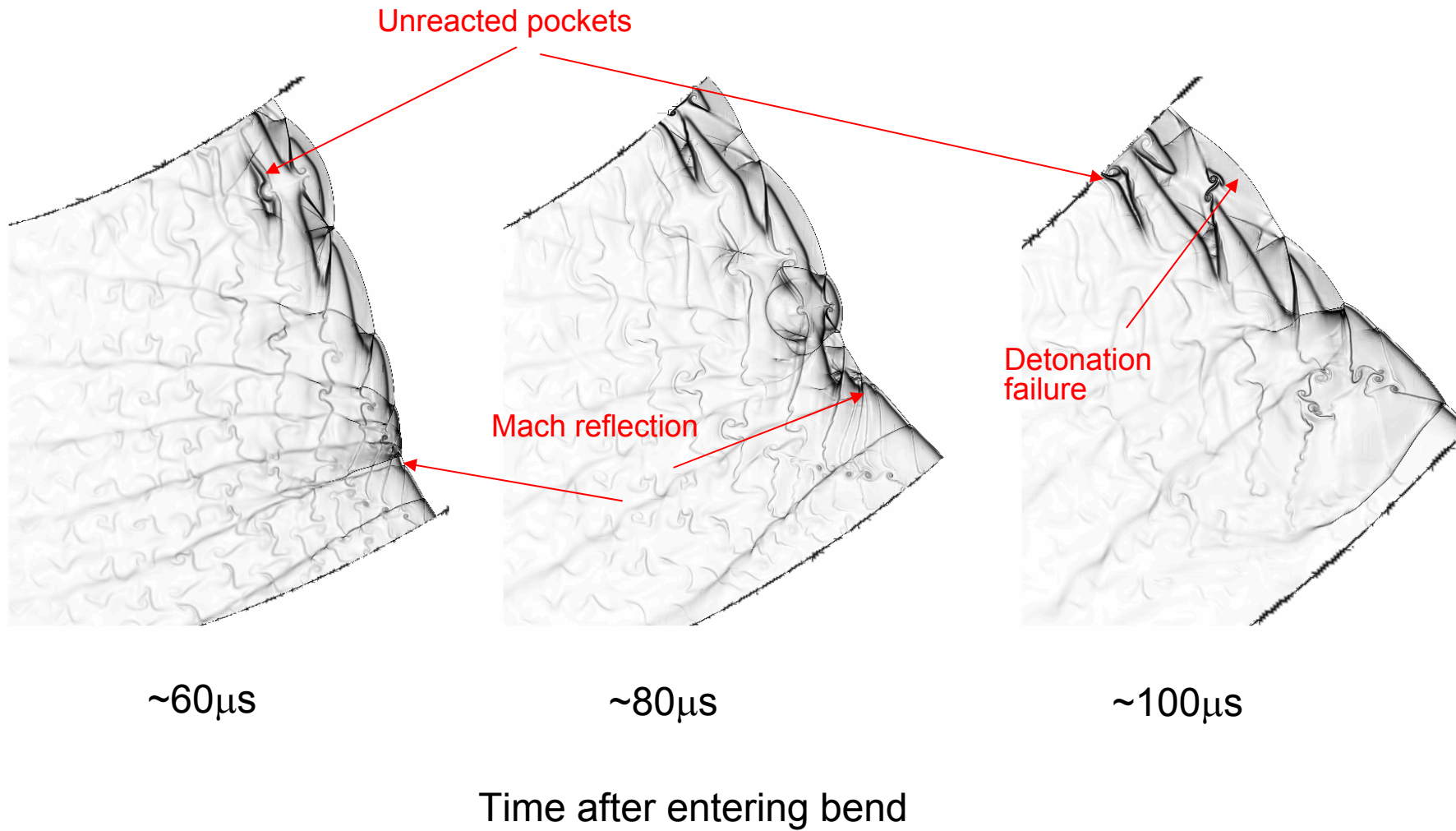
Resolution insufficient to capture
high-energy release in triple points



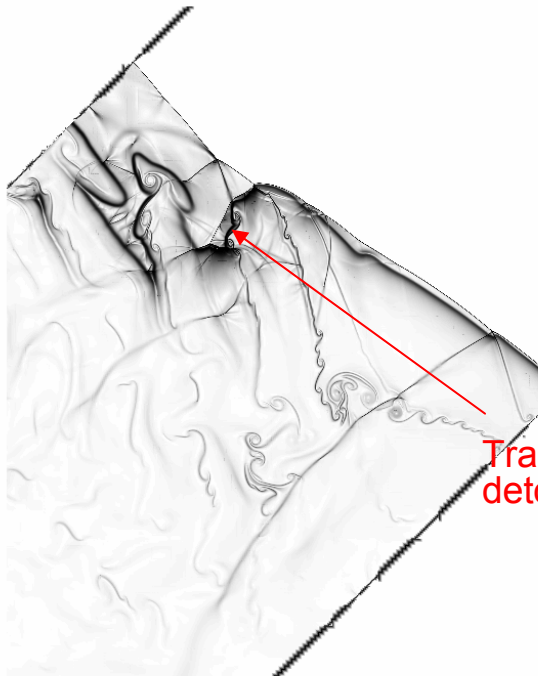
56.2 Pts/ l_{ig}

Resolution sufficient: Number of triple
points becomes approximately the
same as before bend

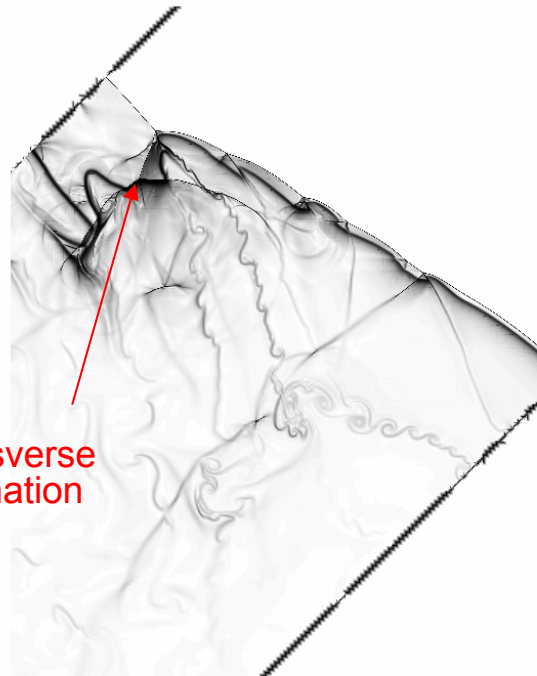
Principal flow phenomena (45° , $56.2 \text{ Pts}/l_{ig}$)



Principal flow phenomena (45° , $56.2 \text{ Pts}/l_{ig}$)



$\sim 120 \mu\text{s}$

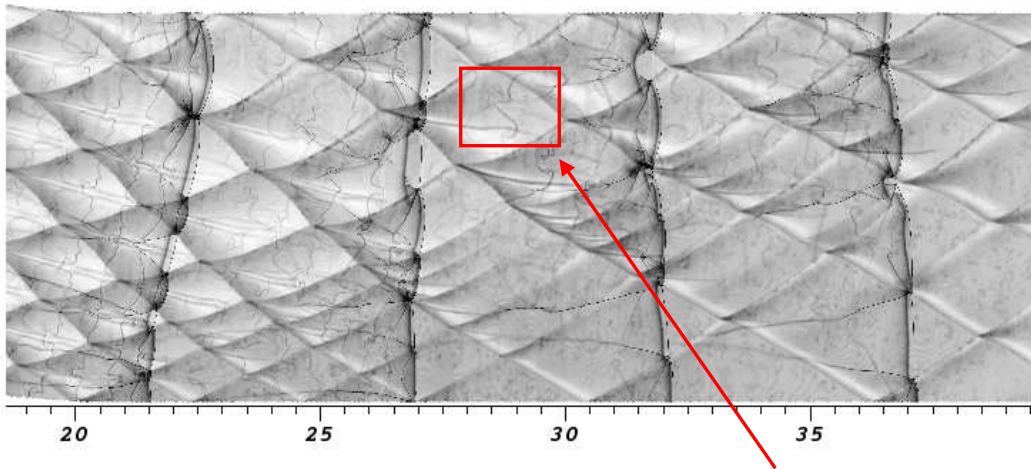


$\sim 130 \mu\text{s}$

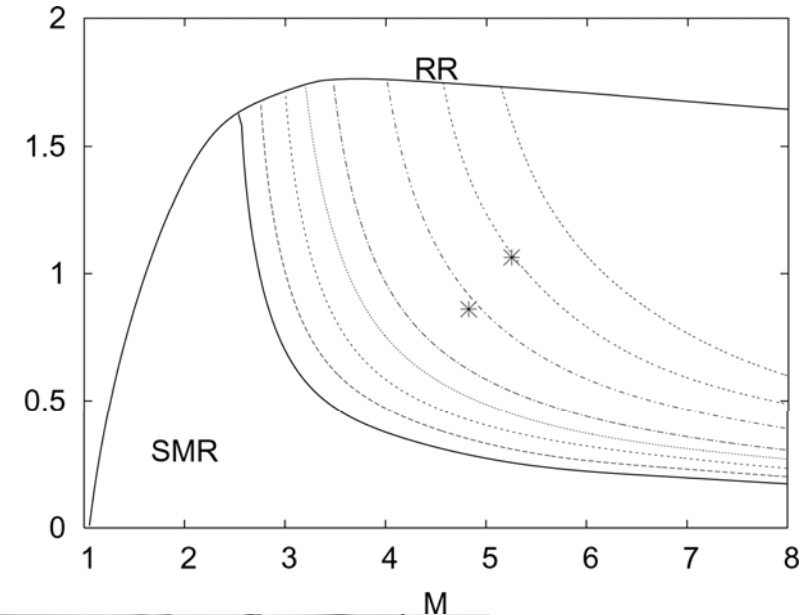
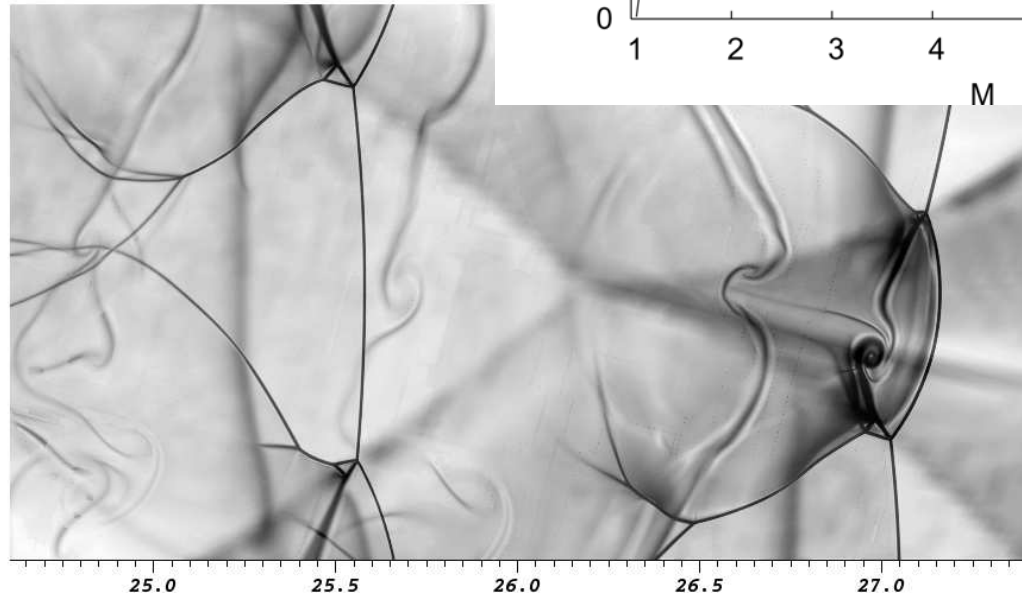


$\sim 140 \mu\text{s}$

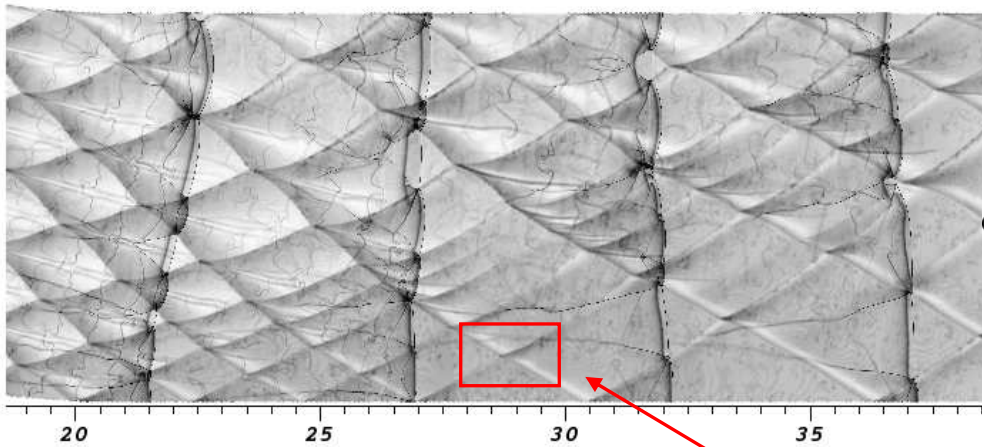
Triple point analysis (15°)



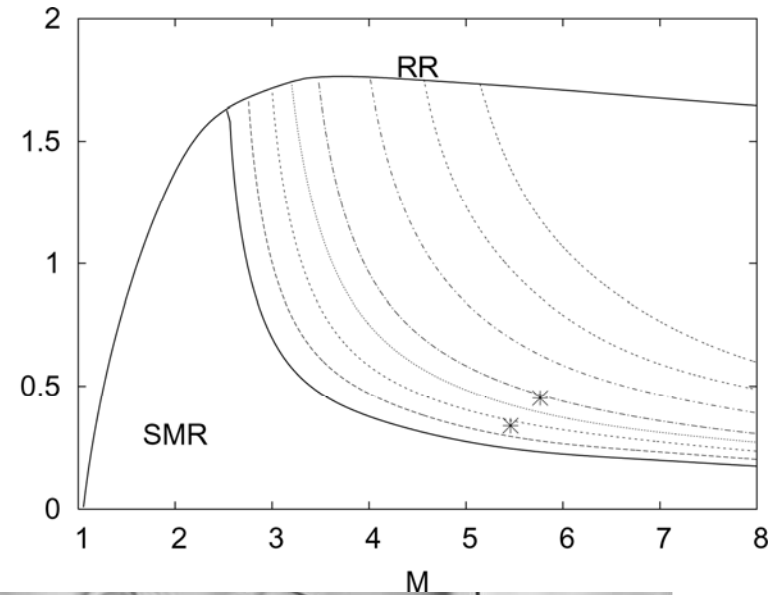
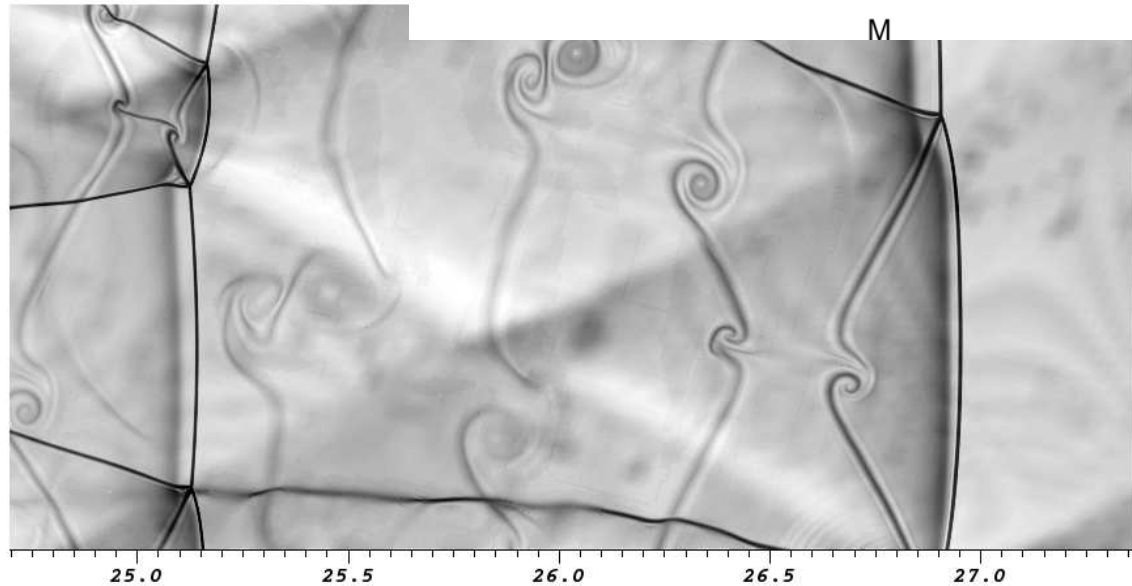
Strengthening of
double Mach
reflection pattern,
trajectory angle
increases



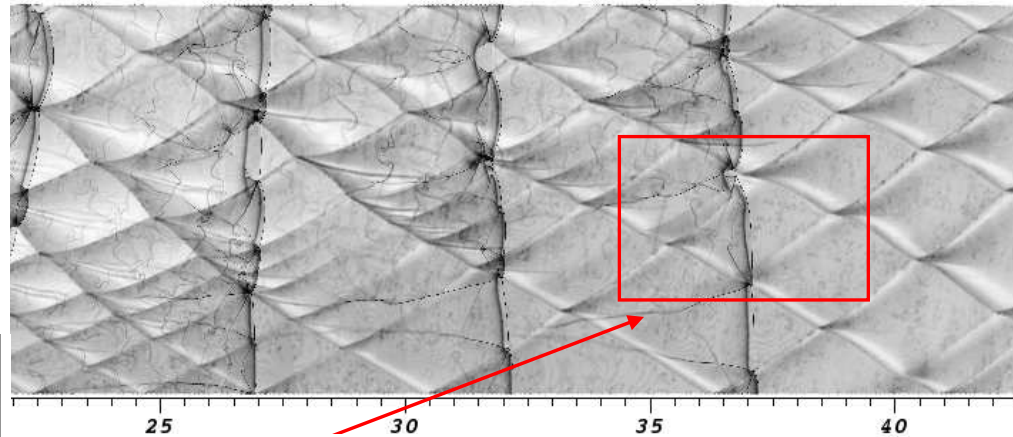
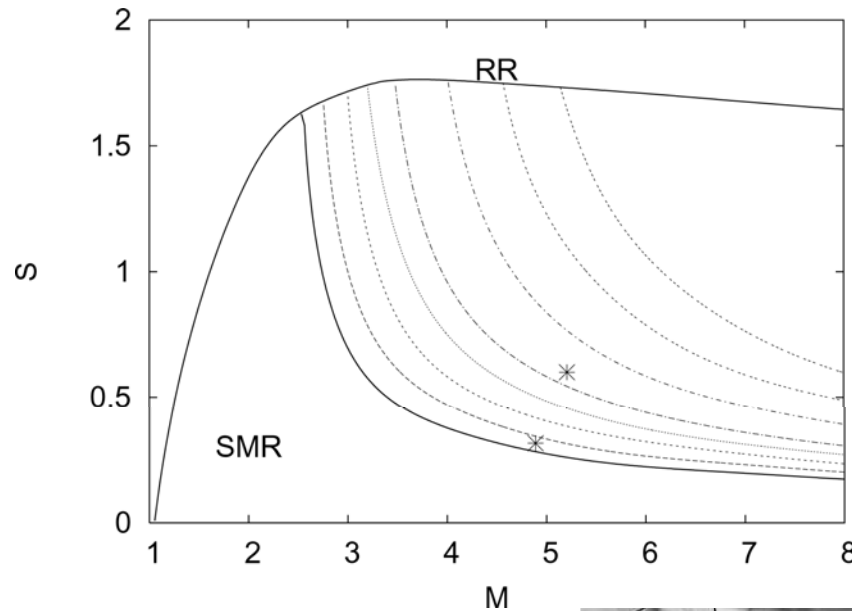
Triple point analysis (15°)



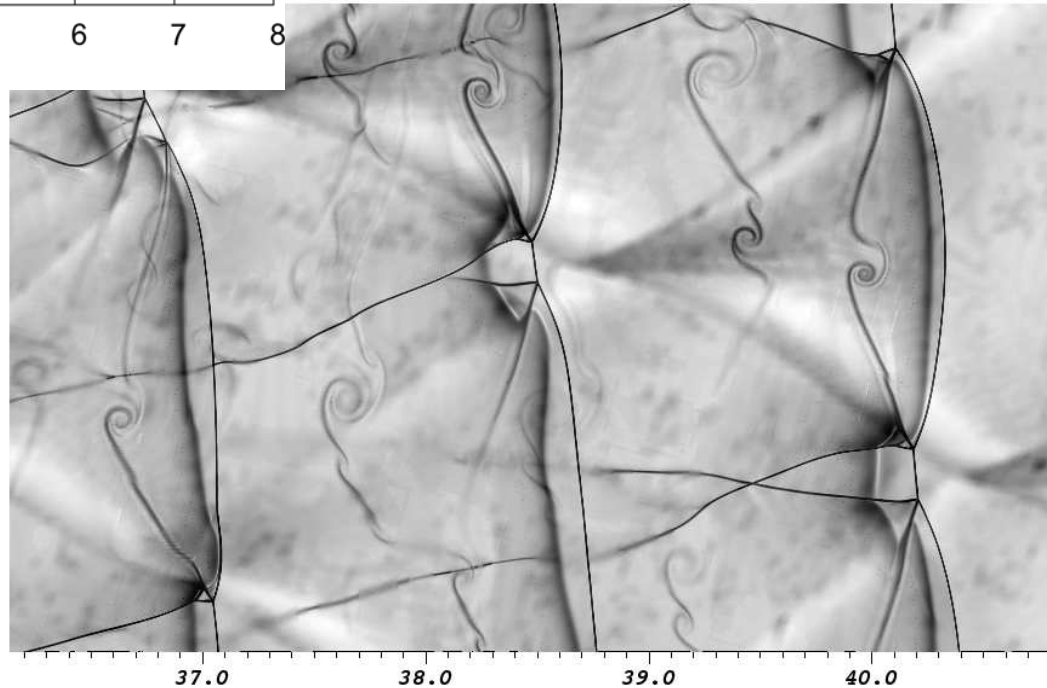
Transitional Mach
reflection pattern,
Straight trajectories



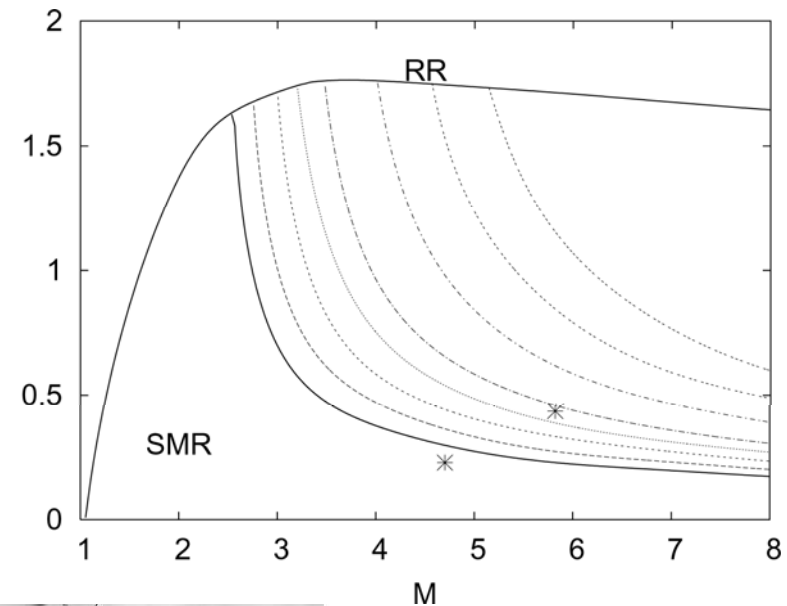
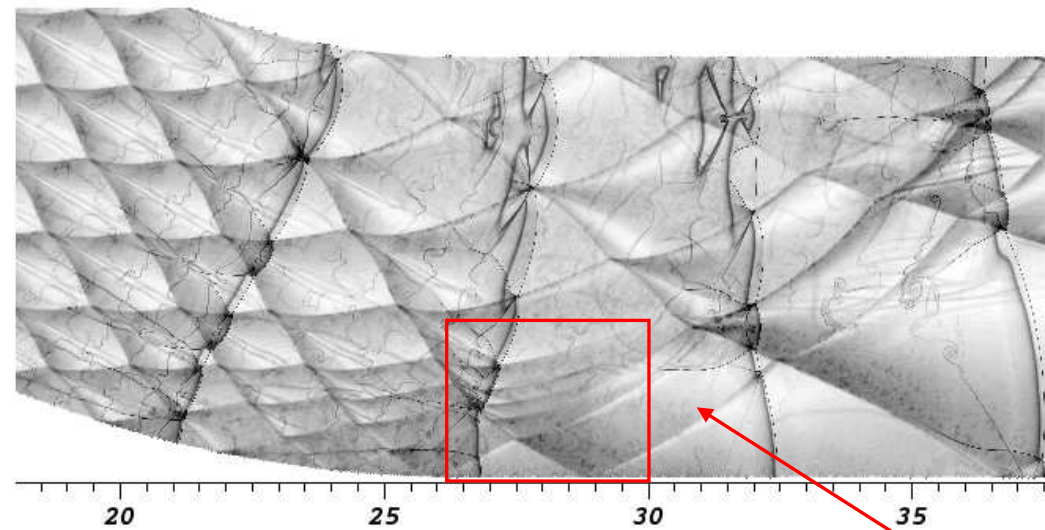
Triple point analysis (15°)



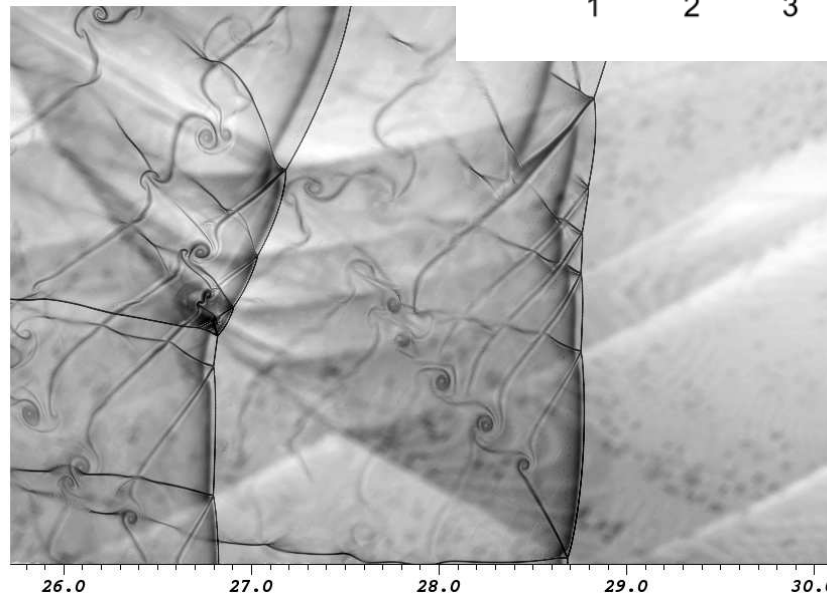
Formation of
transverse wave with
transitional Mach
reflection, double
Mach reflection after
collision



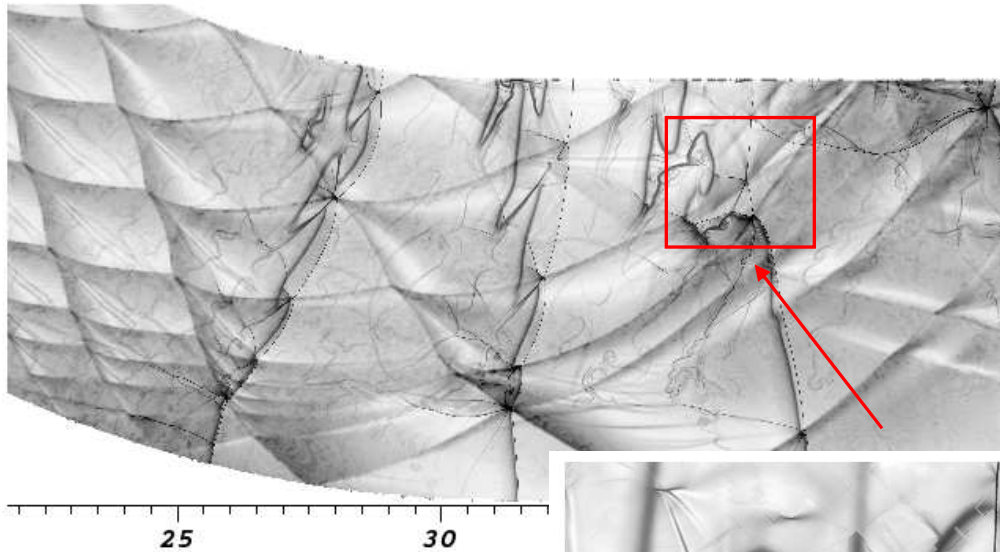
Triple point analysis (30°)



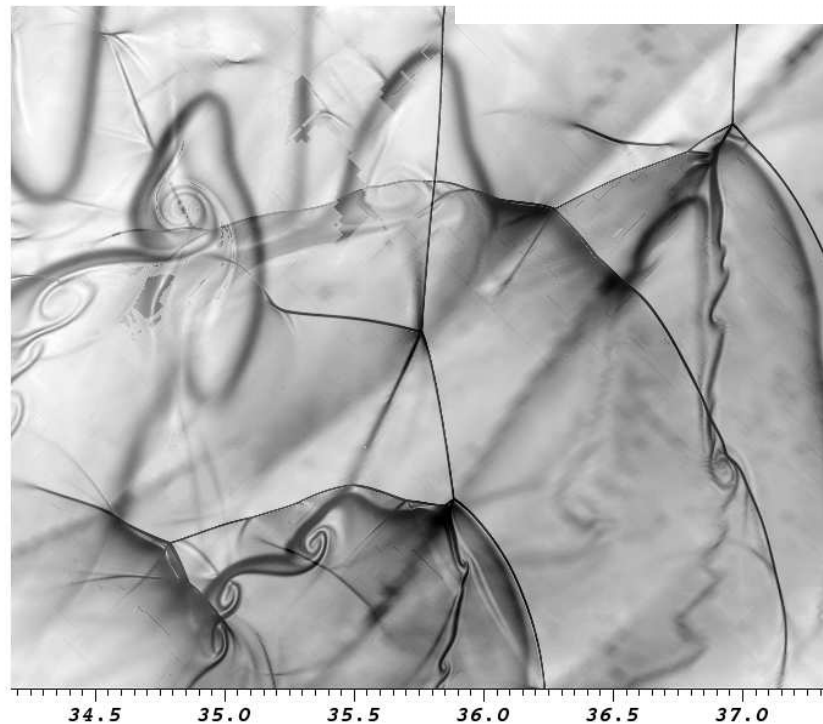
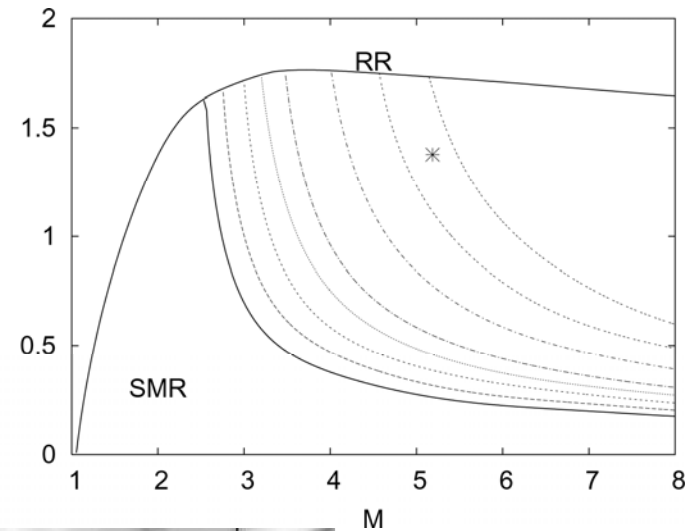
Triple point
quenching and failure
with weakening
transitional Mach
reflection patterns



Triple point analysis (45°)

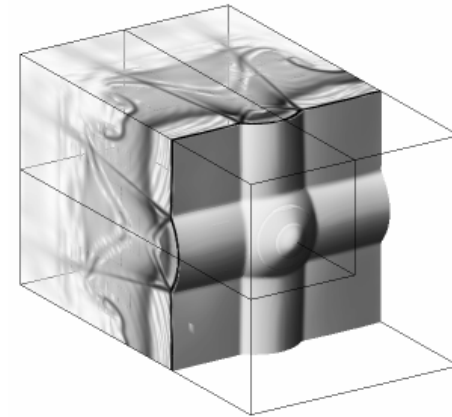
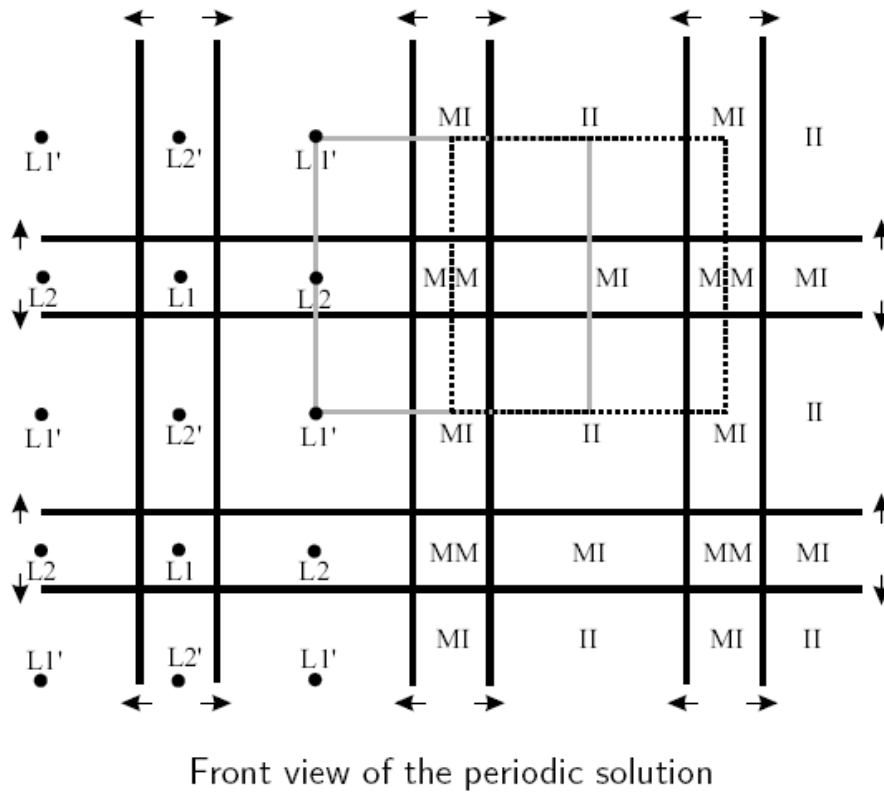


Transverse
detonation with
strong double Mach
reflection pattern,
triple point on
transverse wave

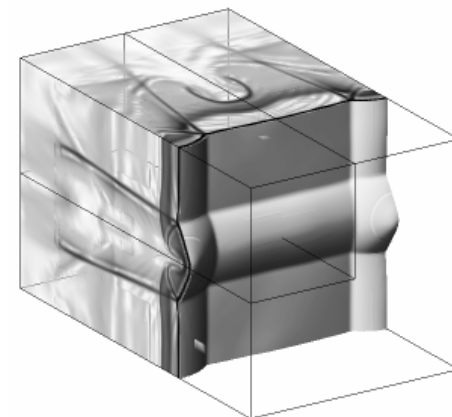


Outlook: Regular cellular structures in 3D

- Regular Chapman-Jouguet detonation for $\text{H}_2 : \text{O}_2 : \text{Ar}/2 : 1 : 7$ at $T_0 = 298\text{K}$ and $p_0 = 6.67 \text{ kPa}$, cell width 3 cm
- Unburned gas flows in with CJ velocity



$t = 680 \mu\text{s} + 600 \mu\text{s}$ (Computation 1)

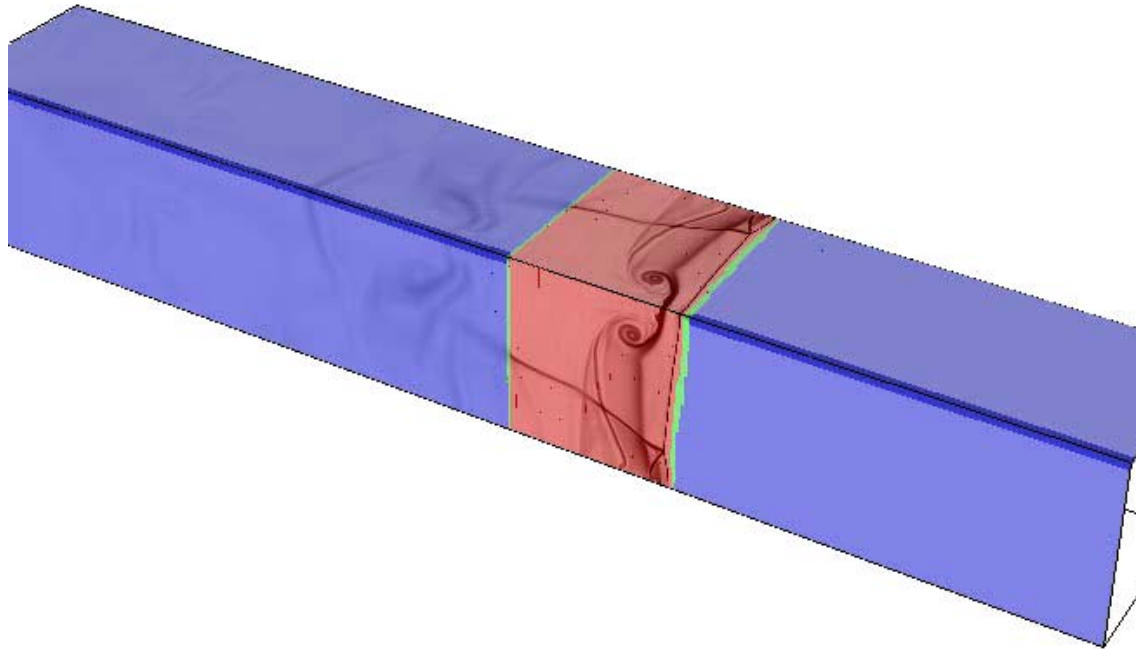


$t = 660 \mu\text{s} + 620 \mu\text{s}$ (Computation 2)

High-resolution simulation

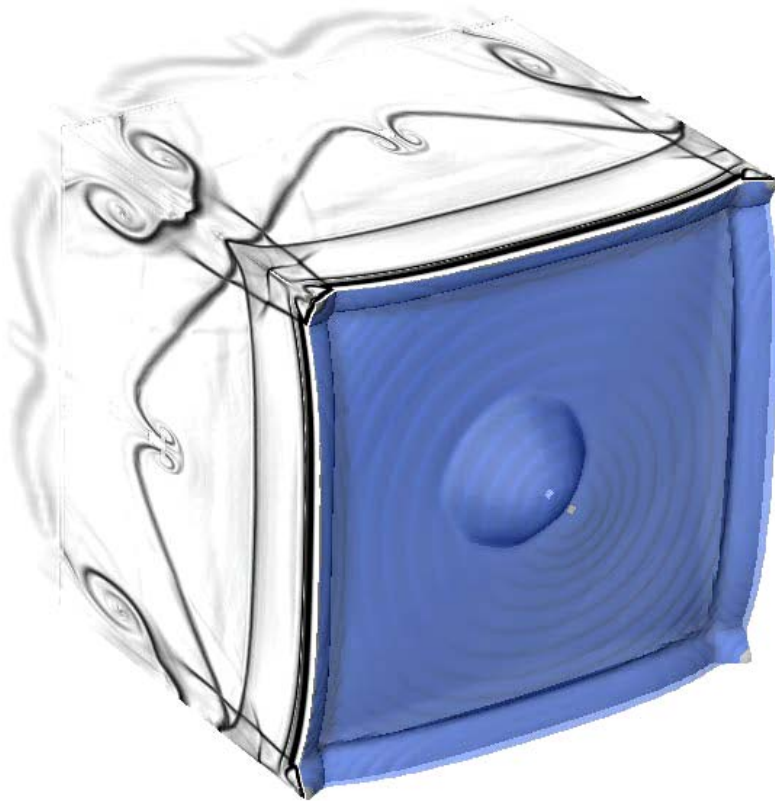
- Simulation of only one quadrant
- 44.8 Pts within induction length
- AMR base grid 400x24x24, 2 additional refinement levels (2, 4)
- Simulation uses ~18M cells instead of 118M (unigrid)
- ~51,000h CPU on 128 CPU Compaq Alpha (LANL QSC)

Task	&
Fluid Dynamics	37.6
Chemical Kinetics	25.1
Boundary Setting	24.4
Reorganization	6.6
Misc.	6.3



Schlieren plot of density of
refinement domains

High-resolution simulation: Results



Schlieren plot of Y_{OH} , iso-surfaces of Y_{OH} and ρ visualize induction length, periodicity exploited for visualization



Transverse wave strength S smaller than in 2D. TMR patterns do occur!

Conclusions

- For particular mixtures, detailed detonation structure simulations with detailed chemistry are possible nowadays in 2D realistic geometries
- Accurate studies for idealized 3D configurations
- Resolution down to the scale of secondary triple points can be provided on parallel capacity computing systems
 - *Key components:*
 - *Operator splitting allows a cell-wise integration of stiff reaction terms*
 - *SAMR provides a sufficient spatial and temporal resolution, savings up to >250*
- Unreactive, thermally perfect shock polar analysis is applicable to explain observed reflection patterns
 - *Shock wave reflection theory is applicable to predict local triple point structure and stability*
 - *Triple point type is determined solely by S and M which can be derived from a single time step*
 - *Still missing: estimate for secondary triple point velocity a_t tailored for detonations to rigorously TMR/DMR transition*
- Observations:
 - *Stable triple point structures in self-sustained detonations seem to exist only in the TMR and DMR, but not in the SMR regime*
 - *A change of the reflection type happens especially in triple point collisions*
- Literature, links to software, papers, etc.: <http://www.csm.ornl.gov/~r2v>